THE HIGGS BOSON & BEYOND

Tao Han PITT PACC, Univ. of Pittsburgh TsingHua U. / CFHEP, Beijing Joint Colloquium National Taiwan University, Dec. 8, 2015





2013 Nobel Laureate

© The Nobel Foundation. Photo: Lovisa Engblom.

François Englert and Peter W. Higgs "for the <u>theoretical discovery</u> of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

THE DISCOVERY: July 4th, 2012: **A NEUTRAL BOSON DECAY TO TWO PHOTONS**



Summer 2015 update: 50 for both fermion coupling $h \rightarrow \tau\tau$ & bosonic coupling $WW \rightarrow h$



- it's neutral, a boson
- it's spin-0, parity-even
- it couples to mass, non-universally All indications point to the SM Higgs !

A milestone discovery: It is a brand new class!





Mosaic of the CMS and ATLAS detectors (as in 2007), part of the Large Hadron Collider at CERN. In 2012, research teams used these detectors to fingerprint decay products from the long-sought Higgs boson and determine its mass successfully testing a key prediction of the standard model of particle physics. Photo: Maximilian Brice and Claudia Marcelland/CERN



50 years theoretical work ... 25 years experimental work ...

Congratulations to our CMS colleagues in Taiwan !

THE NATURE OF FORCES:

The Four Fundamental Interactions

long range $\sim (G_N m_1 m_2)/r^2$

Gravity Strong Véak Néak short range ~ e^{-mr}/r^2 All forces in the world can be attributed to these four interactions!

All known physics

$$W = \int_{k < \Lambda} [\mathcal{D}g \dots] \exp\left\{\frac{i}{\hbar} \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G}R - \frac{1}{4}F^2 + \bar{\psi}i \not\!\!D\psi - \lambda\phi\bar{\psi}\psi + |D\phi|^2 - V(\phi)\right]\right\}$$

amplitude current understanding

quantum mechanics

spacetime

gravity strong & electroweak

matter

Higgs

E&M: Most Successful in Theory & Practice!

$$\mathcal{L} = \left(-\frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right) + \bar{\psi}(i\gamma^{\mu}D_{\mu} - m_{e})\psi$$
$$F^{\mu\nu} = \partial^{\mu}A^{\nu} - \partial^{\nu}A^{\mu}, \quad D_{\mu} = \partial_{\mu} + ieA$$

- At low energies → Maxwell's theory; vector-like coupling by a U_{em}(1) gauge symmetry
- At high energies → Quantum-mechanical, renormalizable, most accurate (in science!): a part of trillion

 $a_e^{theo} = 0.001159652181643(763)$ $a_e^{exp} = 0.00115965218073(28)$

• QED becomes strongly interacting asymptotically (screening effects) $\alpha(Q_0^2) = \alpha(Q_0^2)$





At ultra-violet (UV) \rightarrow theory is invalid.

The strong force: SU_c(3) Quantum Chromo-Dynamics Successful Theory, Challenging in Practice!

$$\mathcal{L} = \left(-\frac{1}{4}F^{a}_{\mu\nu}F^{a\mu\nu}\right) + \sum_{f} \bar{q}_{f}(i\gamma^{\mu}\partial_{\mu} - g_{s}\gamma^{\mu}A_{\mu} - m_{f})q_{f}$$

 $O\mu A_{\nu} \neq i g_s |A_{\mu}, A_{\nu}|$

$$A^{\mu}(x) = \sum_{1}^{6} A(x)^{\mu}_{a} T^{a}, \quad [T^{a}, T^{b}] = if_{abc}T^{b}$$

 At short distances/high energies → asymptotically free (anti-screening effects)

$$\alpha_s(Q^2) = \frac{12\pi}{(33 - 2n_f)\ln(Q^2/\Lambda^2)}$$

 $\Gamma_{\mu\nu} - 0\mu A_{\nu}$

Highly predictable at high energies: Crucial for HEP, early Universe ...

