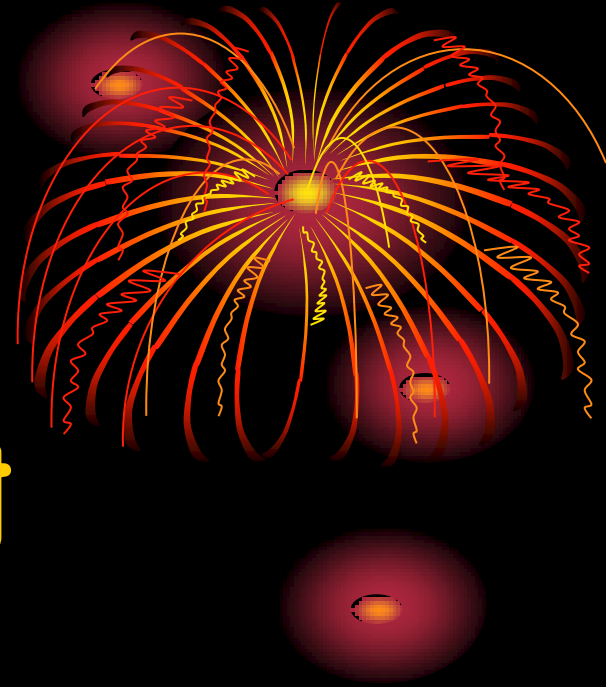


# 11 反射、折射、干涉、繞射

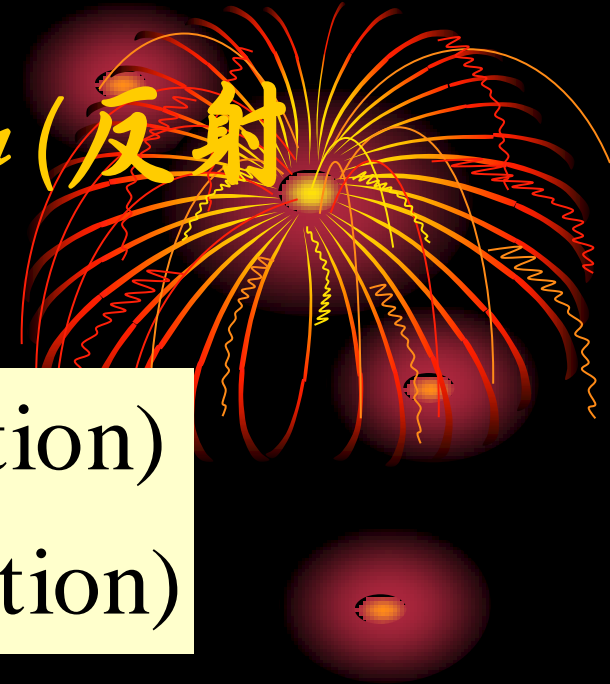


# Sections

1. 反射 (reflection)與折射 (refraction)
2. 干涉 (interference)
3. 繞射 (diffraction)

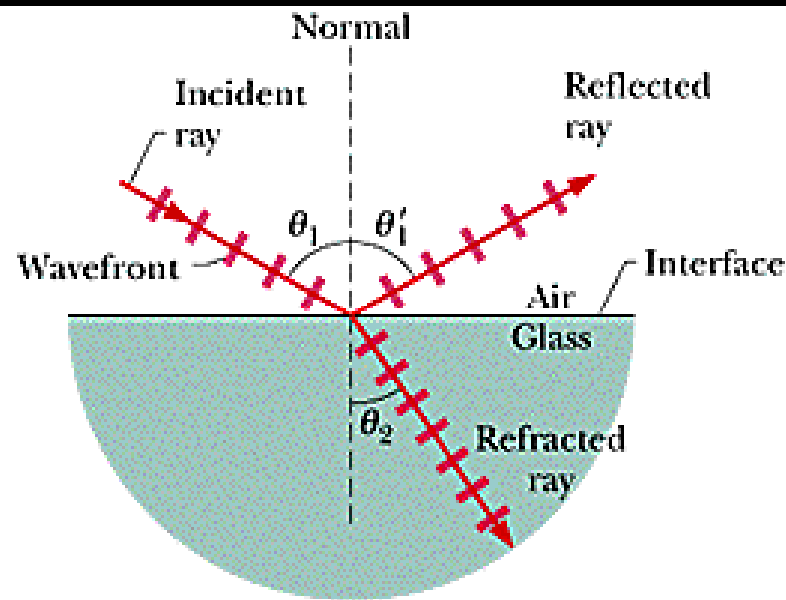
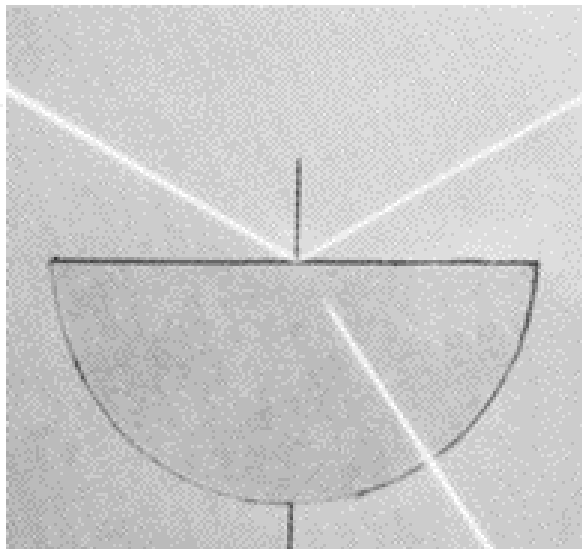


# 11-1 Reflection and Refraction (反射與折射)



$$\theta_1' = \theta_1 \quad (\text{reflection})$$

$$n_2 \sin \theta_2 = n_1 \sin \theta_1 \quad (\text{refraction})$$



# 折射率



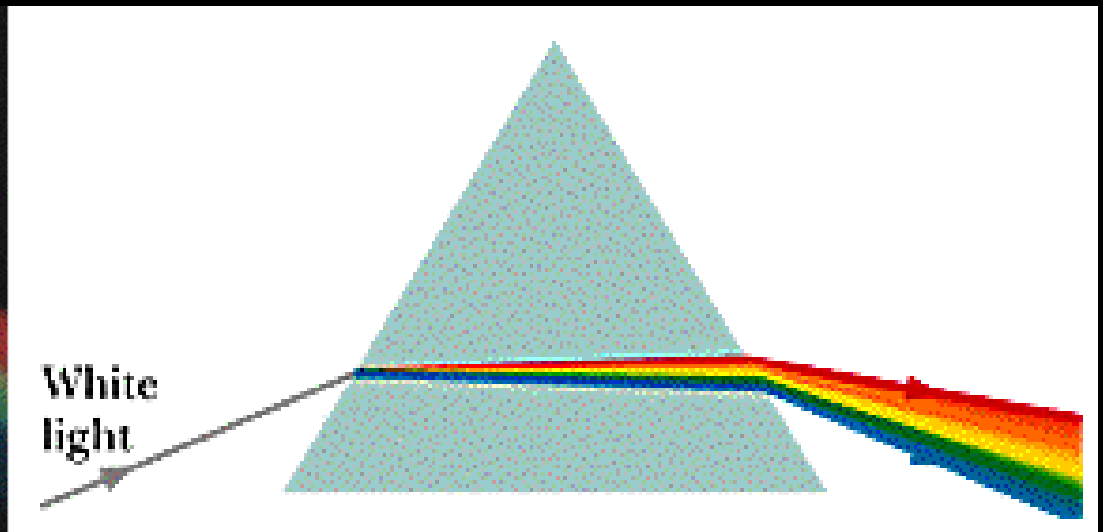
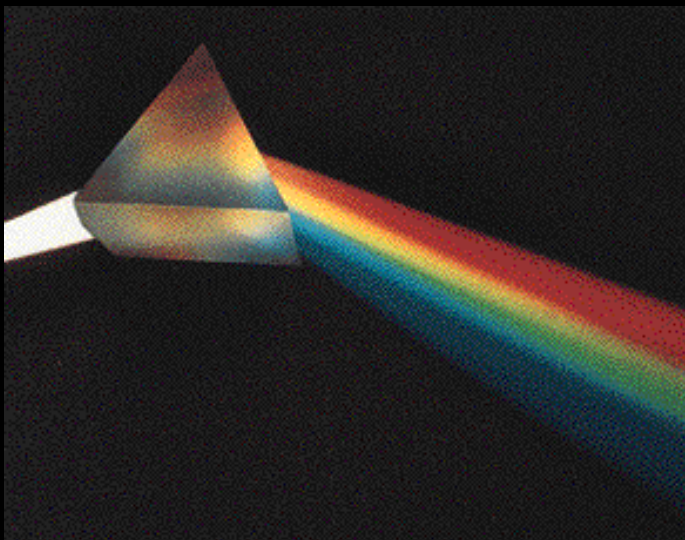
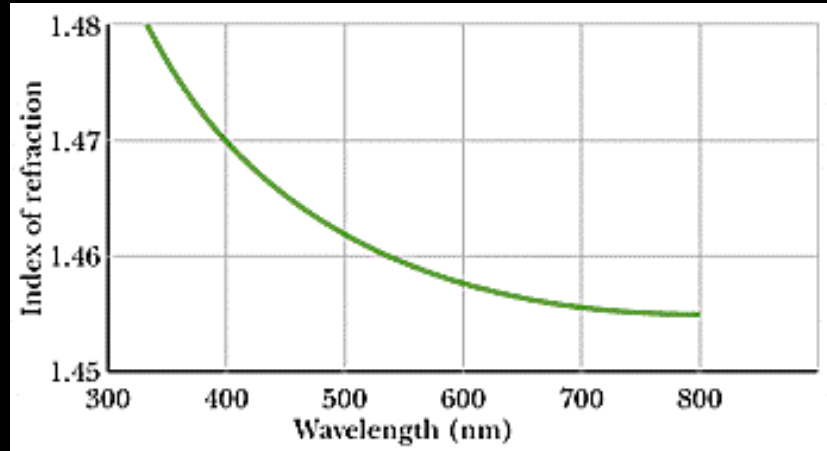
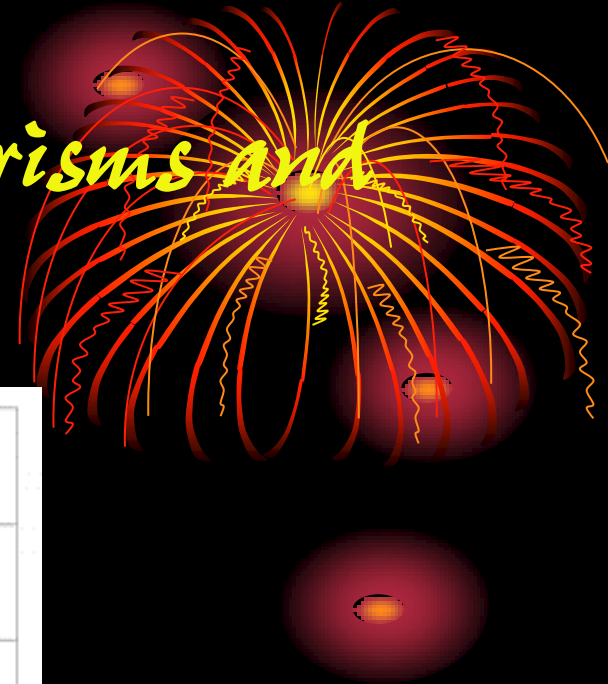
- The Index of Refraction
- The Stealth Aircraft F-117A

TABLE 34-1 SOME INDICES OF REFRACTION<sup>a</sup>

MEDIUM	INDEX	MEDIUM	INDEX
Vacuum	exactly 1	Typical crown glass	1.52
Air (STP) <sup>b</sup>	1.00029	Sodium chloride	1.54
Water (20° C)	1.33	Polystyrene	1.55
Acetone	1.36	Carbon disulfide	1.63
Ethyl alcohol	1.36	Heavy flint glass	1.65
Sugar solution (30%)	1.38	Sapphire	1.77
Fused quartz	1.46	Heaviest flint glass	1.89
Sugar solution (80%)	1.49	Diamond	2.42

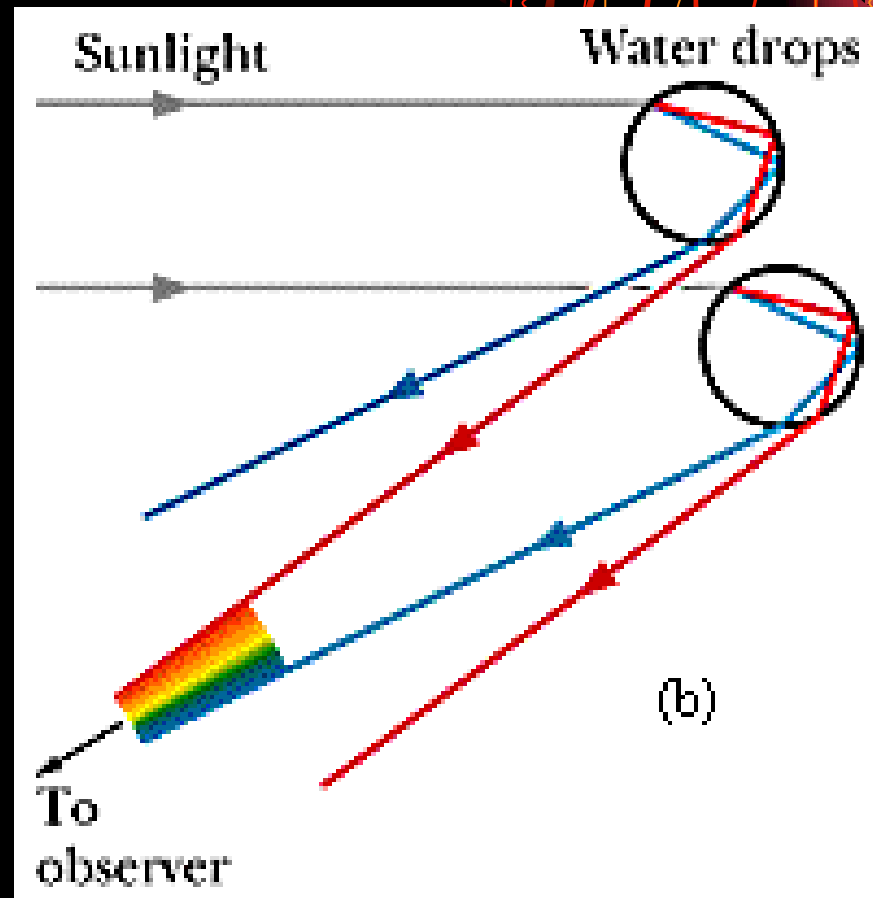


# Chromatic Dispersion - prisms and gratings

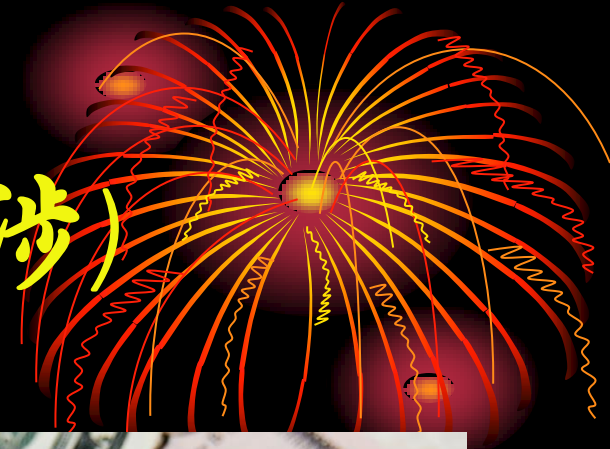




# Rainbows



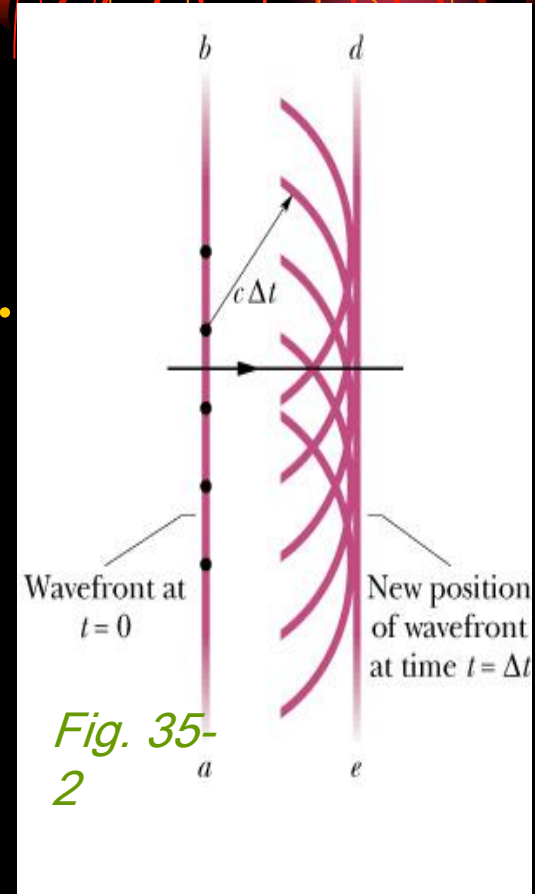
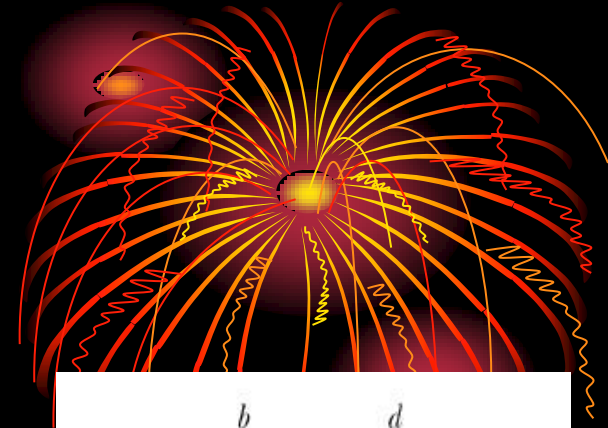
# 11-2 Interference - (干涉)



