Polymorphic Layered MoTe₂ from Semiconductor, Topological Insulator, to Weyl Semimetal

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ABSTRACT: Large size (~2 cm) single crystals of layered MoTe₂ in both 2H- and 1T'-types were synthesized using TeBr₄ as the source of Br₂ transport agent in chemical vapor transport growth. The crystal structures of the as-grown single crystals were fully characterized by X-ray diffraction, Raman spectroscopy, scanning transmission electron microscopy, scanning tunneling microscopy (STM), and electrical resistivity (ρ) measurements. The resistivity ρ(T), magnetic susceptibility χ(T), and heat capacity Cᵥ(T) measurement results reveal a first order structural phase transition near ~240 K for 1T'-MoTe₂, which has been identified to be the orthorhombic Td-phase of MoTe₂ as a candidate of Weyl semimetal. The STM study revealed different local defect geometries found on the surface of 2H- and Td-types of MoTe₆ units in trigonal prismatic and distorted octahedral coordination, respectively.

INTRODUCTION

The two-dimensional (2D) layered transition metal dichalcogenides (TMDs), MX₂ (M = Ta, Mo, or W and X = Se, S, or Te), have attracted significant attention due to their rich physical properties including superconductivity, large thermoelectric effect, and anomalous magnetoresistance.²⁻⁵ These 2D materials have shown promising potential for next-generation devices as an alternative to graphene as a 2D monolayer material.⁶,⁷ In search of materials having both a band gap like a typical semiconductor and a high charge carrier mobility like graphene, the recently demonstrated MoTe₂-based high mobility device opened a new area for future applications.⁸

The recent theoretical prediction and following discovery of WTe₂ as a topological type-II Weyl semimetal⁹ with extremely large magnetoresistance⁹ have triggered extensive attention to uncover the origin of these intriguing emerging physical properties identified in TMDs, which arise from rich variations of intralayer packing of MXₓ octahedral or trigonal prismatic units and interlayer stacking via van der Waals (vdW) interlayer interactions.¹⁰⁻¹¹ The two most common structures of TMDs are the 2H-type (space group P6₃/mmc with MX₆ in trigonal prismatic coordination) and the 1T-type (space group P3m1 with edge-shared MX₆ in octahedral coordination).¹² In addition to the common 2H-type structure, MoTe₂ has also been found in a variation of 1T-type (α), namely the 1T'-type (β), which is composed of buckled layers of edge-sharing MoTe₆octahedra in monoclinic symmetry of space group P2₁/m.¹³⁻¹⁴ Although the name 1T'-type has been discussed theoretically in parallel to those of 1T- and 1H-types by assuming there is one layer per unit cell,¹⁵ the experimentally identified two layers per unit cell following the space group P2₁/m operation suggests that the correct categorization for the recently discussed 1T'-type should be named 2T'-type following the convention of TMDs.¹⁶ So as not to be confused with those discussed in recent publications, we will continue to use the term 1T' in this paper.

Recent theoretical investigations showed that the separation of Weyl points as well as the length of a Fermi arc can be...
adiabatically tuned as a function of Mo-doped WTe$_2$,\textsuperscript{17} and the orthorhombic Td phase of pristine MoTe$_2$ also possesses type-II Weyl nodes,\textsuperscript{18} which have much larger Weyl node pair separations than those found in WTe$_2$.\textsuperscript{9,18} Hence, Td-MoTe$_2$ is expected to provide more convincing experimental evidence via angle-resolved photoemission spectroscopy compared to that of WTe$_2$. On the basis of the transport studies of single crystal with the aim of growing large and high quality MoTe$_2$ single crystals of controlled structure types, we have demonstrate the successful growth of polymorphic single crystals using a novel chemical vapor transport (CVT) method that involves a mixture of Mo and TeBr$_4$ as the precursor, where TeBr$_4$ decomposed at 420 °C becomes the natural source of Br$_2$ as the transport agent in the CVT growth. By controlling the temperature profile of the source, large crystals (∼2 cm) of MoTe$_2$ have been grown successfully in less than a week. Details of the CVT growth method as well as a comprehensive characterization of the as-grown single crystals are presented, including X-ray diffraction (XRD), Raman spectroscopy, scanning transmission electron microscopy (STEM), scanning tunneling microscopy (STM), and transport measurements.