 At long distances/low energies > 10⁻¹³cm
 → Strongly interacting: quarks condensate (π⁰, π[±]...) & (colorless) hadrons (p⁺, n) formed.
 Short range force by a dynamical mass: e<sup>-m_π r/r²
</sup>





The Weak force: Quark & Lepton Flavor Transitions Beta decay $n \rightarrow p^+ e^- v \rightarrow$ Charged current interaction: W[±] Inspired by EM current-current interactions, Fermi proposed (1934)



force range $\sim \sqrt{G_F} \sim M_W^{-1} \sim 10^{-18} \mathrm{m}$



Weak interaction based on $SU(2) \ge U(1)$: (Glashow, '63)

$$B_{\mu\nu} = \partial_{\mu}B_{\nu} - \partial_{\nu}B_{\mu}$$

$$W^{i}_{\mu\nu} = \partial_{\mu}W^{i}_{\nu} - \partial_{\nu}W^{i}_{\mu} - g\epsilon_{ijk}W^{j}_{\mu}W^{k}_{\nu}$$

$$- \frac{g}{2\sqrt{2}}\sum_{i}\overline{\Psi}_{i}\gamma^{\mu}(1-\gamma^{5})(T^{+}W^{+}_{\mu}+T^{-}W^{-}_{\mu})\Psi_{i}$$

$$- e\sum_{i}q_{i}\overline{\psi}_{i}\gamma^{\mu}\psi_{i}A_{\mu}$$

$$- \frac{g}{2\cos\theta_{W}}\sum_{i}\overline{\psi}_{i}\sqrt{\mu}(g^{i}_{V}-g^{i}_{A}\gamma^{5})\psi_{i}Z_{\mu}.$$

The local gauge symmetry prevents gauge bosons masses! $\frac{1}{2}M_A^2 A_\mu A^\mu \rightarrow \frac{1}{2}M_A^2 (A_\mu - \frac{1}{e}\partial_\mu \alpha)(A^\mu - \frac{1}{e}\partial^\mu \alpha) \neq \frac{1}{2}M_A^2 A_\mu A^\mu$ Pauli's rejection to the Yang-Mills theory. **The Weak force: Quark & Lepton Flavor Transitions** Even worse:

"The Left- and right-chiral electrons carry different Weak charges" (Lee & Yang)

 Simple structure and particle contents: Leptons:

 $\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$, $\begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L$, $\begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$, $e_R, \mu_R, \tau_R, (\nu'_R s ?)$

Quarks:

 $\begin{pmatrix} u \\ d \end{pmatrix}_L$, $\begin{pmatrix} c \\ s \end{pmatrix}_L$, $\begin{pmatrix} t \\ b \end{pmatrix}_L$, u_R , d_R , c_R , s_R , t_R , b_R

Fermion masses also forbidden by gauge symmetry! $-m_e\bar{e}e = -m_e\bar{e}\left(\frac{1}{2}(1-\gamma_5) + \frac{1}{2}(1+\gamma_5)\right)e = -m_e(\bar{e}_Re_L + \bar{e}_Le_R)$

Electroweak gauge theory \rightarrow massless!





THE SPONTANEOUS SYMMETRY BREAKING

"The Lagrangian of the system may display an symmetry, but the ground state does not respect the same symmetry."

Known Example: Ferromagnetism

Above a critical temperature, the system is symmetric, magnetic dipoles randomly oriented. Below a critical temperature, the ground state is a completely ordered configuration in which all dipoles are ordered in some arbitrary direction $SO(3) \rightarrow SO(2)$



Domains Before Magnetization



Domains After Magnetization

Low temperature super-conductivity is another example! The concept of SSB: profound, common.

THE NAMBU-GOLDSTONE THEOREM

"If a continuous symmetry of the system is spontaneously broken, then there will appear a massless degree of freedom, called the Nambu-Goldstone boson."



Except the photon, no massless boson (a long-range force carrier) has been seen in Nature! (Recall Pauli's criticism) The Spontaneous Symmetry Breaking: Brilliant idea & common phenomena, confronts the Nambu-Goldstone theorem! -- A show stopper ?