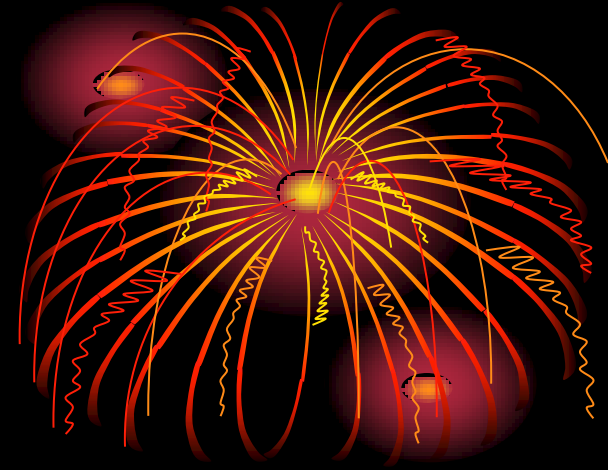
*What produces the blue-green of a Morpho's wing?  
How do colorshifting inks shift colors?*

# Huygens' principle

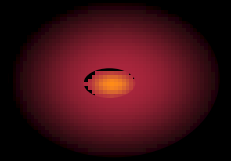
All points on a wavefront serve as point sources of spherical secondary wavelets. After a time  $t$ , the new position of the wavefront will be that of a surface tangent to these secondary wavelets.

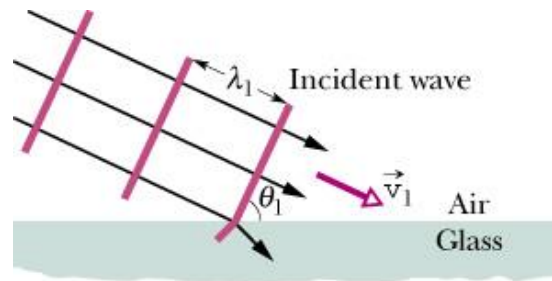




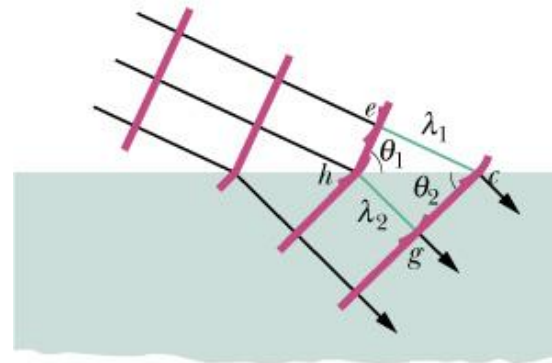


*Law of Refraction from Huygens' principle*





(a)



(b)

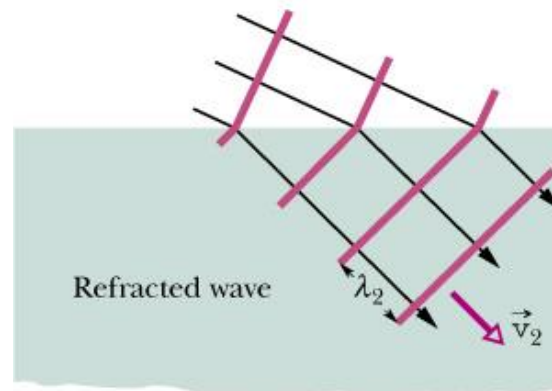
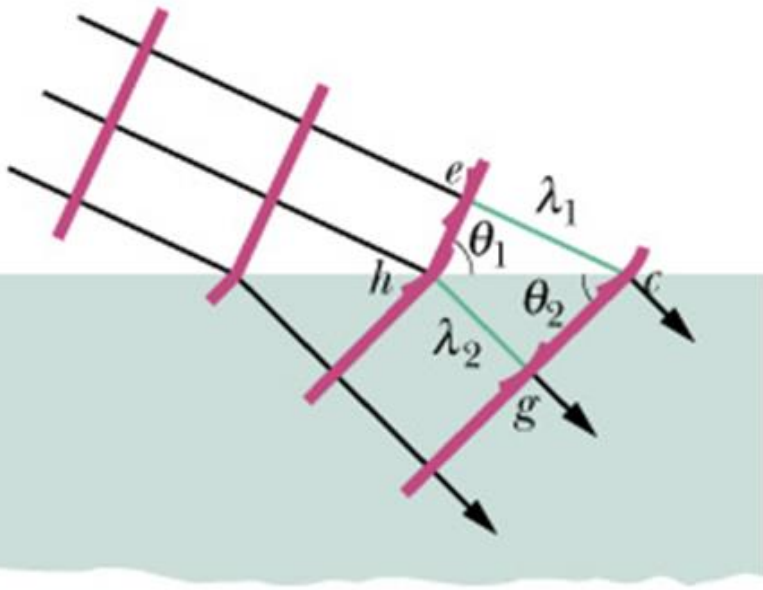


Fig. 35-3 (c)



$$t_{ec} = t_{hg} = \frac{\lambda_1}{v_1} = \frac{\lambda_2}{v_2} \rightarrow \frac{\lambda_1}{\lambda_2} = \frac{v_1}{v_2}$$

$$\sin \theta_1 = \frac{\lambda_1}{hc} \quad (\text{for triangle } hce)$$

$$\sin \theta_2 = \frac{\lambda_2}{hc} \quad (\text{for triangle } hcg)$$

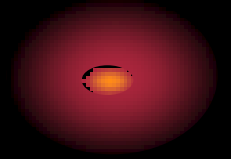
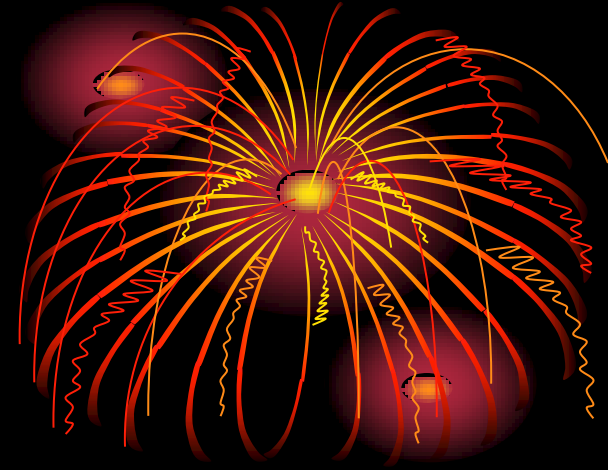
$$\frac{\sin \theta_1}{\sin \theta_2} = \frac{\lambda_1}{\lambda_2} = \frac{v_1}{v_2}$$

$$\text{Index of Refraction: } n = \frac{c}{v}$$

$$n_1 = \frac{c}{v_1} \quad \text{and} \quad n_2 = \frac{c}{v_2}$$

$$\frac{\sin \theta_1}{\sin \theta_2} = \frac{c/n_1}{c/n_2} = \frac{n_2}{n_1}$$

$$\text{Law of Refraction: } n_1 \sin \theta_1 = n_2 \sin \theta_2$$



*Phase Difference, Wavelength and  
Index of Refraction*

# Wavelength and Index of Refraction

$$\frac{\lambda_n}{\lambda} = \frac{v}{c} \rightarrow \lambda_n = \lambda \frac{v}{c} \rightarrow \lambda_n = \frac{\lambda}{n}$$

$$f_n = \frac{v}{\lambda_n} = \frac{c/n}{\lambda/n} = \frac{c}{\lambda} = f$$

The frequency of light in a medium is the same as it is in vacuum



# Phase Difference

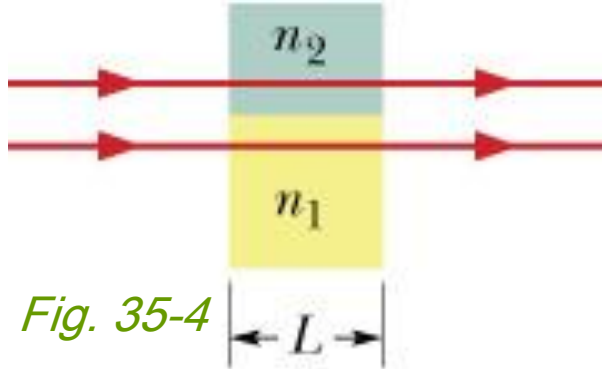


Fig. 35-4

Since wavelengths in  $n_1$  and  $n_2$  are different, the two beams may no longer be in phase

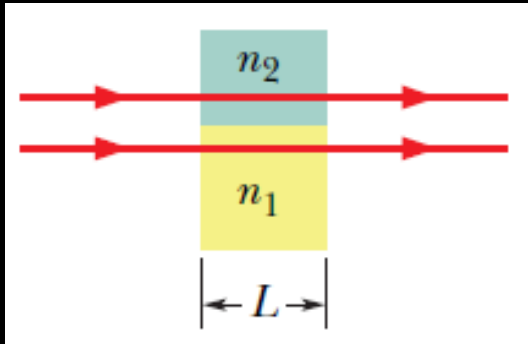
$$\text{Number of wavelengths in } n_1: N_1 = \frac{L}{\lambda_{n_1}} = \frac{L}{\lambda/n_1} = \frac{Ln_1}{\lambda}$$

$$\text{Number of wavelengths in } n_2: N_2 = \frac{L}{\lambda_{n_2}} = \frac{L}{\lambda/n_2} = \frac{Ln_2}{\lambda}$$

$$\text{Assuming } n_2 > n_1: N_2 - N_1 = \frac{Ln_2}{\lambda} - \frac{Ln_1}{\lambda} = \frac{L}{\lambda}(n_2 - n_1)$$

$N_2 - N_1 = 1/2$  wavelength  $\rightarrow$  destructive interference

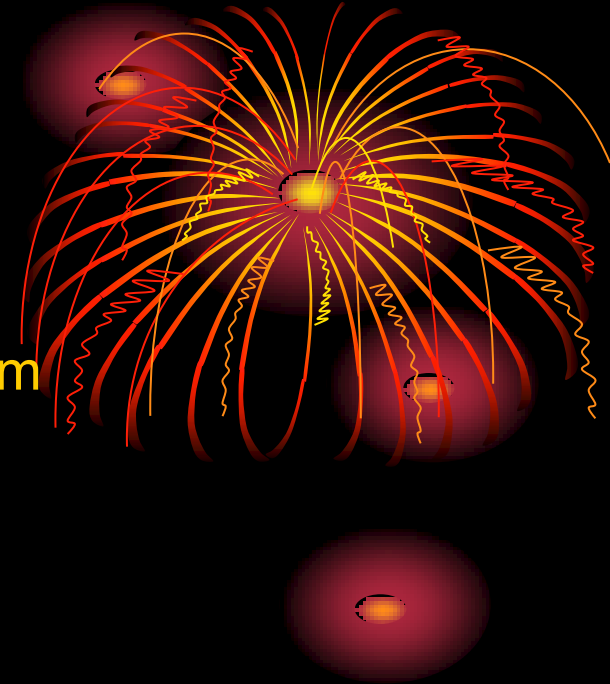
# Ex.11-1 35-1



wavelength 550.0 nm

$n_2=1.600$  and

$L = 2.600 \text{ m}$

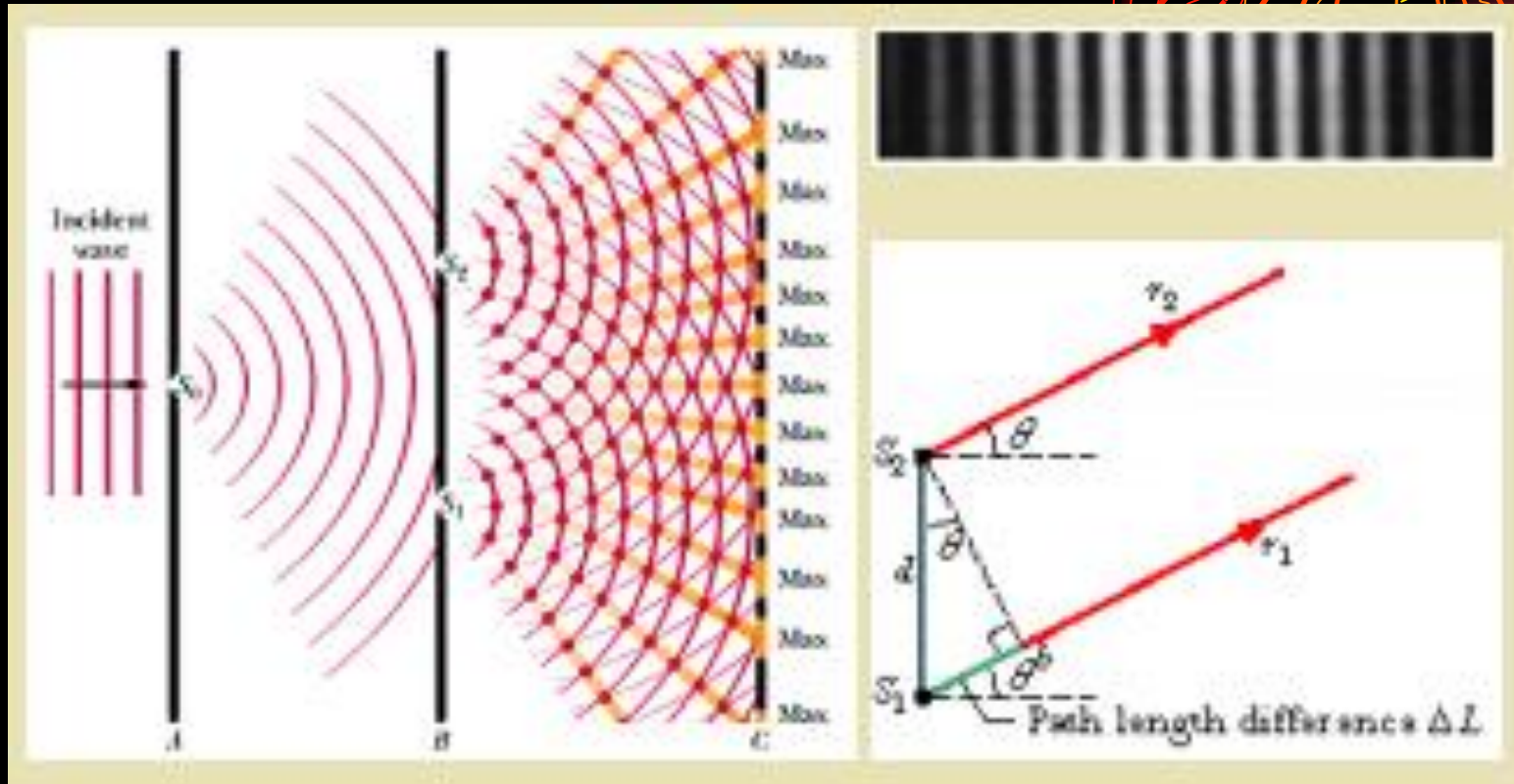


$$\begin{aligned} N_2 - N_1 &= \frac{L}{\lambda} (n_2 - n_1) \\ &= \frac{2.600 \times 10^{-6} \text{ m}}{5.500 \times 10^{-7} \text{ m}} (1.600 - 1.000) \\ &= 2.84. \end{aligned} \quad (\text{Ans})$$

phase difference = 17.8 rad  $\approx$  1020°

effective phase difference = 0.84 wavelength

# Young's Experiment



$$d \sin \theta = m\lambda, \quad \text{for } m = 0, 1, 2, \dots \quad (\text{maxima—bright fringes}).$$

$$d \sin \theta = (m + \frac{1}{2})\lambda, \quad \text{for } m = 0, 1, 2, \dots \quad (\text{minima—dark fringes}).$$