### EXPERIMENTAL SECTION

#### Synthesis and Crystal Growth.

The chemical vapor transport method was employed using TeBr$_4$ as Br$_2$ source, which acts as the transport agent and allows an effective and faster vapor transport to produce the necessary supersaturation of the expected final product. A three-zone muffle furnace was fabricated for this purpose, having typical temperature profiles for the growth of single crystals of 2H-MoTe$_2$ and 1T-MoTe$_2$, as shown in Figure S1. A stoichiometric amount of Mo/Te = 1:2 (6N purity for Te and 5N for Mo) was sealed into an evacuated quartz ampule and heated for 2 days at 750 °C. Approximately 10 g of the prereacted MoTe$_2$ powder was placed together with a variable amount of TeBr$_4$ (purity 4N) (120 mg) at one end of the silica ampule (length of 40 cm with inner diameter of 1.8 cm and outer diameter of 2.0 cm). Bromine concentration in the range of 2.6–3.7 mg was yielded sufficiently from ∼5 mg of TeBr$_4$ per cm$^3$, which provided high transport rates of nearly 150 mg per day. All preparation steps before the quartz tube evacuated flame sealing were carried out in an argon-filled glovebox of oxygen, and the water level was kept below ∼1 ppm. The loaded ampule was evacuated and flame sealed before loading into the tube furnace for CVT growth. For the growth of 1T-MoTe$_2$ single crystals, the end of the ampule containing the prereacted material was held at 1050 °C, and the growth end was maintained at a temperature near 950 °C with a temperature gradient near 2.5 °C/cm for a week. Shiny 1T-MoTe$_2$ single crystals of sizes up to 17 × 10 × 2 mm$^2$ were obtained, as shown in the inset of Figure 1(a). For the growth of 2H-MoTe$_2$ single crystals, the prereacted material was held at 800 °C, and the crystals were grown at the end maintained at a temperature of 750 °C with a temperature gradient near 1.25 °C/cm for approximately a week. Shiny 2H-MoTe$_2$ single crystals of sizes up to 21 × 16 × 2 mm$^2$ were obtained, as shown in the inset of Figure 1(b).

The crystal structure and phase purity of the samples were checked by synchrotron X-ray powder diffraction (SXRD) using a wavelength of λ = 0.619927 Å (BL01C2, NSRRC, Taiwan). Raman spectroscopy was performed in the backscattering configuration using 632 nm excitation lasers. The field-cooled (FC) and zero-field-cooled (ZFC) magnetization curves were measured in a commercial vibrating sample magnetometer (VSM, Quantum Design, USA) in the presence of an applied magnetic field of 10 kOe. The heat capacity ($C_p$) and transport measurements were carried out using the physical properties measurement system (PPMS, Quantum Design, USA).

### RESULTS AND DISCUSSION

The room-temperature SXRD patterns of the pulverized as-grown MoTe$_2$ single crystal are shown in Figure S2. All diffraction peaks for single crystals grown at the low temperature (750 °C) profile can be indexed with space group P6$_3$/mmc (No. 194) of hexagonal 2H-type,\textsuperscript{12} whereas the crystals grown at high temperature (950 °C) can be indexed with space group $P2_1/m$ (No. 11) of monoclinic 1T'-type at room temperature.\textsuperscript{15} In addition, the low temperature ($T = 200$ K) diffraction peaks are indexed with space group $Pmm2_1$ (No. 31) of orthorhombic Td-MoTe$_2$. The refined lattice parameters are listed in Table 1 and found to be consistent with the values reported in the literature.\textsuperscript{2,13,21–25} The diffraction patterns for the as-grown single crystals with planes of preferred orientation of (00l) reflections only are shown in Figure 1.