THE MAGIC IN 1964

"If a LOCAL gauge symmetry is spontaneously broken, then the gauge boson acquires a mass by absorbing the Goldstone mode."

BROKEN SYMMETRY AND THE MASS OF GAUGE VECTOR MESONS*

F. Englert and R. Brout Faculté des Sciences, Université Libre de Bruxelles, Bruxelles, Belgium (Received 26 June 1964)

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS PLB

P.W.HIGGS Tait Institute of Mathematical Physics, University of Edinburgh, Scotland

Received 27 July 1964

BROKEN SYMMETRIES AND THE MASSES OF GAUGE BOSONS

Peter W. Higgs

PRL

Tait Institute of Mathematical Physics, University of Edinburgh, Edinburgh, Scotland (Received 31 August 1964)

GLOBAL CONSERVATION LAWS AND MASSLESS PARTICLES*

G. S. Guralnik,[†] C. R. Hagen,[‡] and T. W. B. Kibble PRL Department of Physics, Imperial College, London, England (Received 12 October 1964)

An illustrative (original) Model: $\mathscr{L} = |\mathscr{D}^{\mu}\phi|^{2} - \mu^{2}|\phi|^{2} - |\lambda|(\phi^{*}\phi)^{2} - \frac{1}{4}F_{\mu\nu}F^{\mu\nu},$

where

$$\phi = \frac{\phi_1 + i\phi_2}{\sqrt{2}}$$

is a complex scalar field⁴ and as usual

$$\mathscr{D}_{\mu} \equiv \partial_{\mu} + i q A_{\mu}$$

and

$$F_{\mu\nu}\equiv\partial_{\nu}A_{\mu}-\partial_{\mu}A_{\nu}.$$

The Lagrangian (5.3.1) is invariant under U(1) rotations

$$\phi \to \phi' = e^{i\theta}\phi$$

and under the local gauge transformations

$$\phi(x) \rightarrow \phi'(x) = e^{iq\alpha(x)}\phi(x),$$

 $A_{\mu}(x) \rightarrow A'_{\mu}(x) = A_{\mu}(x) - \partial_{\mu}\alpha(x).$

[¶]C. Quigg, Gauge Theories of the Strong ...

An illustrative (original) Model: After the EWSB, parameterized in terms of $\langle \phi \rangle_0 = v/\sqrt{2}, \qquad \phi = e^{i\zeta/v}(v+\eta)/\sqrt{2}$ $\approx (v+\eta+i\zeta)/\sqrt{2}.$

Then the Lagrangian appropriate for the study of small oscillations is

$$\mathscr{L}_{so} = \frac{1}{2} [(\partial_{\mu} \eta)(\partial^{\mu} \eta) + 2\mu^{2} \eta^{2}] + \frac{1}{2} [(\partial_{\mu} \zeta)(\partial^{\mu} \zeta)]$$
$$-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + q\nu A_{\mu} (\partial^{\mu} \zeta) + \frac{q^{2} \nu^{2}}{2} A_{\mu} A^{\mu} + \cdots$$

The gauge field acquires a mass, mixes with the Goldstone boson. Upon diagonalization: $\frac{q^2v^2}{2}\left(A_{\mu}+\frac{1}{qv}\partial_{\mu}\zeta\right)\left(A^{\mu}+\frac{1}{qv}\partial^{\mu}\zeta\right),$

a form that pleads for the gauge transformation

$$A_{\mu} \rightarrow A'_{\mu} = A_{\mu} + \frac{1}{qv} \partial^{\mu} \zeta,$$

which corresponds to the phase rotation on the scalar field

$$\phi \rightarrow \phi' = e^{-i\zeta(x)/\nu}\phi(x) = (\nu+\eta)/\sqrt{2}.$$

the resultant Lagrangian is then: $\mathscr{L}_{so} = \frac{1}{2} [(\partial_{\mu} \eta)(\partial^{\mu} \eta) + 2\mu^{2} \eta^{2}] - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{q^{2} \nu^{2}}{2} A'_{\mu} A'^{\mu}$

• an η -field, with (mass)² = $-2\mu^2 > 0$; the Higgs boson!