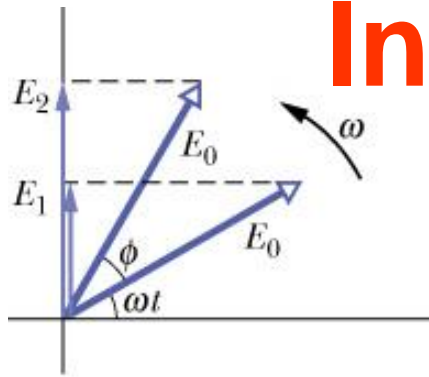
# Coherence

Two sources to produce an interference that is stable over time, if their light has a *phase relationship* that does not change with time:  $E(t)=E_0\cos(\omega t+\phi)$

**Coherent sources:** Phase  $\phi$  must be well defined and constant. When waves from coherent sources meet, stable interference can occur — laser light (produced by cooperative behavior of atoms)

**Incoherent sources:**  $\phi$  jitters randomly in time, no stable interference occurs — sunlight

# Intensity and phase



(a)

$$E(t) = E_0 \sin \omega t + E_0 \sin(\omega t + \phi) = ?$$

$$E = 2(E_0 \cos \beta) = 2E_0 \cos \frac{1}{2} \phi$$

$$E^2 = 4E_0^2 \cos^2 \frac{1}{2} \phi$$

$$\frac{I}{I_0} = \frac{E^2}{E_0^2} = 4 \cos^2 \frac{1}{2} \phi \rightarrow I = 4I_0 \cos^2 \frac{1}{2} \phi$$

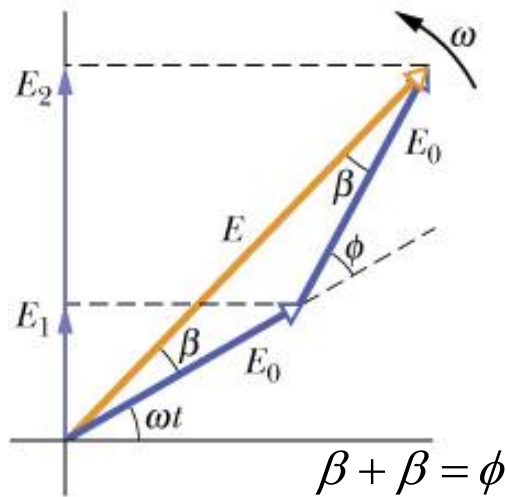
**Eq. 35-22**

$$\frac{\left( \begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right)}{2\pi} = \frac{\left( \begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right)}{\lambda}$$

$$\left( \begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right) = \frac{2\pi}{\lambda} \left( \begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right)$$

$$\phi = \frac{2\pi}{\lambda} (d \sin \theta)$$

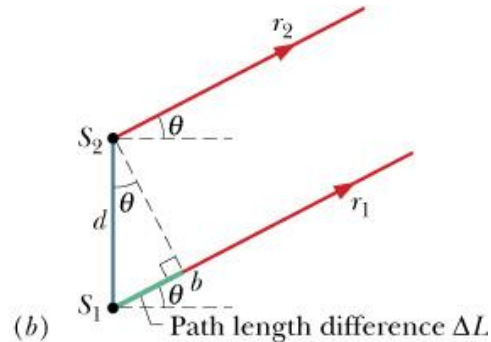
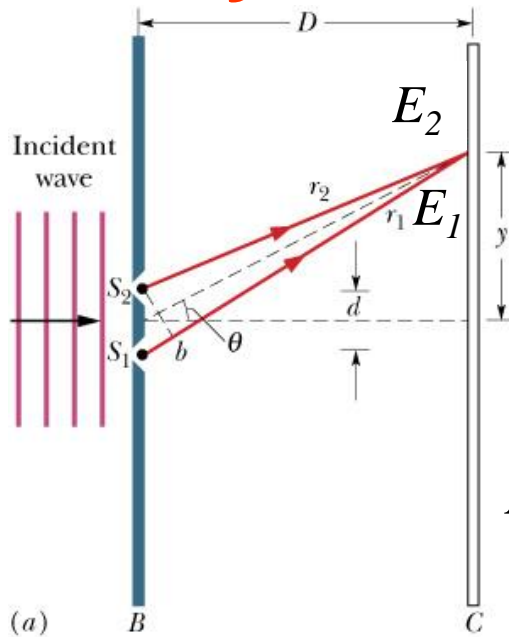
**Eq. 35-23**



**Fig. 35-13** (b)



# Intensity in Double-Slit Interference



$$E_1 = E_0 \sin \omega t \quad \text{and} \quad E_2 = E_0 \sin(\omega t + \phi)$$

$$I = 4I_0 \cos^2 \frac{1}{2} \phi$$

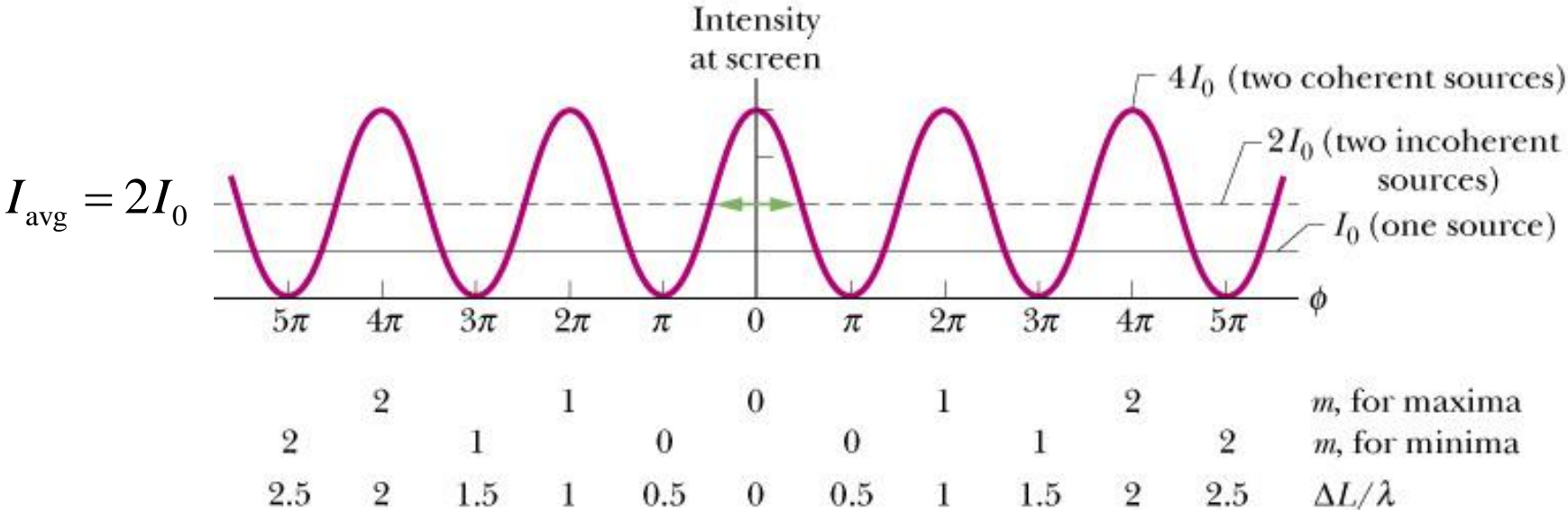
$$\phi = \frac{2\pi d}{\lambda} \sin \theta$$

maxima when:  $\frac{1}{2} \phi = m\pi$  for  $m = 0, 1, 2, \dots \rightarrow \phi = 2m\pi = \frac{2\pi d}{\lambda} \sin \theta$

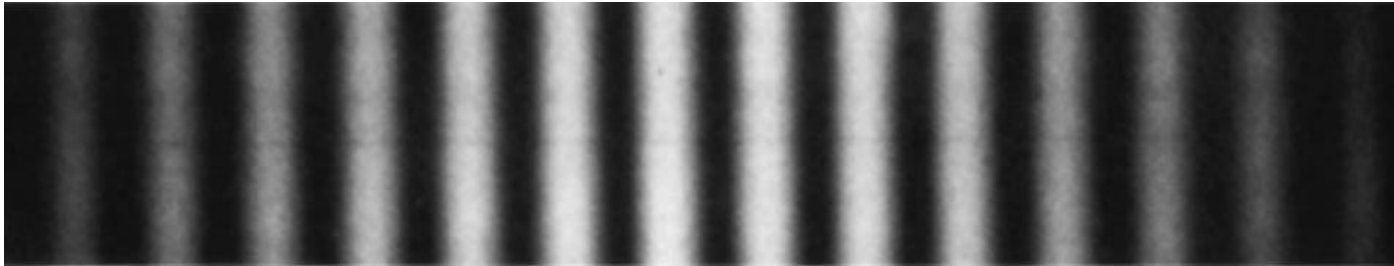
$\rightarrow d \sin \theta = m\lambda$  for  $m = 0, 1, 2, \dots$  (maxima)

minima when:  $\frac{1}{2} \phi = (m + \frac{1}{2})\pi \rightarrow d \sin \theta = (m + \frac{1}{2})\lambda$  for  $m = 0, 1, 2, \dots$  (minima)

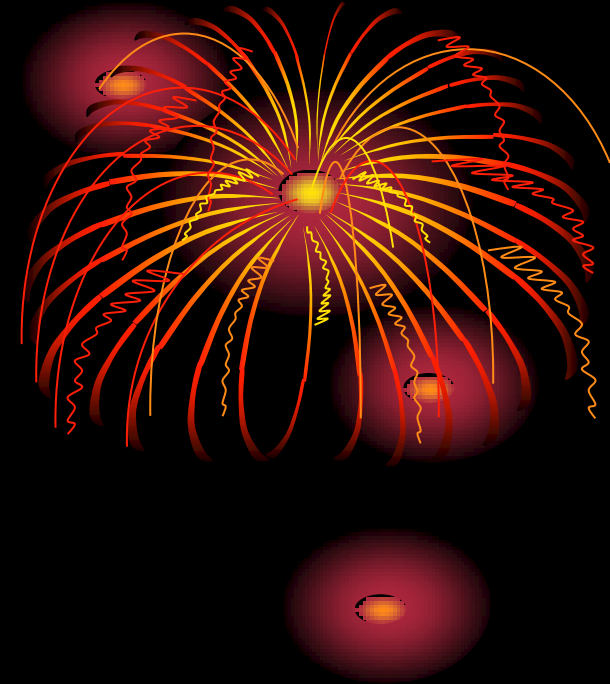
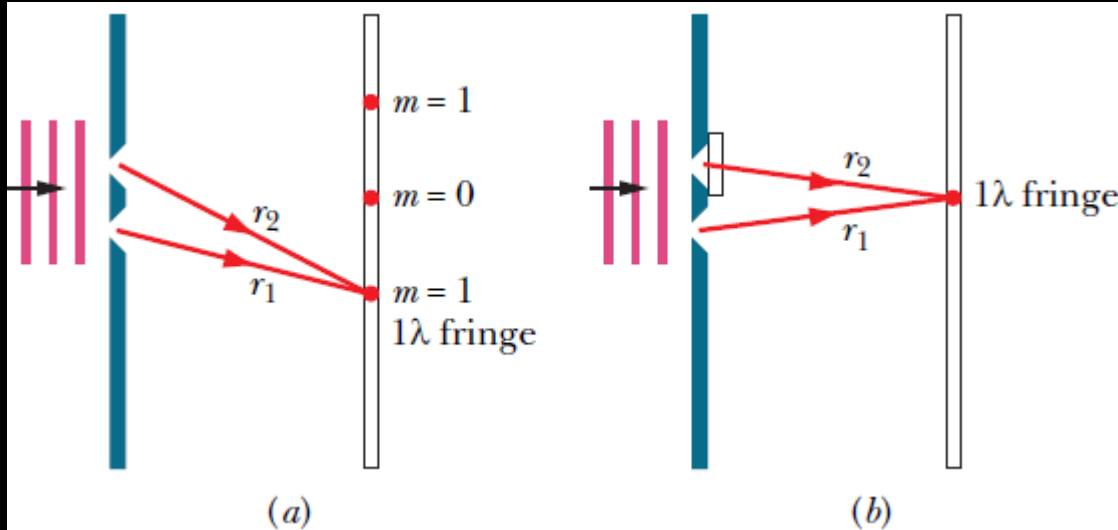
# Intensity in Double-Slit Interference



*Fig. 35-12*



# Ex.11-2 35-2



$$N_2 - N_1 = \frac{L}{\lambda} (n_2 - n_1)$$

$$L = \frac{\lambda(N_2 - N_1)}{n_2 - n_1} = \frac{(600 \times 10^{-9} \text{ m})(1)}{1.50 - 1.00}$$

