Confinement/Deconfinement Phase Transitions in Strongly Coupled Anisotropic Theories

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Short talk: therefore bottom-up presentation

- Strongly coupled anisotropic theory.
 - ✓ Unequal diagonal energy momentum tensor elements for the spatial dimensions. Unequal pressures.
- How the theory looks like and how to obtain it?
 - ✓ 4d SU(N) gauge theory in the large N_c -limit.
 - $\checkmark\,$ Its dynamics are affected by a scalar operator $\mathcal{O}\sim \textit{TrF}^2.$
 - ✓ Anisotropy is introduced by another operator $\tilde{\mathcal{O}} \sim \theta(x_3) TrF \wedge F$ with a space dependent coupling.
 - ✓ On the gravity dual side a "backreacting" scalar field depending on spatial directions axion; and a non-trivial dilaton.
 - Eventually an Einstein-Axion-Dilaton theory in 5 dimensions with a non-trivial potential.

Theory Evolution:

Non-Confining:

(Azeyanagi, Li, Takayanagi, 2009; Mateos, Trancanelli, 2011; Jain, Kundu, Sen, Sinha, Trivedi, 2015;...)

Confining:

(D.G., Gursoy, Pedraza, 2017)

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- May have heard already about such theories since they violate well known universal relations:
 - ✓ Shear Viscosity over entropy density ratio parametrically lower than KSS bound: 1/4π. (Rebhan,Steineder 2011; Jain, Samanta, Trivedy 2015; D.G., Gursoy, Pedraza, 2017;...)
 - ✓ Langevin Coefficient inequality for heavy quark motion in the anisotropic plasma $\kappa_L > < \kappa_T$.

(Gursoy, Kiritsis, Mazzanti, Nitti 2010; D.G, Soltanpanahi, 2013a, 2013b) \checkmark ...

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Confinement/Deconfinement Phase transitions

 The free energy of the theories vs the temperature *T* for different anisotropy (α/j=0,1,3):



- Horizontal Axis: Confining Phase.
- Upper Branch: Black hole A:Deconfining Plasma Phase.
- Lower Branch: Black hole B:Deconfining Plasma Phase.
- a/j ≃ 2: A critical value above which a richer structure in the phase diagram exist.

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• The T_c is reduced in presence of anisotropies of the theory.

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So far: Findings and Proposal

- The *Tc*(α) decrease with α, resembling the phenomenon of inverse magnetic catalysis where the confinement-deconfinement temperature decreases with the magnetic field B.
- No charged fermionic degrees of freedom in our case; our plasma is neutral.
- Our findings suggest that the anisotropy by itself could instead be the cause of lower T_c in presence of anisotropies.

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η/s for our theory: Dependence on the Temperature.



• The ratio depends on the temperature at $\alpha/T \gg 1$ as

$$\frac{\eta_{13}}{s} \sim \left(\frac{T}{\tilde{\alpha}|1+3z-\theta|}\right)^{2-\frac{2}{z}}$$

• The range of the temperature power is $[0,\infty)$.

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The Anisotropic Theory

Consider a generalized (reduced) Einstein-Axion-Dilaton action with a potential for the dilaton and an arbitrary coupling between the axion and the dilaton:

$$S = \frac{1}{2\kappa^2} \int d^5 x \sqrt{-g} \left[R - \frac{1}{2} (\partial \phi)^2 + V(\phi) - \frac{1}{2} Z(\phi) (\partial \chi)^2 \right].$$

The eoms read

$$\begin{split} R_{\mu\nu} &- \frac{1}{2} R g_{\mu\nu} = \frac{1}{2} \partial_{\mu} \phi \partial_{\nu} \phi + \frac{1}{2} Z(\phi) \partial_{\mu} \chi \partial_{\nu} \chi - \frac{1}{4} g_{\mu\nu} (\partial \phi)^2 - \frac{1}{4} g_{\mu\nu} Z(\partial \chi)^2 + \frac{1}{2} g_{\mu\nu} V(\phi) , \\ \frac{1}{\sqrt{-g}} \partial_{\mu} \left(\sqrt{-g} g^{\mu\nu} \partial_{\nu} \phi \right) &= \frac{1}{2} \partial_{\phi} Z(\phi) (\partial \chi)^2 - V'(\phi) , \\ \frac{1}{\sqrt{-g}} \partial_{\mu} \left(\sqrt{-g} g^{\mu\nu} \partial_{\nu} \chi \right) &= 0 . \end{split}$$