• a massive vector field A'_{μ} , with mass = qv

• no ζ -field.

• By virtue of a gauge choice - the unitary gauge, the ζ -field disappears in the spectrum: a massless photon "swallowed" the massless NG boson! Degrees of freedom count: **Before EWSB:** After: 2 (scalar)+2 (gauge pol.); 1 (scalar)+3 (gauge pol.) • Two problems provide cure for each other! massless gauge boson + massless NG boson → massive gauge boson + no NG boson This is truly remarkable!

A FEW OBSERVATIONS A. The Higgs mechanism ≠ a Higgs boson ! From theoretical point of view, **3 Nambu-Goldstone bosons were all we need!** A non-linear realization of the gauge symmetry: $U = \exp\{i\omega^{i}t^{i}/v\}, \quad D_{\mu}U = \partial_{\mu}U + igW_{\mu}^{i}\frac{\tau^{i}}{2}U - ig'UB_{\mu}\frac{\tau^{3}}{2}$ $\mathcal{L} = \frac{v^2}{2} [D^{\mu} U^{\dagger} D_{\mu} U] \to \frac{v^2}{4} (\sum g^2 W_i^2 + g'^2 B^2)$ The theory is valid to a unitarity bound ~ 2 TeV The existence of a light, weakly coupled Higgs boson carries important message for our understanding & theoretical formulation in & beyond the SM -UV completion / renormalizibility.

B. λ: a "New Force" The Higgs potential: $V = -\mu^2 / \phi/^2 + \lambda \phi/^4$ It represents a weakly coupled new force (a fifth force):

- In the SM, λ is a free parameter, now measured: $\lambda = m_{\rm H}^2 / 2v^2 \approx 0.13$
 - Is it fundamental or induced?
 - In SUSY, it is related to the gauge couplings tree-level: $\lambda = (g_L^2 + g_Y^2)/8 \approx 0.3/4 \leftarrow a$ bit too small
 - In composite/strong dynamics, harder to make λ big enough.
 (due to the loop suppression by design)
 Already possess challenge to BSM theories.

At higher energies, λ is NOT asymptotically free. It blows up at a high-energy scale (the Landau pole), unless it starts from small (or zero \rightarrow triviality). For $m_{H} = 126$ GeV, rather light:



The SM can be a consistent 500 perturbative theory up to M_{pl} ! allowing M_N , M_{GUT} , ... 400 M_H [GeV/*c*²] Triviality FW 300 Precision 200 Top-Yukawa drags the vacuum 126 meta-stable, 100 *r*acuum is absolute minimum New physics below 10⁷⁻¹¹ GeV? 0 5 19 7 15 17 3 13 $\log_{10} \Lambda [GeV]$ The new coupling λ very important!

600

C. Electroweak Super-Conductivity

Normal phase \Rightarrow $E^2 = p^2 c^2$

Long-range force



⇐Superconducting phase $E^2 = p^2 c^2 + m^2 c^4$ gap leads to $\sim exp(-r/\lambda)$ $\lambda \sim m^{-1}$ penetration depth

In "conventional" electro-magnetic superconductivity: $m_{\gamma} \sim m_e/1000$, $T_c^{em} \sim \mathcal{O}(\text{few } K)$. BCS theory.

In "electro-weak superconductivity":

В

T>Tc

 $m_w \sim G_F^{-\frac{1}{2}} \sim 100 \text{ GeV}, \ T_c^w \sim 10^{15} K!$

The Higgs potential is of the Landau-Ginsburgh form, but it represents a new fundamental interaction.

"... most of the grand underlying principles have been firmly established. An eminent physicist remarked that the future truths of physical science are to be looked for in the sixth place of decimals." ---- Albert Michelson (1894)

Michelson–Morley experiments (1887): "the moving-off point for the theoretical aspects of the second scientific revolution"

Will History repeat itself (soon)?