2H-MoTe$_2$ crystallizes in a layer-type structure with Mo atoms surrounded by six Te atoms in a trigonal prismatic coordination similar to that of 2H-MoS$_2$.\textsuperscript{12} 2H-MoTe$_2$ has been characterized by a stacking sequence of XyXyX... (xy: Te atom; y,x: Mo atom) with all monatomic planes in 2D hexagonal close packing,\textsuperscript{16} as shown in Figure 2(a). Alternatively, the 2H-type is easier to describe using the face-sharing hexagonal close packing of trigonal prismatic MoTe$_6$ units in each layer, and these layers are held together by the weak van der Waals force that gives rise to the quasi-two-dimensional character, as shown in Figure 2(d). On the other hand, 1T'-MoTe$_2$ could be viewed as a distortion of the typical 1T-type with layers composed of edge-sharing MoTe$_6$...
octahedra, which is demonstrated by the monoclinic distortion of $\beta \approx 93.9^\circ$ from the original hexagonal lattice, as shown in Figure 2(b). The distortion from $1T$ to $1T'$ not only breaks the in-plane hexagonal symmetry with the added mirror symmetry, the Te monolayers above and below the Mo monolayer also rearrange in a zigzag deviation along the $z$-direction relative to that of the flat Mo monolayer. The low temperature orthorhombic $Td$-MoTe$_2$ phase differs from the $1T'$-MoTe$_2$ phase by a small deviation from the right angle for $\beta$, i.e., the monoclinic angle of $1T'$ has a slightly larger ($\beta = 93.55^\circ$) than that of $\beta = 90^\circ$ for $Td$ with orthorhombic symmetry, as shown in Figure 2(c). It should also be noted that the structure symmetry of the $1T'$-WTe$_2$ reported in the literature is identical to the $Td$-MoTe$_2$ phase instead of the $1T'$-MoTe$_2$.27

The electronic structure of $2H$-MoTe$_2$, $1T'$-MoTe$_2$, and $Td$-MoTe$_2$ were calculated using a $30 \times 30 \times 6$ and $8 \times 16 \times 4$ Table 1. Crystal Structures and Lattice Parameters of MoTe$_2$ Single Crystals

<table>
<thead>
<tr>
<th>crystal type</th>
<th>space group</th>
<th>$a$ (Å)</th>
<th>$b$ (Å)</th>
<th>$c$ (Å)</th>
<th>angle</th>
</tr>
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<tr>
<td>$2H$-MoTe$_2$</td>
<td>$P6_3/mmc$</td>
<td>3.4474(2)</td>
<td>3.4474(2)</td>
<td>9.3528(3)</td>
<td>$\beta = 90^\circ; \gamma = 120^\circ$</td>
</tr>
<tr>
<td>$1T'$-MoTe$_2$</td>
<td>$P2_1/m$</td>
<td>3.4692(4)</td>
<td>6.3348(4)</td>
<td>13.8835(2)</td>
<td>$\beta \approx 93.9^\circ$</td>
</tr>
<tr>
<td>$Td$-MoTe$_2$</td>
<td>$Pmn2_1$</td>
<td>3.4783(4)</td>
<td>6.3563(2)</td>
<td>13.8935(2)</td>
<td>$\beta = 90^\circ$</td>
</tr>
</tbody>
</table>

Figure 2. Crystal structures of (a) $2H$-MoTe$_2$, (b) $1T'$-MoTe$_2$, and (c) $Td$-MoTe$_2$; the MoTe$_6$ unit has a trigonal prismatic coordination for $2H$ and a distorted octahedral coordination for $1T'$ and $Td$. MoO$_6$ coordination and projection views for $2H$-MoTe$_2$ along the [001] and $Td$-MoTe$_2$ along the [100] directions are shown in (d) and (e), respectively.

Figure 3. Band structures and density of states of (a) $2H$-MoTe$_2$, (b) $1T'$-MoTe$_2$, and (c) $Td$-MoTe$_2$. Blue and green lines indicate the Mo $d$-orbital and Te $p$-orbital, respectively. The corresponding first Brillouin zones (BZ) of all samples are shown below. The insets (b1) and (b2) of $1T'$-MoTe$_2$ are enhanced images around $E_F$, which show a well-defined continuous energy gap throughout the whole BZ. The insets (b1) and (b2) of $Td$-MoTe$_2$ show the enhanced band structure around $E_F$, which shows a pair of Weyl points along $\Gamma - X$ and is fully gapped along the $\Gamma - S$ direction.
Monkhorst–Pack k-mesh over the Brillouin zone (BZ), respectively, as shown in Figure 3. The first-principles calculations were based on the generalized gradient approximation (GGA)\textsuperscript{28} using the full-potential projected augmented wave method\textsuperscript{29} as implemented in the VASP package\textsuperscript{30}. The spin–orbit coupling (SOC) was included self-consistently. An indirect band gap of $\sim 0.81$ eV is obtained for 2H-MoTe\textsubscript{2} between the valence band maximum at $\Gamma$ point and conduction band minimum at the center of $\Gamma$-$K$, which is consistent with the transport measurement results. All of the bands are found to be spin degenerate due to the spatial inversion symmetry. The Mo-4d states dominate the energy band from $E_g$ to approximately $-1.0$ eV, whereas the Te-5p states are of lower energy. It is noteworthy that the valence band energy splitting at K point is primarily due to the SOC instead of the interlayer interaction (Figure 3(a)), which is similar to that found in 2H-MoS\textsubscript{2} and 2H-WS\textsubscript{2}\textsuperscript{31}.