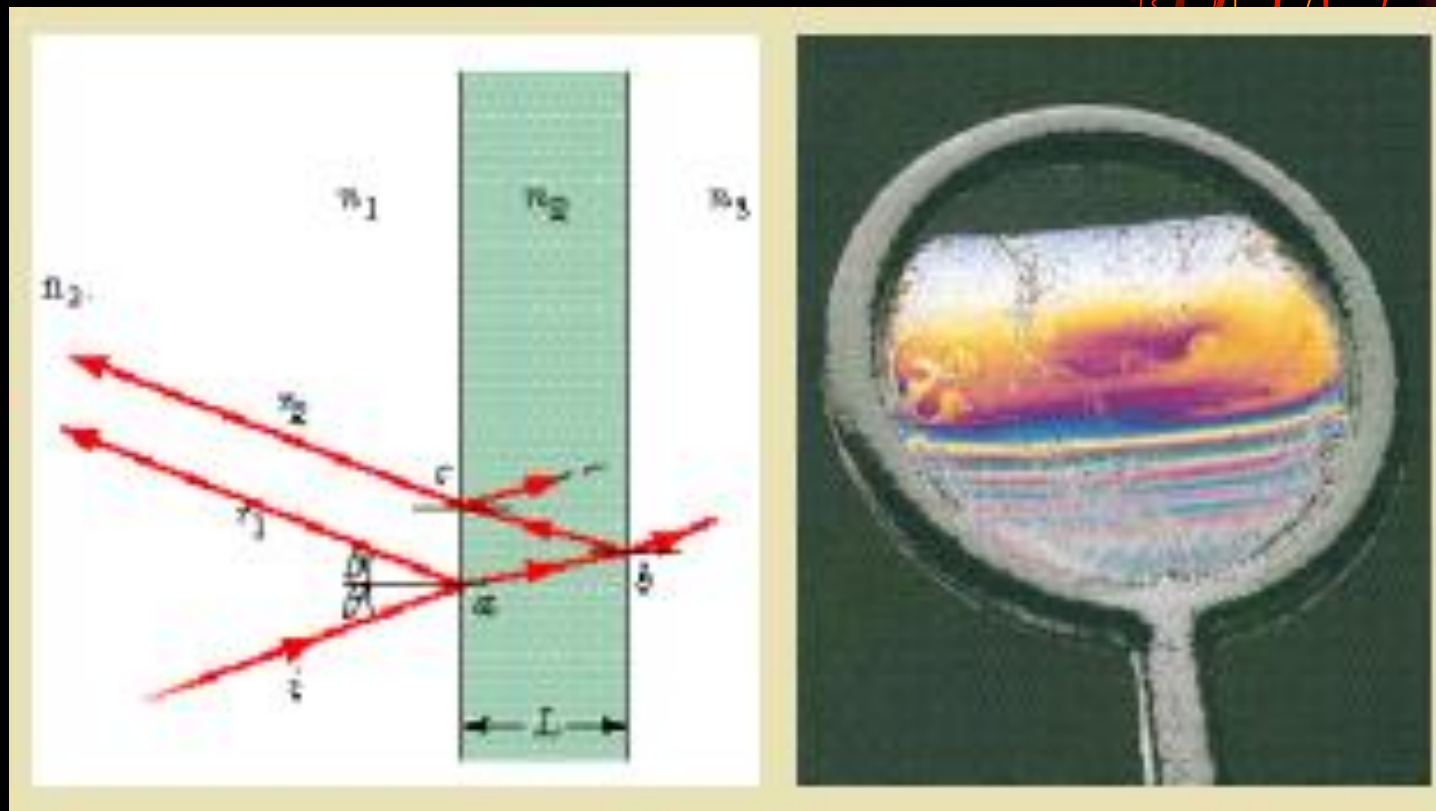
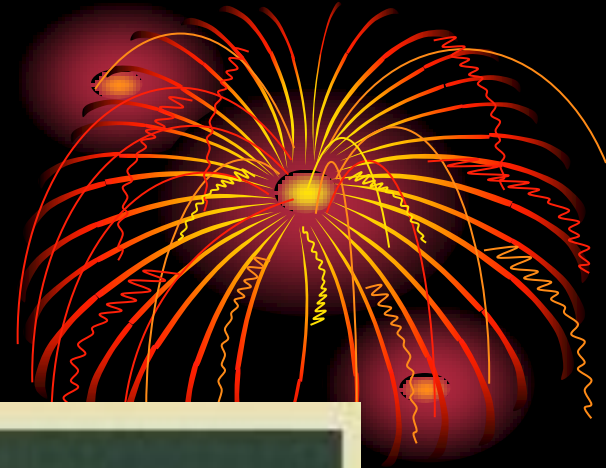
$$= 1.2 \times 10^{-6} \text{ m.} \quad (\text{Answer})$$

wavelength 600 nm

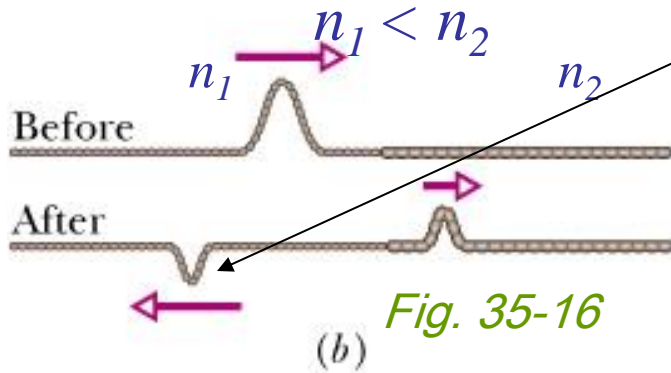
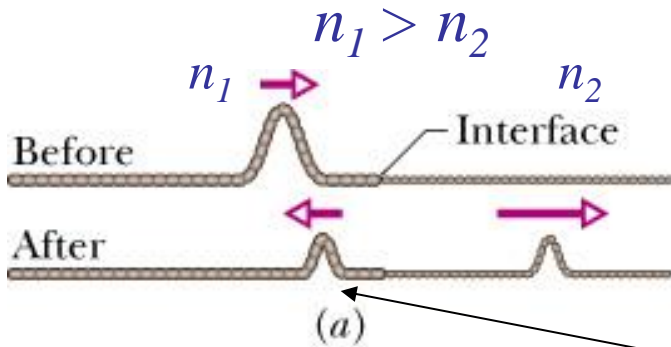
$n_2=1.5$  and

$m = 1 \rightarrow m = 0$

# Interference from Thin Films



# Reflection Phase Shifts



**Reflection**  
Off lower index  
Off higher index

**Reflection Phase Shift**  
0  
0.5 wavelength



# Phase Difference in Thin-Film Interference

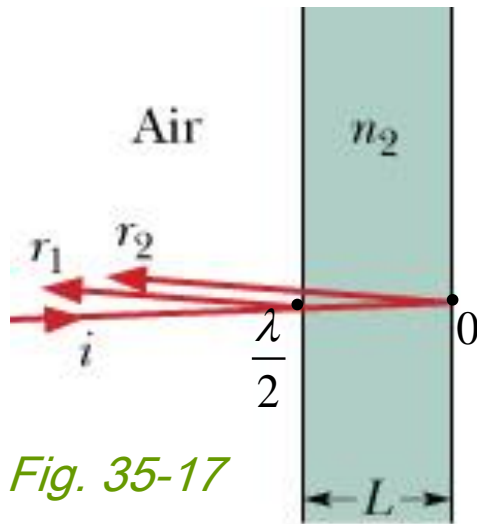


Fig. 35-17

Three effects can contribute to the phase difference between  $r_1$  and  $r_2$ .

1. Differences in reflection conditions
2. Difference in path length traveled.
3. Differences in the media in which the waves travel. One must use the wavelength in each medium ( $\lambda / n$ ), to calculate the phase.

# Equations for Thin-Film Interference

$\frac{1}{2}$  wavelength phase difference to difference in reflection of  $r_1$  and  $r_2$

$$2L = \frac{\text{odd number}}{2} \times \text{wavelength} = \frac{\text{odd number}}{2} \times \lambda_{n_2} \quad (\text{in-phase waves})$$

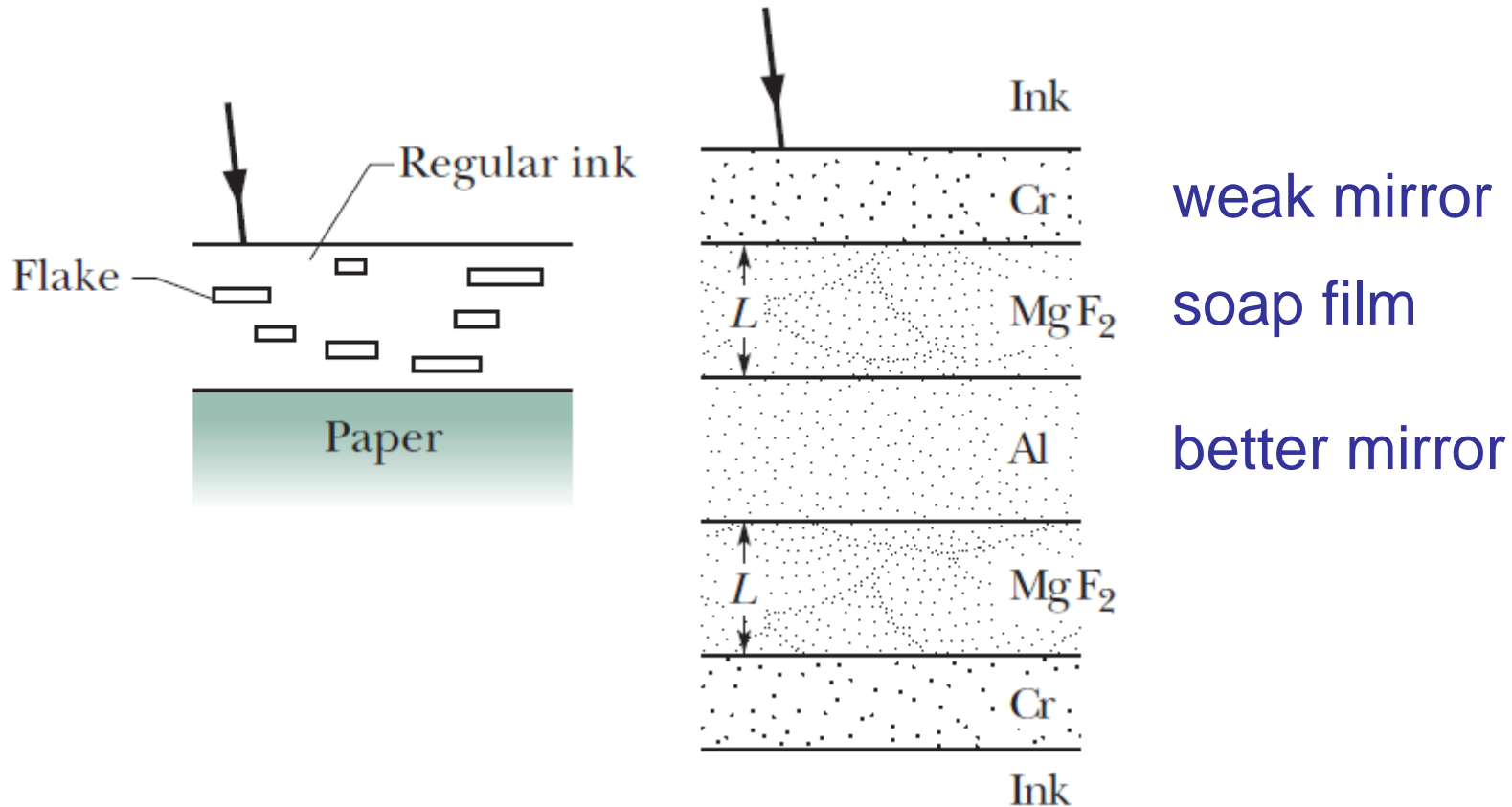
$$2L = \text{integer} \times \text{wavelength} = \text{integer} \times \lambda_{n_2} \quad (\text{out-of-phase waves})$$

$$\lambda_{n_2} = \frac{\lambda}{n_2}$$

$$2L = \left(m + \frac{1}{2}\right) \frac{\lambda}{n_2} \quad \text{for } m = 0, 1, 2, \dots \quad (\text{maxima-- bright film in air})$$

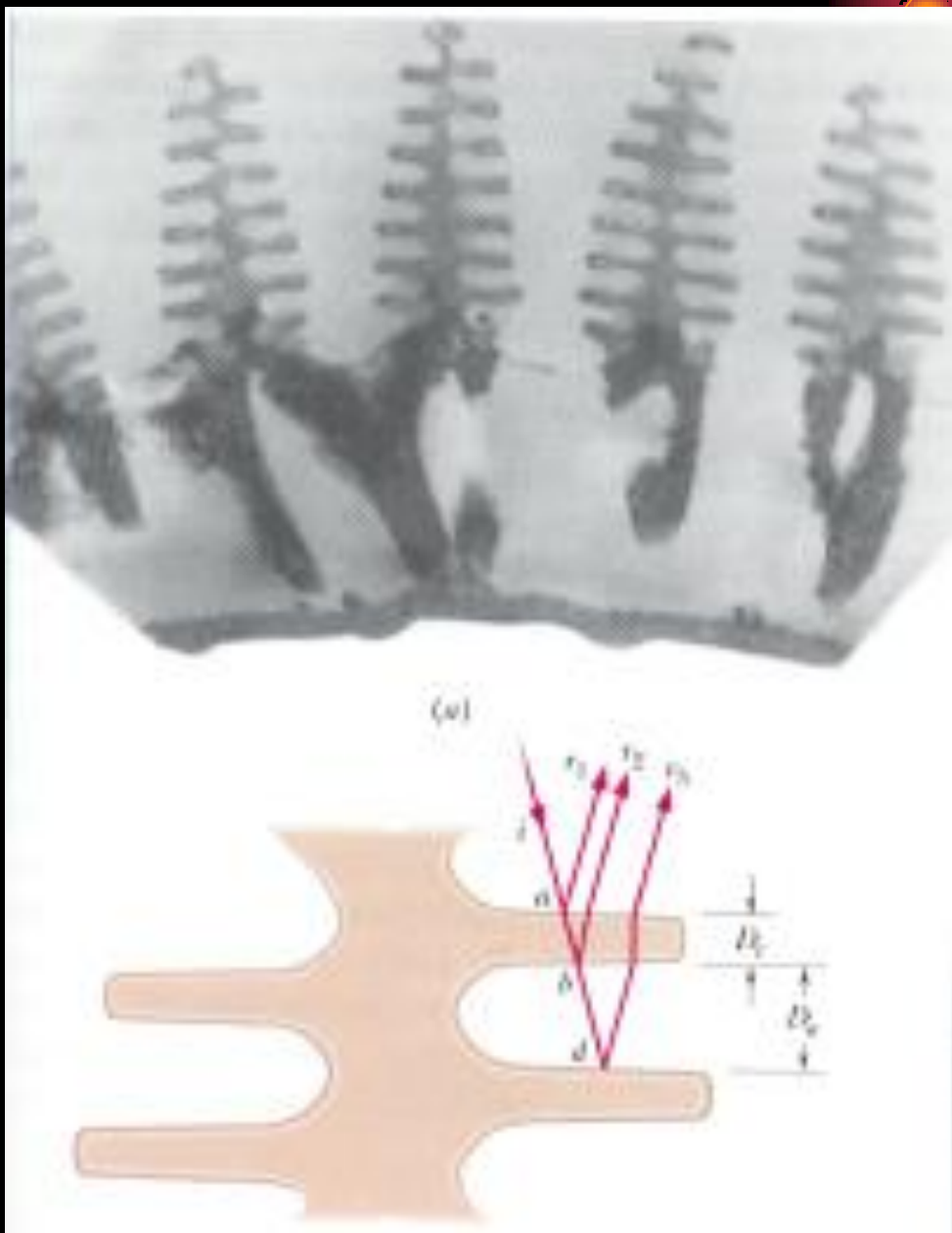
$$2L = m \frac{\lambda}{n_2} \quad \text{for } m = 0, 1, 2, \dots \quad (\text{minima-- dark film in air})$$