Nima Arkani-Hamed (Director of CFHEP, Beijing)



Danny franketer (and frankete

NEW ERA: Under the Higgs lamp post

The "Observation" papers: Now 3600 cites each!



Vast scope of topics, from interpretations, explorations in & beyond the SM; applications in astronomy, cosmology, CC; strings/branes, to "Philosophical Perspectives"

Question 1: The Nature of EWSB? In the SM: V.(#) $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \mu^2 H^2 + (\lambda v H^3) + \frac{\lambda}{4} H^4$ Re(d) Fully determined at the weak scale: $v = (\sqrt{2}G_F)^{-1/2} \approx 246 \text{ GeV} \quad m_H \approx 126 \text{ GeV}$ $m_H^2 = 2\mu^2 = 2\lambda v^2 \quad \Rightarrow \quad \mu \approx 89 \text{ GeV}, \quad \lambda \approx \frac{1}{2}.$ It is a weakly coupled new force, underwent a 2nd order phase transition. Is there anything else?

Question 1: The Nature of EWSB ?



With new physics near the EW scale: $V(h) \rightarrow m_h^2(h^{\dagger}h) + \frac{1}{2}\lambda(h^{\dagger}h)^2 + \frac{1}{3!\Lambda^2}(h^{\dagger}h)^3 \rightarrow \lambda_{hhh} = (7/3)\lambda_{hhh}^{SM}$ $\rightarrow \frac{1}{2}\lambda(h^{\dagger}h)^2 \log\left[\frac{(h^{\dagger}h)}{m^2}\right] \qquad \Rightarrow \lambda_{hhh} = (5/3)\lambda_{hhh}^{SM}$

 $\frac{\lambda(h^{+}h)^{2} \text{ term could be made "-":}}{\text{leading to EW phase transition strong 1st order!}$ $\xrightarrow{\rightarrow} O(1) \text{ deviation on } \lambda_{\text{hhh}}$

Question 2: The "Naturalness"

"... scalar particles are the only kind of free particles whose mass term does not break either an internal or a gauge symmetry." Ken Wilson, 1970



If $\Lambda^2 \gg m_H^2$, then unnaturally large cancellations must occur.

The Higgs mass fine-tune: $\frac{\delta m_H}{m_H} \sim 1\% (1 \text{ TeV}/\Lambda)^2$

Natural: O(1 TeV) new physics, associated with ttH. Unknown: Deep UV-IR correlations? Agnostic: Multiverse/anthropic?

"Naturalness" in perspective:



Unbelievable! 4 mm² / 20 cm² ~ 10⁻³ fine-tune.

"Naturalness" -> TeV scale new physics.



Question 4: The "Flavor Puzzle"

- Particle mass hierarchy
- Patterns of quark, neutrino mixings
- New CP-violation second second

Higgs Yukawa couplings as the pivot!



The Higgs as pivot for "seesaw": $m_{\nu} \sim \kappa \frac{\langle H^0 \rangle^2}{M}$

Type I seesaw: $M = M_{N_{r}}$ right-handed (sterile) N_{R}^{i} $H \rightarrow NN, N \rightarrow Hv, ...$

Type II seesaw: $M = M_{H++}$, a Higgs triplet Φ_3 $H^{++} \rightarrow l_i^+ l_j^+$

Type III seesaw: $M = M_T$, a fermionic triplet T_3 : $T^+ \rightarrow H l_i^+, T^0 \rightarrow W^{\pm} l$

Watch out: $H^0 \rightarrow \mu \tau (l_i^+ l_j^-)$ for BSM flavor physics!