The electronic structure of 1T-MoTe\textsubscript{2} (Figure 3(b)) is dramatically different from that of 2H-MoTe\textsubscript{2}. The finite density of states (DOS) at $E_g$ indicates a metallic phase, and this result is consistent with a metallic behavior of resistivity as shown in Figure 8. The orbital characters around $E_g$ are mainly derived from Mo-4d and Te-5p orbitals. Mo electron-like and Te hole-like band edges cross each other near $E_g$, which results in a band inversion and being gapped due to SOC. Hence, even though 1T-MoTe\textsubscript{2} has a metallic ground state, the valence band remains separated from the conduction band by a continuous energy gap throughout the entire BZ in the presence of SOC (below, Figure 3(b)). The band inversion feature and well-defined energy gap implies that 1T-MoTe\textsubscript{2} might be characterized as a $Z_2$ topological insulator (below, Figure 3(b))\textsuperscript{15}.

The third structure of MoTe\textsubscript{2} is orthorhombic Td phase. Contrary to monoclinic 1T-MoTe\textsubscript{2} phase that possesses spatial inversion symmetry, the crystal inversion symmetry has been broken in Td-MoTe\textsubscript{2}, which results from slight lattice structure distortion. Figure 3(c) shows the band structure of Td-MoTe\textsubscript{2} with SOC. The electronic structure exhibits significant band splitting due to the breaking of spatial inversion symmetry except for the time-reversal invariant points. Similar to the 1T phase, Td-MoTe\textsubscript{2} shows a metallic ground state, whereas the orbital characters around $E_g$ are mainly contributed from Mo-4d and Te-5p orbitals. Zooming in around $E_g$ along $\Gamma$-$X$ (below, Figure 3(c)) shows that two singly generated bands cross each other and form two Weyl nodes with opposite chirality, which is consistent with a recent theoretical prediction\textsuperscript{18}.

The scanning transmission electron microscopy (STEM) images of 2H-MoTe\textsubscript{2} and 1T-MoTe\textsubscript{2} were obtained in 200 kV JEOL 2100F equipped with a probe corrector for the spherical aberration. The TEM specimens were prepared by mechanical exfoliation and ion beam thinning (Gatan Precision Ion Polishing System II, PIPS II). Figure 4(a) is the high-angle annular dark-field (HAADF) STEM image of 2H-MoTe\textsubscript{2} along the [0001] zone axis. The intensity of the HAADF-STEM image is sensitively proportional to the atomic number $Z$, so that the positions of Mo atoms appear brighter in the figure. From the selected-area diffraction pattern (inset of Figure 4(a)), the (1100) lattice planes are labeled, and the plane spacing is found to be near 3.0 Å, which is consistent with the 2H-MoTe\textsubscript{2} hexagonal lattice structure of lattice constant $a = 3.52$ Å. The HAADF-STEM image of the 1T-MoTe\textsubscript{2} phase is shown in Figure 4(b) along the [001] zone axis. The (100) and (010) lattice plane spacings, as labeled in the figure, are consistent with the lattice constants of $a = 6.33$ Å and $b = 3.48$ Å for the 1T-MoTe\textsubscript{2} phase.