# Color Shifting by Paper Currencies, paints and Morpho Butterflies

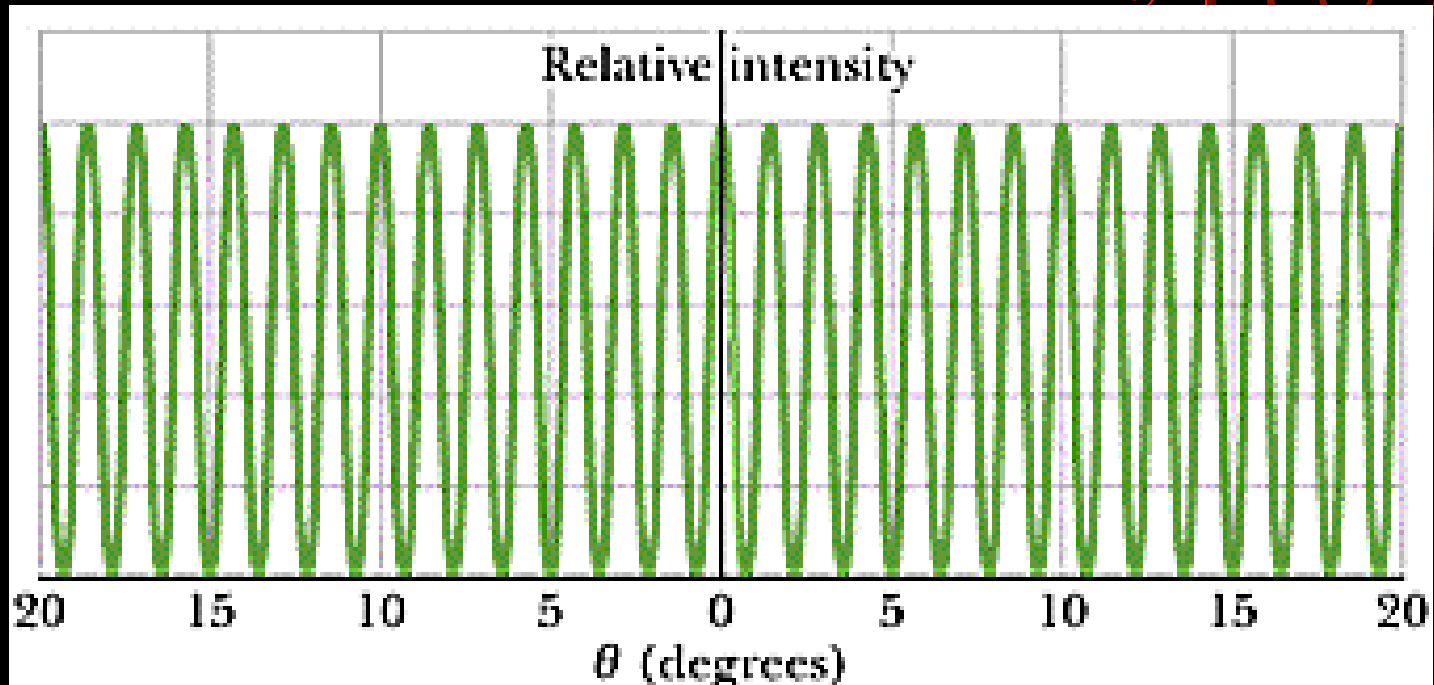
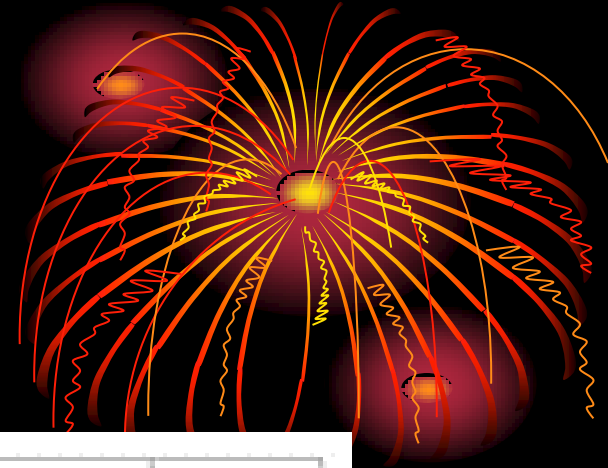


looking directly down : red or red-yellow  
tilting : green

# 大藍魔爾蝴蝶



# 雙狹縫干涉之強度



# Ex.11-3 35-3 Brightened reflected light from a water film



$$2L = \frac{\text{odd number}}{2} \times \frac{\lambda}{n_2},$$

thickness 320 nm

$n_2=1.33$

$$2L = (m + \frac{1}{2}) \frac{\lambda}{n_2}.$$

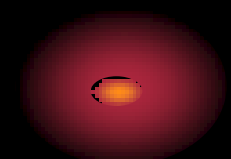
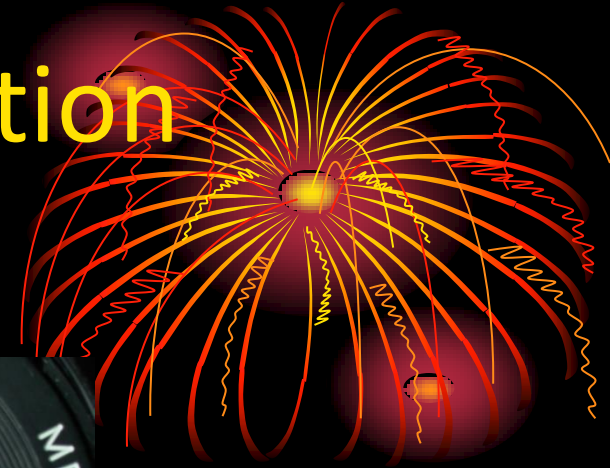
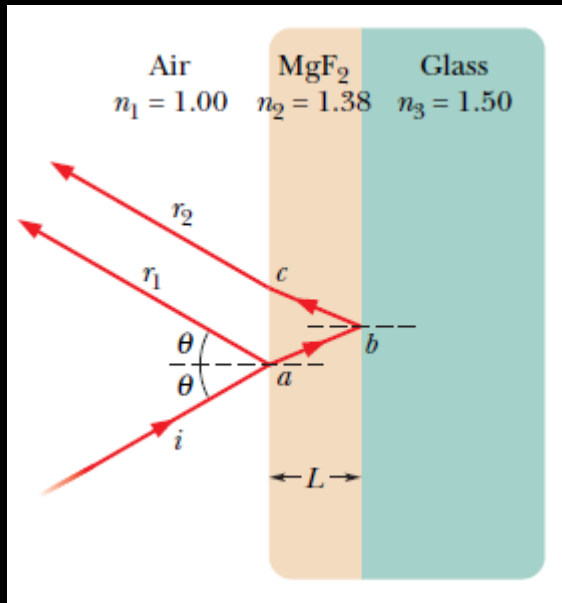
$$\lambda = \frac{2n_2L}{m + \frac{1}{2}} = \frac{(2)(1.33)(320 \text{ nm})}{m + \frac{1}{2}} = \frac{851 \text{ nm}}{m + \frac{1}{2}}.$$

$m = 0$ , 1700 nm, infrared

$m = 1$ , 567 nm, yellow-green ←

$m = 2$ , 340 nm, ultraviolet

# Ex.11-4 35-4 anti-reflection coating



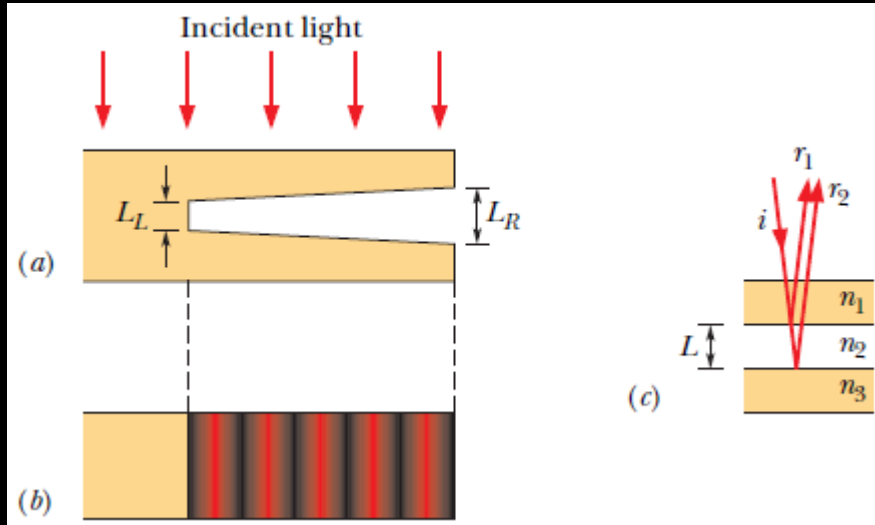
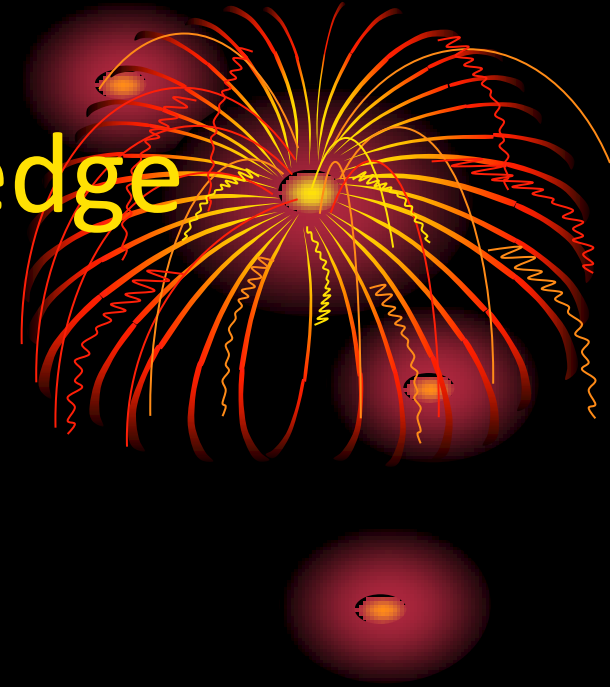
$$2L = \frac{\text{odd number}}{2} \times \frac{\lambda}{n_2}.$$

$$L = \left(m + \frac{1}{2}\right) \frac{\lambda}{2n_2}, \quad \text{for } m = 0, 1, 2, \dots$$

$$L = \frac{\lambda}{4n_2} = \frac{550 \text{ nm}}{(4)(1.38)} = 99.6 \text{ nm}.$$



# Ex.11-5 35-5 thin air wedge



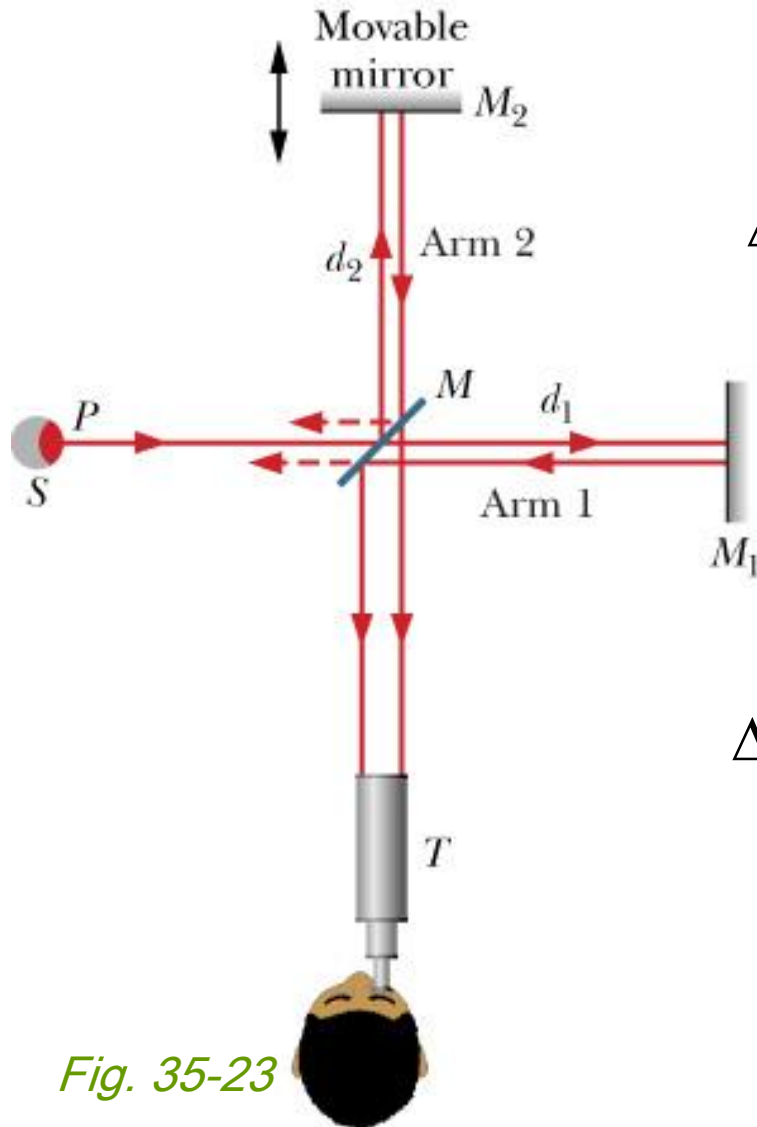
$$2L = \text{integer} \times \frac{\lambda}{n_2},$$

$$2L = m \frac{\lambda}{n_2}, \quad \text{for } m = 0, 1, 2, \dots$$

$$L_L = (m_L) \frac{\lambda}{2n_2}, \quad L_R = (m_L + 5) \frac{\lambda}{2n_2}.$$

$$\begin{aligned} \Delta L = L_R - L_L &= \frac{(m_L + 5)\lambda}{2n_2} - \frac{m_L\lambda}{2n_2} = \frac{5}{2} \frac{\lambda}{n_2} \\ &= 1.58 \times 10^{-6} \text{ m.} \quad (\text{Answer}) \end{aligned}$$

# Michelson Interferometer



$$\Delta L = 2d_1 - 2d_2 \quad (\text{interferometer})$$

$$\Delta L_m = 2L \quad (\text{slab of material of thickness } L \text{ placed in front of } M_1)$$

Fig. 35-23

# Determining Material thickness L

$$N_m = \frac{2L}{\lambda_m} = \frac{2Ln}{\lambda} \quad (\text{number of wavelengths in slab of material})$$

$$N_a = \frac{2L}{\lambda} \quad (\text{number of wavelengths in same thickness of air})$$

$$N_m - N_a = \frac{2Ln}{\lambda} - \frac{2L}{\lambda} = \frac{2L}{\lambda} (n-1) \quad (\text{difference in wavelengths for paths with and without thin slab})$$

# Problem 35-81

In Fig. 35-49, an airtight chamber of length  $d = 5.0 \text{ cm}$  is placed in one of the arms of a Michelson interferometer. (The glass window on each end of the chamber has negligible thickness.) Light of wavelength  $\lambda = 500 \text{ nm}$  is used. Evacuating the air from the chamber causes a shift of 60 bright fringes. From these data and to six significant figures, find the index of refraction of air at atmospheric pressure.

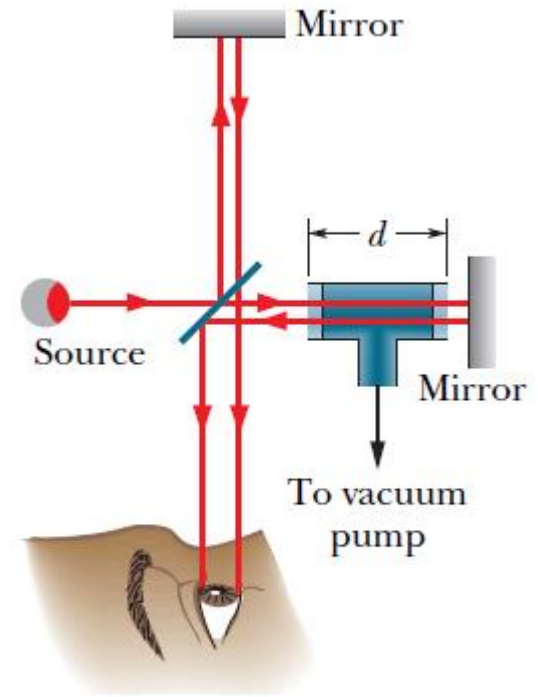


FIG. 35-49 Problem 81.