COLLISION COURSE

Nature News, July '14

Particle physicists around the world are designing colliders that are much larger in size



Facility	HL-LHC	ILC	ILC(LumiUp)	CLIC	TLEP (4 IPs)	HE-LH	C VLHC
$\overline{s} \; ({ m GeV})$	$14,\!000$	250/500/1000	250/500/1000	350/1400/3000	240/350	33,000	100,000
$\mathcal{L}dt \; (\mathrm{fb}^{-1})$	3000/expt	250 + 500 + 1000	1150 + 1600 + 2500	500 + 1500 + 2000	10,000+2600	3000	3000
$dt \ (10^7 \mathrm{s})$	6	3+3+3	(ILC 3+3+3) + 3+3+3	3.1+4+3.3	5+5	6	6

Higgs-Factory: Mega (10⁶) Higgs Physics



ILC: $E_{cm} = 250 (500) \text{ GeV}, 250 (500) \text{ fb}^{-1}$

- Model-independent measurement: ILC Report: 1308.6176 $\Gamma_{\rm H} \sim 6\%$, $\Delta m_{\rm H} \sim 30 \text{ MeV}$ (HL-LHC: assume SM, $\Gamma_{\rm H} \sim 5-8\%$, $\Delta m_{\rm H} \sim 50 \text{ MeV}$)
- TLEP 10⁶ Higgs: $\Gamma_{\rm H} \sim 1\%$, $\Delta m_{\rm H} \sim 5$ MeV. TLEP Report: 1308.6176

THE NEXT ENERGY FRONTIER: 100 TEV HADRON COLLIDER



Snowmass QCD Working Group: 1310.5189 Arkani-Hamed, TH, Mangano, LT Wang, 1511.06495

Higgs Self-couplings:



Triple coupling sensitivity: Test the shape of the Higgs potential, and the fate of the EW-phase transition!

Snowmass 1310.8361

100 TeV pp

		HL-LHC	ILC500	ILC500-up	ILC1000	ILC1000-up	CLIC1400	CLIC3000	HE-LHC	VLHC
	$\sqrt{s} \; (\text{GeV})$	14000	500	500	500/1000	500/1000	1400	3000	33,000	100,000
λ (50%) 83% 46% (21%) 13% 21% 10% 20% 8%	$\int \mathcal{L}dt \; (\mathrm{fb}^{-1})$	3000/expt	500	1600^{\ddagger}	500 + 1000	$1600 + 2500^{\ddagger}$	1500	+2000	3000	3000
	λ	50%	83%	46%	21%	13%	21%	10%	20%	8%

Pushing the "Naturalness" limit



DM Searches



WIMP DM: $M_{\rm DM} < 1.8 \,\,{\rm TeV}\left(\frac{g_{\rm eff}^2}{0.3}\right)$



New Particle Searches

Electroweak Resonances: Z',W'

Colored Resonances:



Heavy Higgs bosons: H^0 , H^{\pm}



Mass reach at 100 TeV: ~ 5x over LHC

New (vector-like) leptons



A GRAND PICTURE:



Summary:

- The Higgs boson is a new class, at a pivotal point of energy, intensity intensity, cosmic frontiers. "Naturally speaking": It should not be a lonely solitary particle. Precision Higgs physics: LHC lights the way: g~10%; $\lambda_{\text{HHH}} \sim 50\%$; Br_{inv.}~ 20% Higgs factory/SppC: g~1%; λ_{HHH} < 10%; Br_{inv.} ~ 2%; Γ_{tot} < 6% - CEPC/SppC New physics reach: 6x LHC reach: 10 - 30 TeV \rightarrow fine-tune < 10^{-4} WIPM DM mass ~ 1 – 5 TeV An exciting journey ahead!

energy

Question 2: The "Naturalness"

"... scalar particles are the only kind of free particles whose mass term does not break either an internal or a gauge symmetry." Ken Wilson, 1970



If $\Lambda^2 \gg m_H^2$, then unnaturally large cancellations must occur. Cancelation in perspective: $m_H^2 = 36,127,890,984,789,307,394,520,932,878,928,933,023$ -36,127,890,984,789,307,394,520,932,878,928,917,398 $= (125 \text{ GeV})^2 ! ?$

Natural: O(1 TeV) new physics, associated with ttH. Unknown: Deep UV-IR correlations? Agnostic: Multiverse/anthropic?