Where the functions

$$V(\phi) = 12\cosh(\sigma\phi) - 6\sigma^2 \phi^2, \qquad Z(\phi) = e^{2\gamma\phi},$$

((Gubser, Nellore), Pufu, Rocha 2008a,b)Note: For $\sigma = 0, \gamma = 1$ the action and the solution of eoms, are of IIB supergravity. $(Mateos, Trancanelli, 2011) \rightarrow \infty$

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Let us apply the black hole background ansatz

$$ds^{2} = \frac{e^{-\frac{1}{2}\phi(u)}}{u^{2}} \left(-\mathcal{FB} dt^{2} + dx_{1}^{2} + dx_{2}^{2} + \mathcal{H} dx_{3}^{2} + \frac{du^{2}}{\mathcal{F}}\right),$$
$$\chi = \alpha x_{3}, \qquad \phi = \phi(u),$$

 $\phi(u), \mathcal{B}(u), \mathcal{F}(u), \mathcal{H}(u)$ four functions to be found, and α is the constant anisotropic parameter, u_h is the black hole horizon(related to the temperature of the theory).

Note:

The linear axion simplifies tremendously the system of equations!

Solve the system...

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Solutions: A demonstration

• Fixing (γ, σ) and α and u_h we get the metric flow from boundary to horizon:

$$ds^{2} = \frac{e^{-\frac{1}{2}\phi(u)}}{u^{2}}\left(-\mathcal{FB} dt^{2} + dx_{1}^{2} + dx_{2}^{2} + \mathcal{H} dx_{3}^{2} + \frac{du^{2}}{\mathcal{F}}\right),$$



• In sufficiently high temperatures, $T \gg \alpha$ for $\gamma = 1, \sigma = 0$ the Einstein equations can be solved analytically:

$$\mathcal{F}(u) = 1 - \frac{u^4}{u_h^4} + \alpha^2 \frac{1}{24u_h^2} \left[8u^2(u_h^2 - u^2) - 10u^4 \log 2 + (3u_h^4 + 7u^4) \log \left(1 + \frac{u^2}{u_h^2}\right) \right]^{-2u^2}$$

$$\mathcal{B}(u) = 1 - \alpha^2 \frac{u_h^2}{24} \left[\frac{10u^2}{u_h^2 + u^2} + \log\left(1 + \frac{u^2}{u_h^2}\right) \right], \quad \mathcal{H}(u) = \left(1 + \frac{u^2}{u_h^2}\right)^{\frac{\alpha}{4}} = \left(1 + \frac{u^2}{u_h^2}\right)^{\frac{\alpha}{4}} = \frac{1}{u_h^2} = \frac{1$$

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Conclusions

- ✓ We have presented Confining Anisotropic theories and studied the confinement/deconfinement phase transitions.
- ✓ There are certain stability conditions that constrain the parameters of the theory.
- ✓ The Confinement/Deconfinement phase transitions occur at lower critical Temperature as the anisotropy is increased!
- $\checkmark\,$ We suggest that the anisotropy by itself could instead be a cause of inverse magnetic catalysis.
- ✓ The shear viscosity over entropy density ratio, takes values parametrically lower than $1/4\pi$, and depends on the Temperature as $T^{2-2/z}$.
- ✓ The diffusion (buttery velocity) and growth of chaos occurs faster than isotropic systems.
- Several ways to probe the theory (Mesons, Energy loss of Quarks, Diffusion of Quarks, Speed of Sound, Entanglement Entropy...).

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