# Solution to Problem 35-81

$\phi_1$  the phase difference with air ; 2 : vacuum

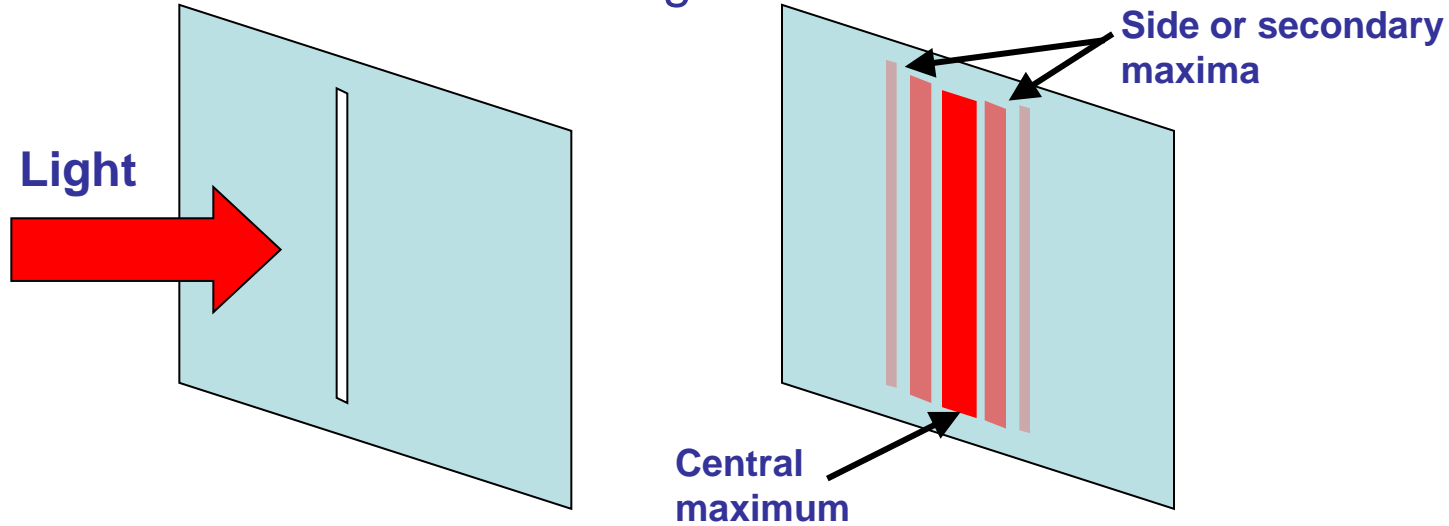
$$\phi_1 - \phi_2 = 2L \frac{2\pi n}{\lambda} - \frac{2\pi}{\lambda} \frac{4\pi(n-1)L}{\lambda}$$

$$\frac{4\pi(n-1)L}{\lambda} = 2N\pi \quad N \text{ fringes}$$

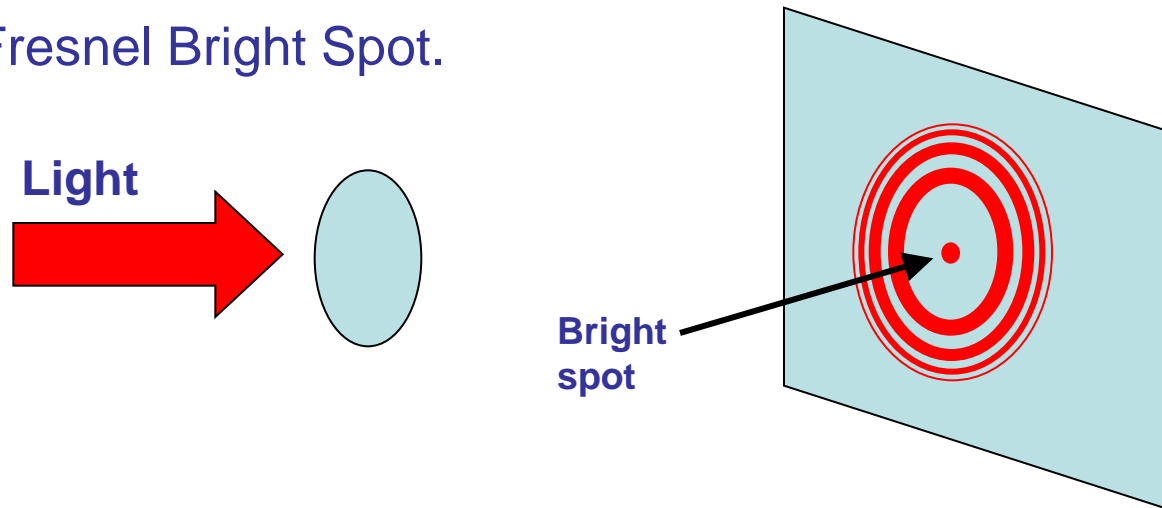
$$n = 1 + \frac{N\lambda}{2L} = 1 + \frac{600 \times 10^{-9} \text{ m}}{2 \times 50 \times 10^{-2} \text{ m}} = 1.00030 .$$

# 11-3 Diffraction and the Wave Theory of Light

Diffraction Pattern from a single narrow slit.

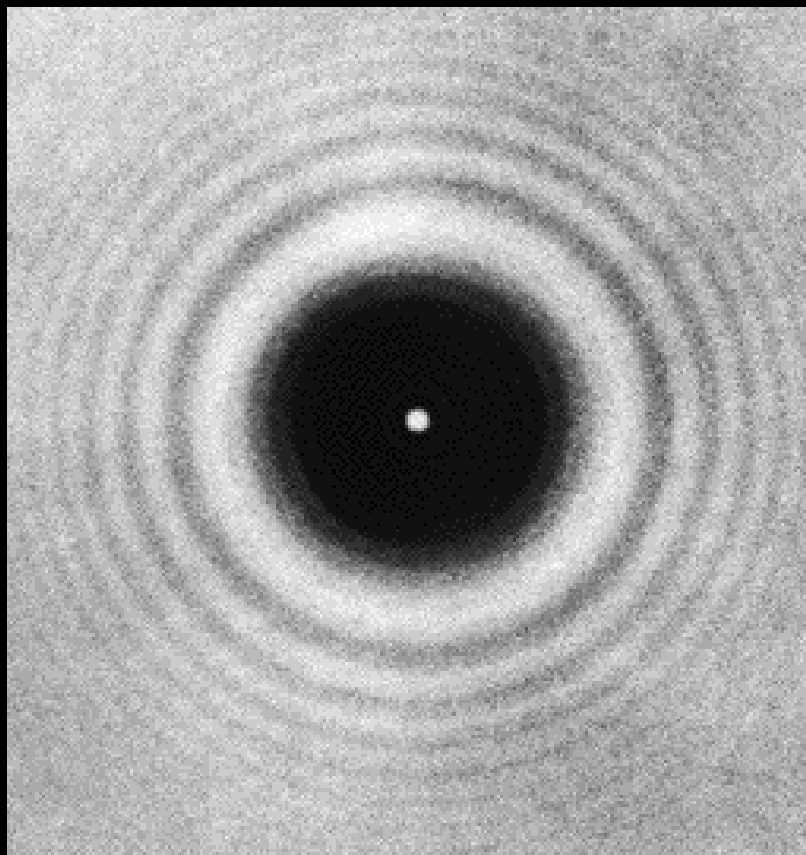
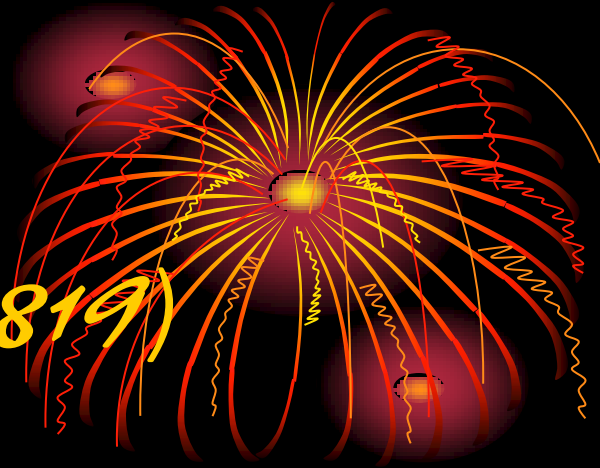


Fresnel Bright Spot.



These patterns cannot be explained using geometrical optics (Ch. 34)!

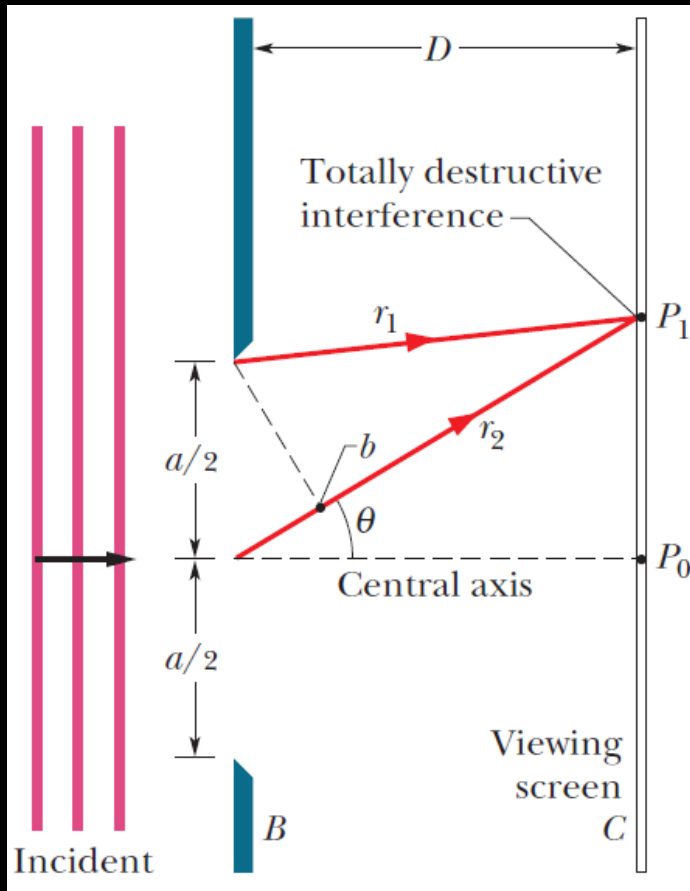
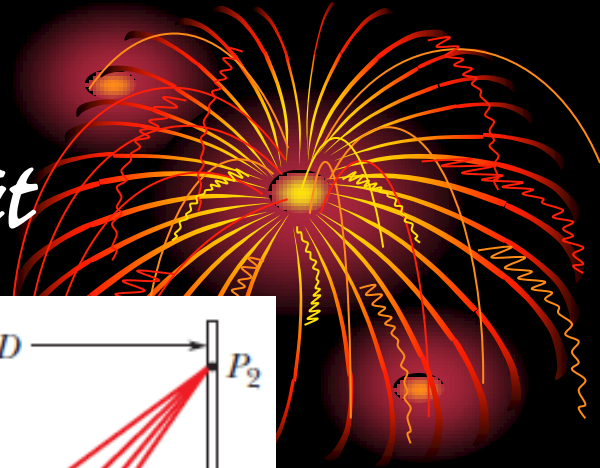
# The Fresnel Bright Spot (1819)



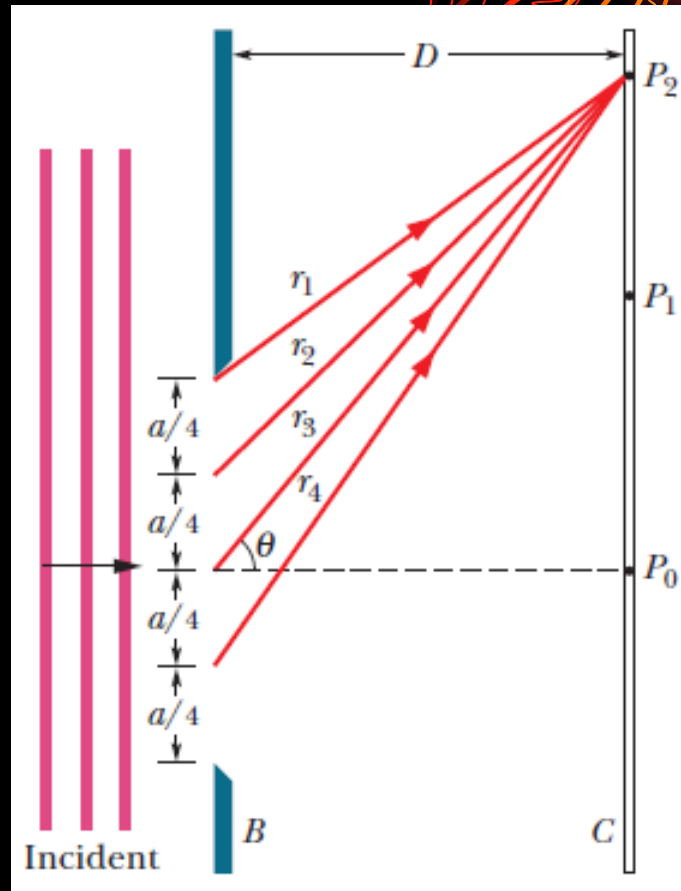
- Newton
  - corpuscle
- Poisson
- Fresnel
  - wave



# Diffraction by a single slit



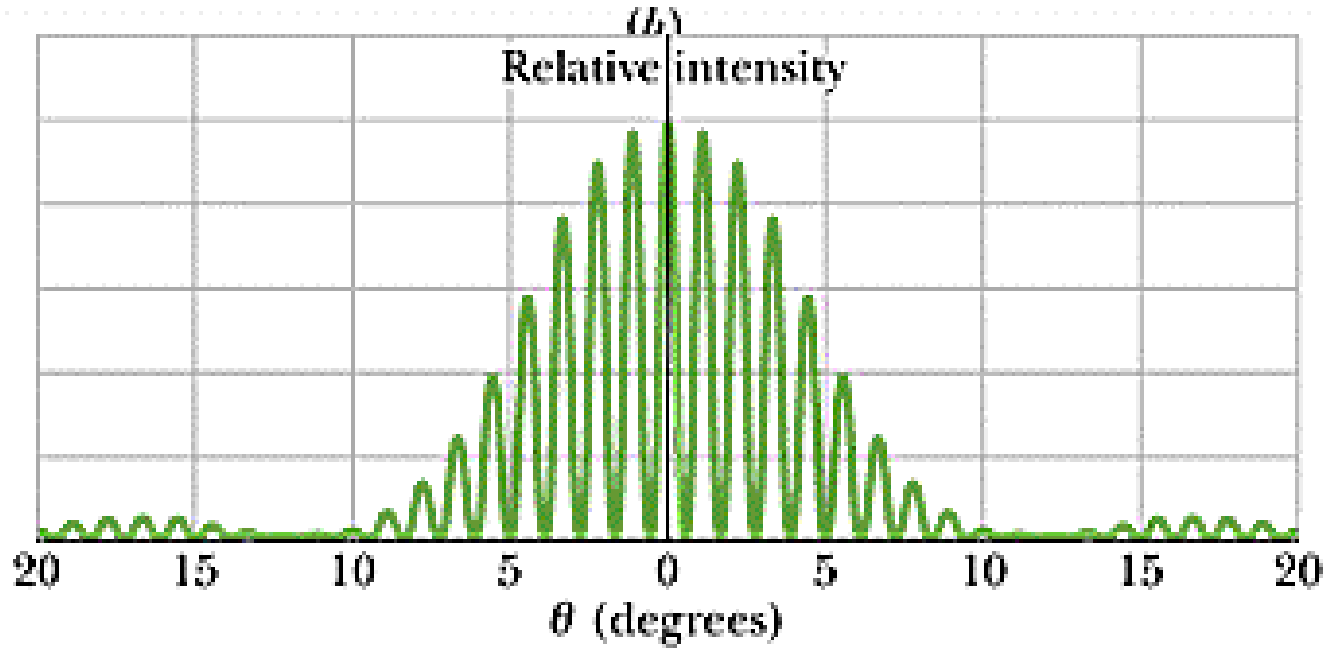
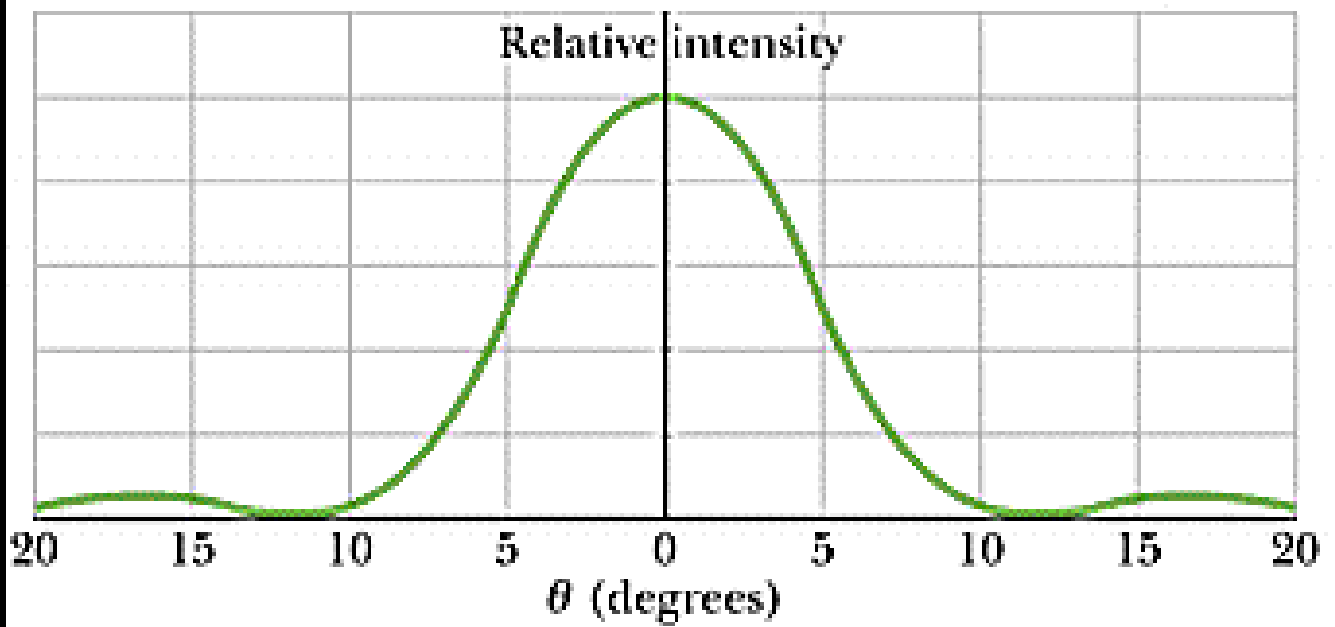
$$a \sin \theta = \lambda \quad (1^{\text{st}} \text{ minima})$$