The scanning tunneling microscopy (STM) measurements were performed for both 2H-MoTe\textsubscript{2} and Td-MoTe\textsubscript{2} at 4.5 K and 1T-MoTe\textsubscript{2} at room temperature using an Omicron LT-STM and an electrochemically etched tungsten tip. $dI/dV$ (V) curves were acquired using the lock-in technique with a bias modulation of 10 mV. Crystals for STM measurements were prepared by in vacuo cleavage at room temperature in a preparation chamber at a base pressure lower than $1 \times 10^{-10}$ mbar before transfer to an STM chamber with a base pressure lower than $5 \times 10^{-11}$ mbar. Figure 5 shows atomically resolved topography images taken on (a, b) 2H-MoTe\textsubscript{2} and (c, d) Td-MoTe\textsubscript{2} crystals and the relationship between the observed atomic corrugations of Te surface layer and the corresponding Te surface lattice modeling. The topography for the 2H structure exhibits a clear hexagonal surface lattice, whereas the Td and 1T structures have rectangular surface lattices that are indistinguishable using STM topography imaging, which can only probe the uppermost atomic layer and cannot discern the difference in angle $\beta$ between the two structures. Tunneling spectroscopy curves acquired at 4.5 K as shown in Figure 5(e) suggest a metallic behavior for the Td-MoTe\textsubscript{2} surface, which is consistent with the band structure calculation of apparent n-type, on the other hand, the 2H-MoTe\textsubscript{2} surface shows a band gap of around 1.13 eV. These results are consistent with the calculation shown above and with previous predictions\textsuperscript{7}. Interestingly, a
comparison of enhanced STM topography of Td-MoTe₂ taken at 4.5 and 77 K, as shown in Figure 5(h), reveals clear protrusions seen at the unit cell edges at 4.5 K (marked by black arrows), which are absent at 77 K. This difference is consistent with a similar one recently observed in WTe₂, which has been attributed to the onset at low temperature of Umklapp interference of quasiparticles induced by Rashba-type splitting. Therefore, the topography at 77 K is expected to better represent the true surface lattice.

Step edges were observed in the as-cleaved MoTe₂ surfaces, as shown in Figure 6, which is helpful to probe the interlayer stacking by exploring the relative orientations of two adjacent MoTe₂ trilayers in stack, as illustrated in Figure 6(b and d). It has been nearly impossible to distinguish the reversed orientation of MoTe₂ units in either trigonal prismatic or octahedral coordination via the surface Te alone, mostly because the Te monolayer on the surface is always arranged in hexagonal close packing. However, we might distinguish the reversed orientation by taking advantage of the different local geometries for the Te vacancy defect sitting on surfaces of reversed orientation. Two types of characteristic defects can be identified in 2H-MoTe₂ and Td-MoTe₂ at 4.5 K (Figure 6), as indicated by the dark arrowhead-like defects of symmetry m for Td-MoTe₂ and the bright three-lobed defects of symmetry 3m for 2H-MoTe₂. In particular, the orientations for the two types of defects are seen reversed from one region to the next region that is one trilayer below, which is consistent with the 2H (P6₃/mmc) and Td (Pmn₂₁) symmetries of required reverse orientation between neighboring layers. Although 2H-MoTe₂ is able to maintain its hexagonal symmetry per layer, Td-MoTe₂ has lost its hexagonal symmetry with the additional in-plane mirror symmetry breaking (Figure 2), which is consistently reflected on the observed defect local symmetry of reversed orientations between neighboring layers exposed to the surface.
i.e., the former has the three-lobe shape of $3m$ symmetry and the latter has the arrowhead-like shape of $m$ symmetry only.

The $2H$ and $1T'$ phases of MoTe$_2$ of various thicknesses have been characterized with Raman spectroscopy at room temperature, as shown in Figure 7. The $2H$-MoTe$_2$ crystal shows Raman peaks between 100 and 300 cm$^{-1}$ (Figure 7(a)), including the prominent peak of the in-plane $E_{2g}$ mode at $\sim 234$ cm$^{-1}$, the out-of-plane $A_{1g}$ mode at $\sim 173$ cm$^{-1}$, and the out-of-plane bulk inactive $B_{2g}$ phonon mode ($\sim 290$ cm$^{-1}$).