$$a \sin \theta = 2\lambda \quad (2^{\text{nd}} \text{ minima})$$

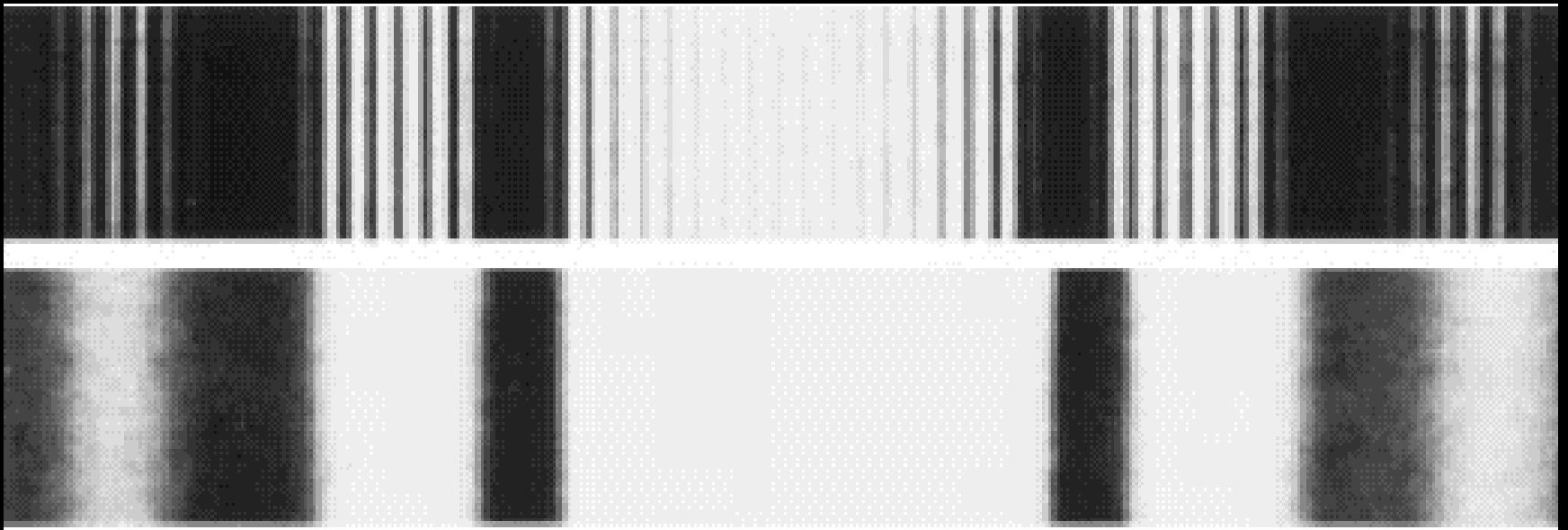
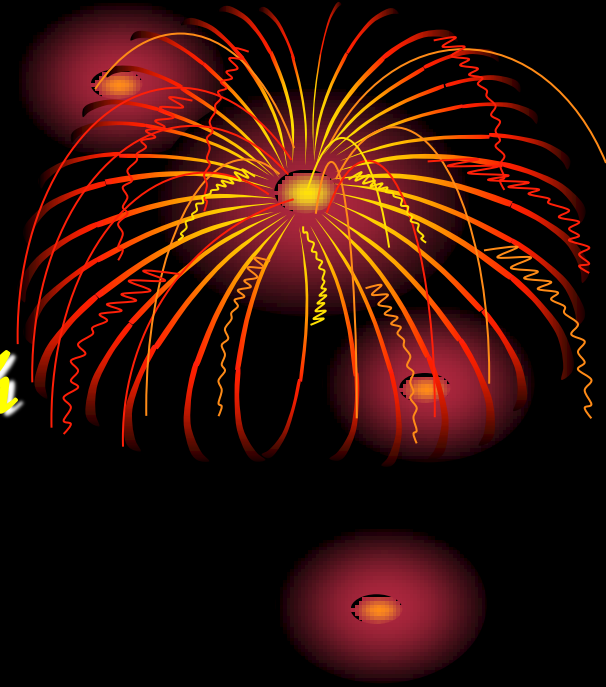
$$a \sin \theta = m\lambda, \quad \text{for } m = 1, 2, 3, \dots \quad (\text{minima — dark fringes}).$$

單狹縫繞射之強度

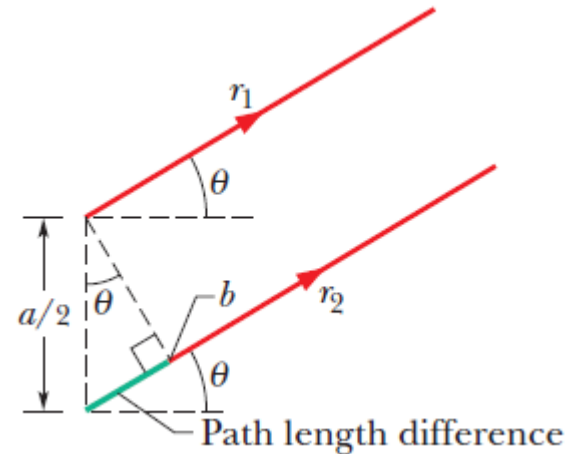
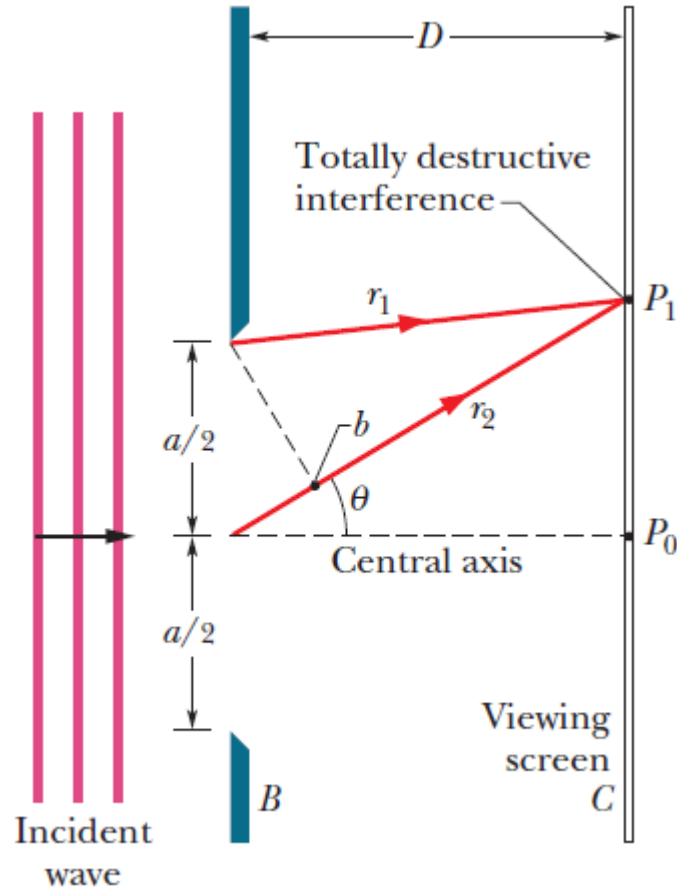


# 雙狹縫與單狹縫

- *Double-slit diffraction (with interference)*
- *Single-slit diffraction*



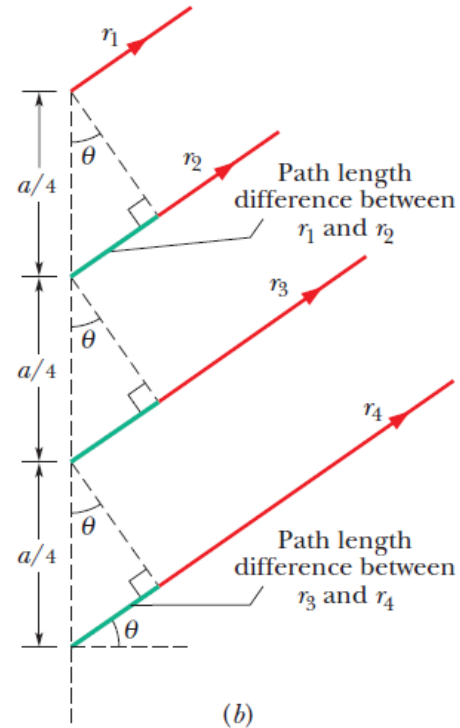
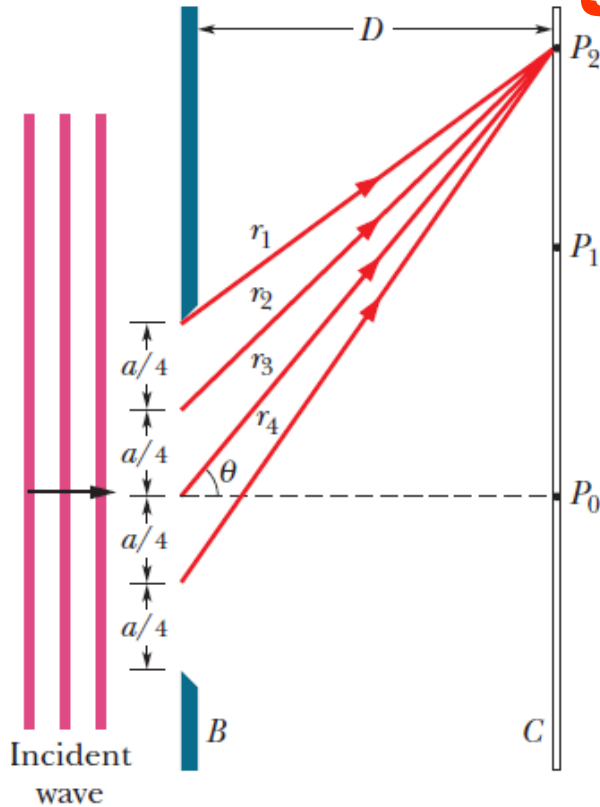
# Diffraction by a Single Slit: Locating the first minimum



$$\frac{a}{2} \sin \theta = \frac{\lambda}{2} \rightarrow a \sin \theta = \lambda$$

(first minimum)

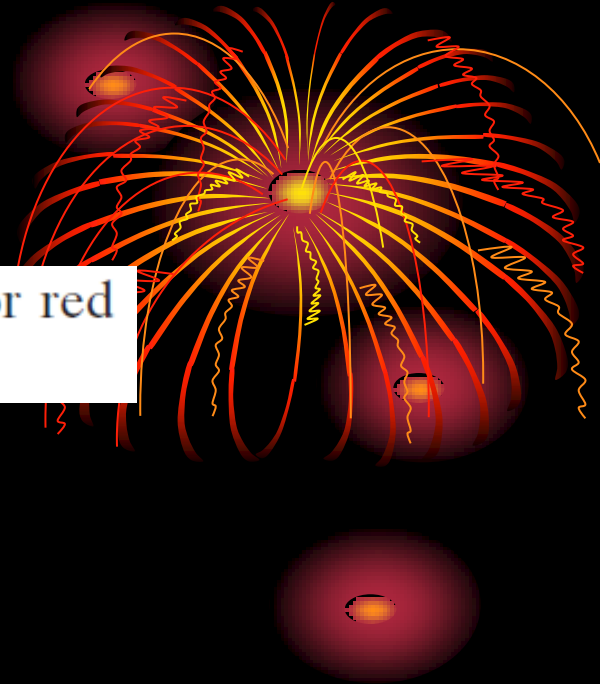
# Diffraction by a Single Slit: Locating the Minima



$$\frac{a}{4} \sin \theta = \frac{\lambda}{2} \quad (a) \quad \rightarrow \quad a \sin \theta = 2\lambda \quad (\text{second minimum})$$

$$a \sin \theta = m\lambda, \quad \text{for } m = 1, 2, 3, \dots \quad (\text{minima-dark fringes})$$

# Ex.11-6 36-1 Slit width



(a) For what value of  $a$  will the first **minimum** for red light of wavelength  $\lambda = 650 \text{ nm}$  appear at  $\theta = 15^\circ$ ?

$$a = \frac{m\lambda}{\sin \theta} = \frac{(1)(650 \text{ nm})}{\sin 15^\circ} \\ = 2511 \text{ nm} \approx 2.5 \mu\text{m}.$$

(b) What is the wavelength  $\lambda'$  of the light whose first side diffraction **maximum** is at  $15^\circ$ , thus coinciding with the first minimum for the red light?

$$a \sin \theta = 1.5\lambda'.$$

$$\lambda' = \frac{a \sin \theta}{1.5} = \frac{(2511 \text{ nm})(\sin 15^\circ)}{1.5} \\ = 430 \text{ nm}.$$

# Intensity in Single-Slit Diffraction, Qualitatively

$$\left( \begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right) = \left( \frac{2\pi}{\lambda} \right) \left( \begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right) \quad \Delta\phi = \left( \frac{2\pi}{\lambda} \right) (\Delta x \sin \theta)$$