The intensity of the Raman modes shows significant thickness dependence from the bulk sample to the nanothickness near $\sim 10$ layers. The presence of $B_{2g}$ mode in the 10-layers sample indicates that the synthesized sample is atomically thin. Besides the observed first-order Raman peaks, several second-order resonant peaks with relatively low intensities at $\sim 144$, 183, and 204 cm$^{-1}$ are also observed in $2H$-MoTe$_2$. The observed characteristic peaks of $1T'$-MoTe$_2$ are $A_g$ (108.2 cm$^{-1}$), $A_g$ (129.2 cm$^{-1}$), $A_g$ (164.8 cm$^{-1}$), and $A_g$ (261.5 cm$^{-1}$), as shown in Figure 7(b). Furthermore, a peak related to the Raman mode ($B_g$) at 180 cm$^{-1}$ was observed at room temperature, whereas this mode disappeared at low temperature, as shown in Figure S3, which could be related to the expected $1T'$ to $Td$ phase transition near $\sim 240$ K. The observed phonon modes match with those reported in the literature and the phase purities for both phases are confirmed in addition to the XRD structural analysis (Figure 1). The most prominent peak for $2H$-MoTe$_2$ is 234 cm$^{-1}$ corresponding to the $E_g$ phonon, in contrast to 166 cm$^{-1}$ corresponding to the $A_g$ phonon for $1T'$-MoTe$_2$, which is consistent with the expected easier excitation for Te atoms in the distorted trilayer along the $c$-direction for the latter.35

The electrical resistivities $\rho(T)$ as a function of temperature for both $1T'$- and $2H$-MoTe$_2$ are shown in Figure 8. Figure 8(a) shows the resistivity of $2H$-MoTe$_2$ as a function of temperature; a slight decrease upon cooling from 300 to 160 K is observed and then increases at lower temperature, which suggests the existence of a thermally excitable narrow gap of
35.3 meV fitted with an Arrhenius law below 50 K (inset of Figure 8(a)). Further, we also fit the data to the variable range hopping mechanism \( \rho = \rho_0 \exp\left(\frac{T_0}{T}\right)^{1/3} \) for two dimensions and found the hopping parameter \( (T_0/3)^{1/3} \) to be equal to 80 K.1/3 This gives rise to a differential activation energy of 31 meV at 50 K. Activation energy obtained from both mechanisms is on the same order. Figure 8(b) shows the Hall resistivity \( \rho_{xy} \) at several temperatures for the 2H-MoTe2 single crystal. \( \rho_{xy} \) data shows a good linearity relative to the magnetic field up to 9 T above 50 K, indicating the dominance of electron carriers. Electronic mobility decreases with rising temperature (inset of Figure 8(b)) of \( \sim 400 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1} \) at 50 K. Because the band gap size of the ideal 2H-MoTe2 has been estimated to be near \( \sim 1.0 \text{ eV} \) (Figure 3), it is possible that the identified narrow gap from the transport study is the gap between the impurity and the conduction bands of n-type doping introduced by Te vacancies.21 The magnetic field of 9 T does not affect the resistivity behavior with apparently no magnetoresistance.

The \( \rho(T) \) curve of 1T’-MoTe2 indicates a metallic behavior of \( \sim 4.4 \times 10^{-4} \text{ Ω cm} \) at room temperature as shown in Figure 8(c). A hysteretic anomaly in resistivity is observed between 240 and 265 K upon the warming and cooling cycles, which has been assigned corresponding to a first-order structural phase transition from the 1T’ phase to a low-temperature orthorhombic Td phase.14 This 1T’-to-Td phase transition can also be identified in the magnetic susceptibility and specific heat measurements, as shown in Figures S4 and S5. The magnetic susceptibilities of single crystal 1T’-MoTe2 as a function of temperature were measured in an applied magnetic field of 10 kOe for field applied along the ab-direction, and a hysteretic \( \chi(T) \) step anomaly near \( \sim 254 \text{ K} \) is found. In addition, the heat capacity \( C_p(T) \) shows an anomaly of onset near \( \sim 250 \text{ K} \). Fitting \( C_p \) with the electronic and phonon contributions in \( C_p/T = \gamma + \beta T^2 \), where \( \gamma \) is the normal-state electronic contribution and \( \beta \) is the lattice contribution to the specific heat, the fitted results yield \( \gamma = 3.22 \text{ mJ mol}^{-1} \text{ K}^{-2} \) and \( \beta = 0.849 \text{ mJ mol}^{-1} \text{ K}^{-4} \) for Td-MoTe2. Considering the saturation value \( (74.93 \text{ J mol}^{-1} \text{ K}^{-1} \text{ at 300 K}) \) as the Dulong–Petit classical limit, the derived Debye temperature \( (\theta = N(12/5)\pi^2 R^0 \Omega^{-3} (R = 8.314 \text{ J mol}^{-1} \text{ K}^{-1}) \) is found to be \( 190 \text{ K} \), which is of the same order and consistent with that found in orthorhombic WTe2.41

The \( \rho(T) \) behavior does not show any noticeable change under a 9 T magnetic field for both 1T’-MoTe2 and Td-MoTe2 (Figure 8(c)), which is surprising compared to WTe2 (orthorhombic Td structure) being identified as a large magnetoresistance material due to perfect electron and hole compensation.5 Such a compensation of electron and holes are absent in Td-MoTe2.20 The low temperature part \( (T < 50 \text{ K}) \) of the \( \rho(T) \) data is well fitted by \( \rho(T) = \rho_0 + AT^2 \) following the Fermi liquid model, as shown by the solid red line in the inset of Figure 8(c). The obtained fitting parameters are the residual resistivity \( \rho_0 = 5.97 \times 10^{-5} \text{ Ω cm} \) and coefficient \( A = 1.42 \times 10^{-8} \text{ Ω cm} \) K\(^{-2}\), which indicates the Fermi liquid-like behavior for Td-MoTe2 single crystal below 50 K.32 In semimetals like 1T’MoTe2, the temperature-dependent resistivity can be expressed as \( \rho = a + bT + cT^2 \), where the first term is from electron-defect scattering, the second term is electron-electron scattering, and the third term is electron-phonon scattering. At low temperature, because of negligible phonon contribution, the resistivity varies almost as \( T^2 \), commonly known as the Fermi liquid behavior. At higher temperature, although the \( T^2 \) term is present, because of large phonon contribution (large value of coefficient \( b \)), the resistivity varies almost linearly. The \( \rho_0 \) as a function of magnetic field at different temperatures for 1T’-MoTe2 (at \( T = 300 \text{ K} \)) and Td-MoTe2 (below 260 K) is shown in Figure 8(d). Below 260 K, \( \rho_0 \) shows a nonlinear behavior with magnetic field for the Td-MoTe2 single crystal. Electrons remain the majority charge carriers in the whole temperature range for both 1T’-MoTe2 and Td-MoTe2 single crystals. Td-MoTe2 has an electron mobility of 133 cm\(^2\) V\(^{-1}\) s\(^{-1}\) at 2 K and decreases with increasing temperature, whereas electron mobility of 1T’-MoTe2 is very small at the level of \( \sim 3 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1} \) at 300 K (inset of Figure 8(d)). The decrease in the carrier mobility of both 2H and 1T’ forms of MoTe2 decreases with temperature. This is mostly due to the increased phonon density at higher temperatures, which results in the smaller \( T \) (average interval between two scattering events). As mobility is directly related to \( T (\mu = qT/m^* \), where \( q \) is the electronic charge and \( m^* \) is the effective mass), mobility also decreases with temperature. The effect of the phonon can easily be seen in the resistivity in metals. In 1T’-MoTe2, due to the negligible electron–phonon scattering contribution at low temperature, the rate of decrease in resistivity decreases and almost saturates. Interestingly, the mobility also saturates at low temperature and decreases thereafter upon increasing temperature. These transport properties are consistent with those estimated from tunneling spectroscopy of STM shown in Figure 5.

### CONCLUSIONS

Large and high-quality single crystals of layered MoTe2 in both 2H- and 1T’-types have been grown using a CVT method, and the low temperature Td phase can also be generated from the as-grown 1T’-MoTe2. The three phases of distinctly different layer structures have been confirmed rigorously with XRD, Raman spectroscopy, STEM, and STM studies. The physical properties have also been compared through electrical resistivity, magnetic susceptibility, and specific heat measurements. Scanning tunneling spectroscopy confirms the reversed stacking relationship between the adjacent MoTe2 trilayers via the low symmetry of the point defects. The metallic behavior of Td-MoTe2 is found to be consistent with the predicted Weyl semimetal phase, and the n-type semiconducting behavior for 2H-MoTe2 has also been confirmed via resistivity and STM tunneling spectroscopy.

### ASSOCIATED CONTENT

3 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.chemmater.6b04363.

Crystal growth profiles, refined room-temperature SXRD patterns, Raman spectra for 1T’-MoTe2 single crystals with various temperatures, magnetic susceptibilities, and heat capacity data (PDF)

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The authors declare no competing financial interest.

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