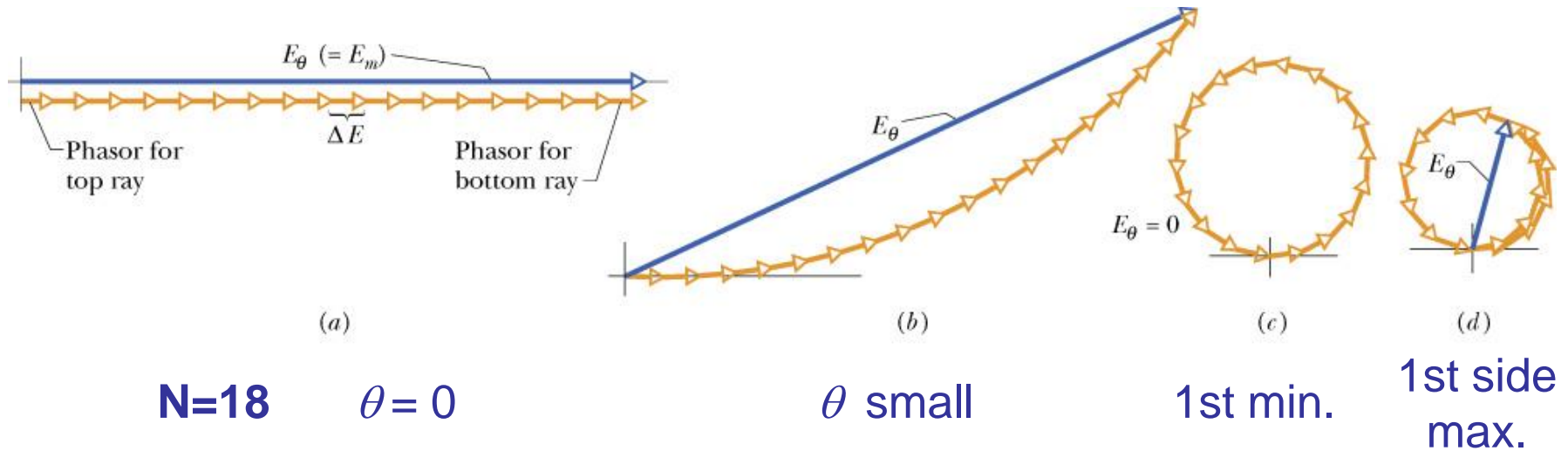


Fig. 36-7



# Intensity and path length difference

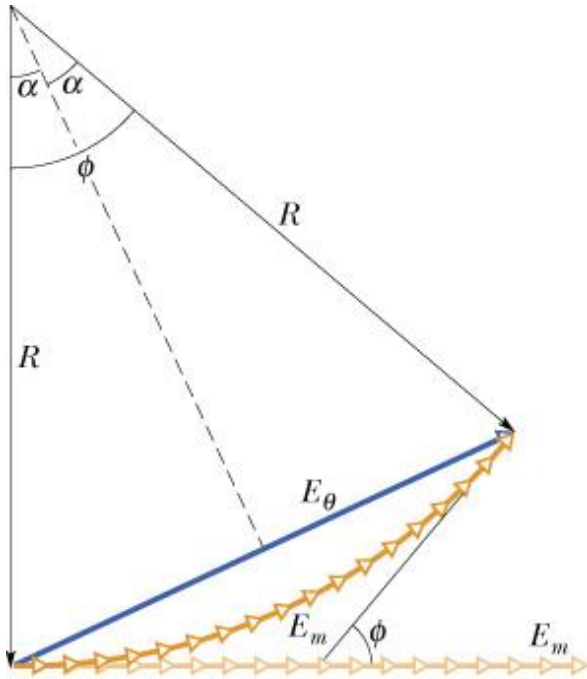


Fig. 36-9

$$\sin \frac{1}{2} \phi = \frac{E_{\theta}}{2R} \quad \phi = \frac{E_m}{R}$$

$$E_{\theta} = \frac{E_m}{\frac{1}{2} \phi} \sin \frac{1}{2} \phi$$

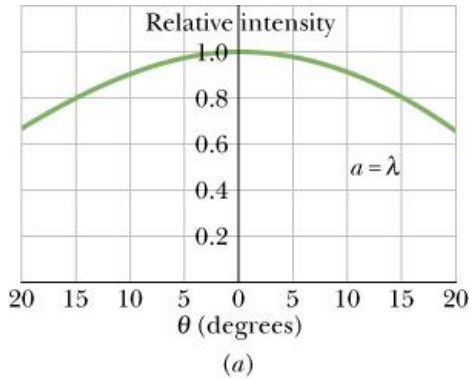
$$\frac{I(\theta)}{I_m} = \frac{E_{\theta}^2}{E_m^2} \rightarrow I(\theta) = I_m \left( \frac{\sin \alpha}{\alpha} \right)^2$$



$$\phi = \left( \frac{2\pi}{\lambda} \right) (a \sin \theta)$$



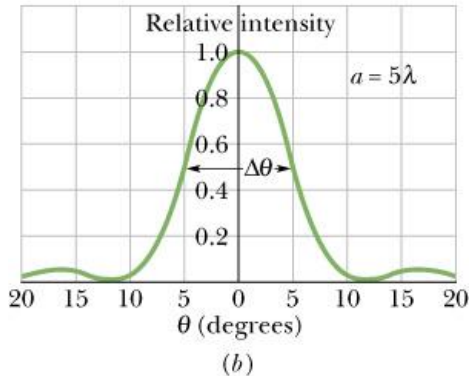
# Intensity in Single-Slit Diffraction, Quantitatively



Here we will show that the intensity at the screen due to a single slit is:

$$I(\theta) = I_m \left( \frac{\sin \alpha}{\alpha} \right)^2 \quad (36-5)$$

$$\text{where } \alpha = \frac{1}{2} \phi = \frac{\pi a}{\lambda} \sin \theta \quad (36-6)$$



In Eq. 36-5, minima occur when:

$$\alpha = m\pi, \quad \text{for } m = 1, 2, 3, \dots$$

If we put this into Eq. 36-6 we find:

$$m\pi = \frac{\pi a}{\lambda} \sin \theta, \quad \text{for } m = 1, 2, 3, \dots$$

$$\text{or } a \sin \theta = m\lambda, \quad \text{for } m = 1, 2, 3, \dots$$

(minima-dark fringes)

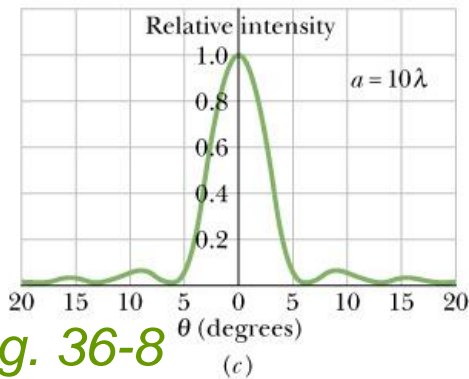


Fig. 36-8

# Ex.11-7 36-2

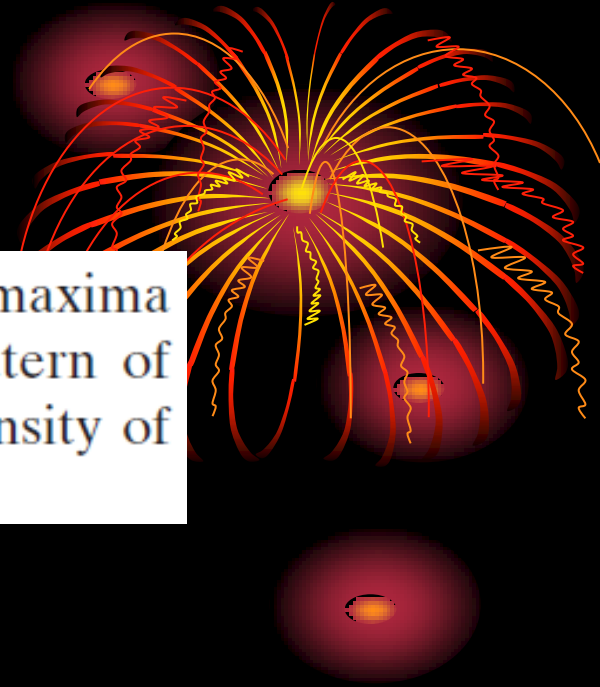
Find the intensities of the first three secondary maxima (side maxima) in the single-slit diffraction pattern of Fig. 36-1, measured as a percentage of the intensity of the central maximum.

$$\alpha = \left( m + \frac{1}{2} \right) \pi, \quad m = 1, 2, 3, \dots$$

$$\frac{I}{I_m} = \left( \frac{\sin \alpha}{\alpha} \right)^2 = \left( \frac{\sin(m + \frac{1}{2})\pi}{(m + \frac{1}{2})\pi} \right)^2, \quad \text{for } m = 1, 2, 3, \dots$$

$$\begin{aligned} \frac{I_1}{I_m} &= \left( \frac{\sin(1 + \frac{1}{2})\pi}{(1 + \frac{1}{2})\pi} \right)^2 = \left( \frac{\sin 1.5\pi}{1.5\pi} \right)^2 \\ &= 4.50 \times 10^{-2} \approx 4.5\%. \end{aligned}$$

$$\frac{I_2}{I_m} = 1.6\% \quad \text{and} \quad \frac{I_3}{I_m} = 0.83\%$$



# Diffraction by a Circular Aperture

Distant point source, e.g., star

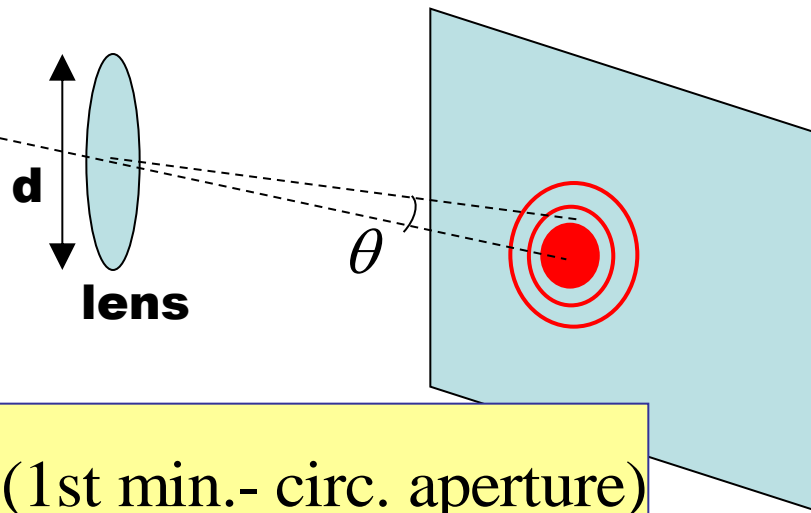
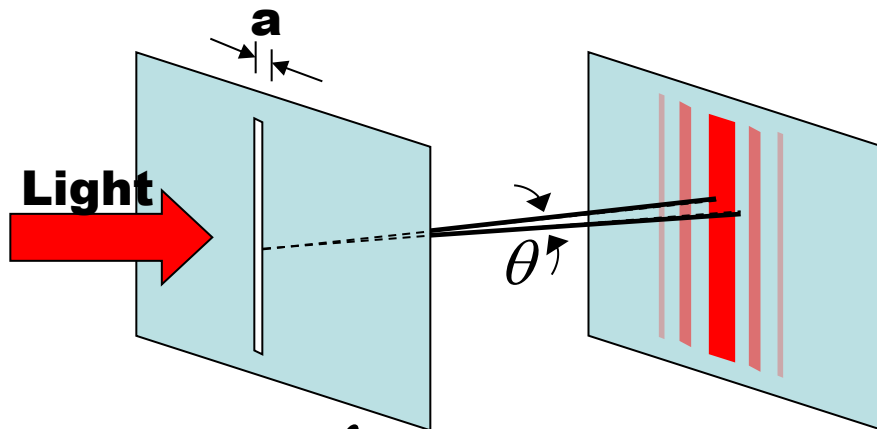
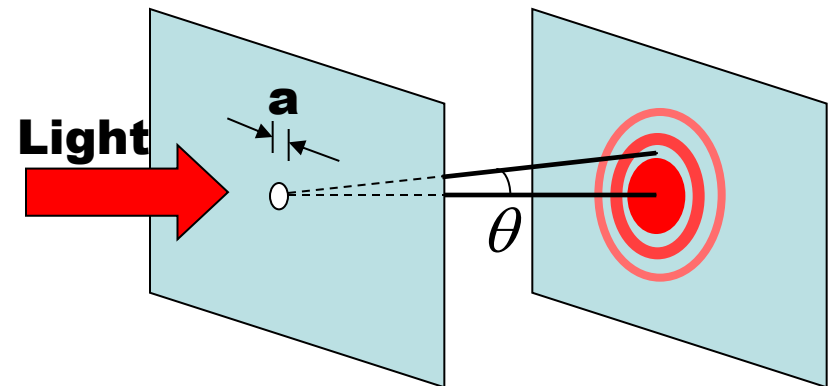


Image is not a point, as expected from geometrical optics! Diffraction is responsible for this image pattern

$$\sin \theta = 1.22 \frac{\lambda}{d} \text{ (1st min.- circ. aperture)}$$

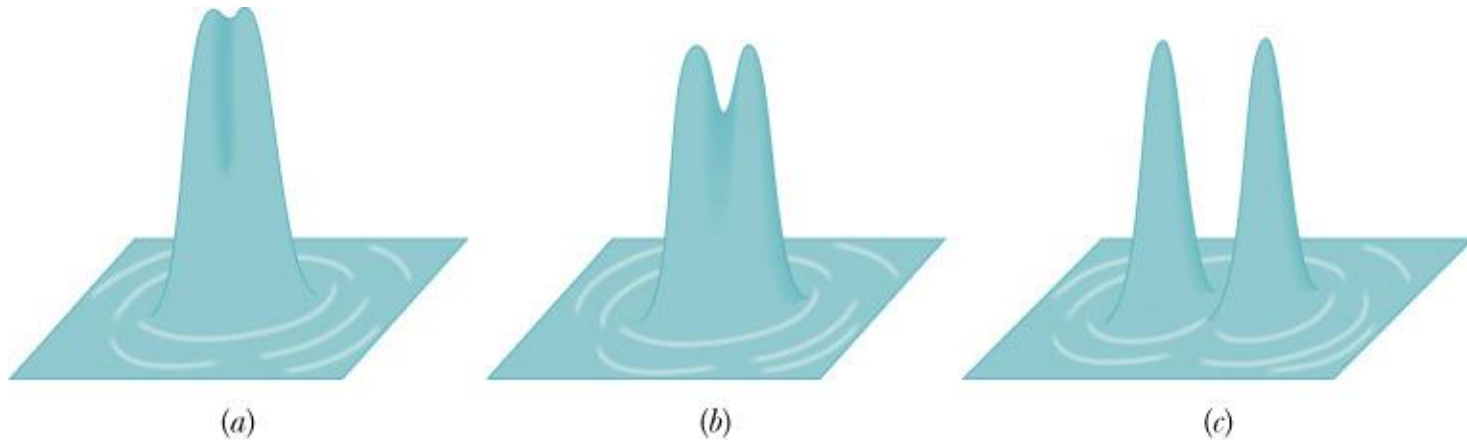


$$\sin \theta = 1.22 \frac{\lambda}{a} \text{ (1st min.- single slit)}$$



# Resolvability

Rayleigh's Criterion: two point sources are barely resolvable if their angular separation  $\theta_R$  results in the central maximum of the diffraction pattern of one source's image is centered on the first minimum of the diffraction pattern of the other source's image.

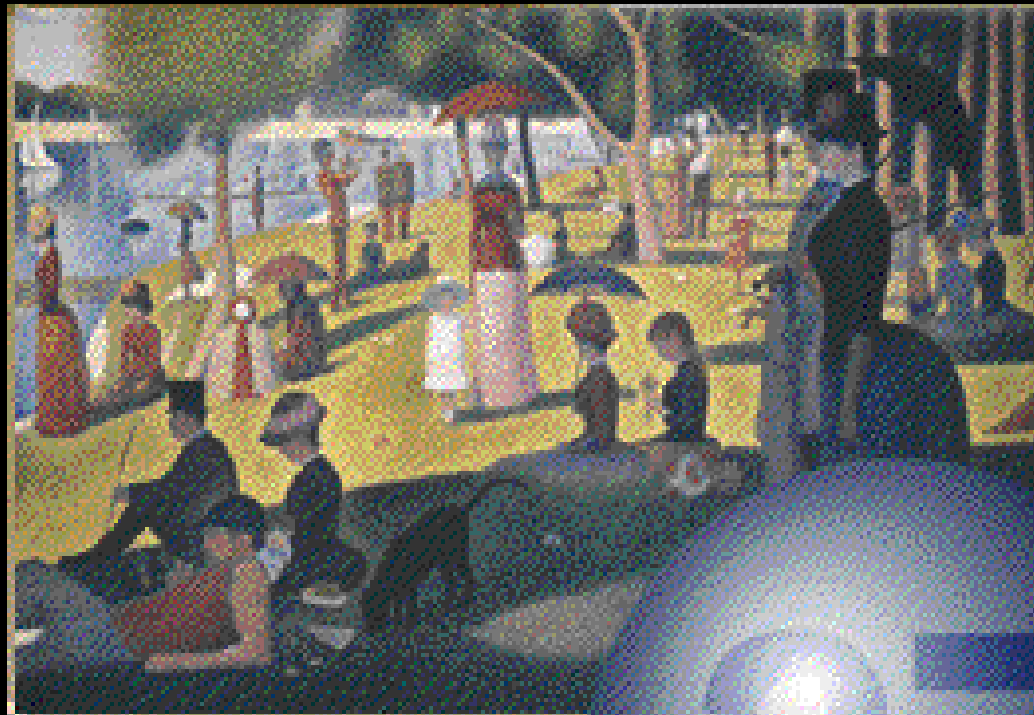


$$\theta_R = \sin^{-1} \left( 1.22 \frac{\lambda}{d} \right) \overset{\theta_R \text{ small}}{\approx} 1.22 \frac{\lambda}{d} \quad (\text{Rayleigh's criterion})$$

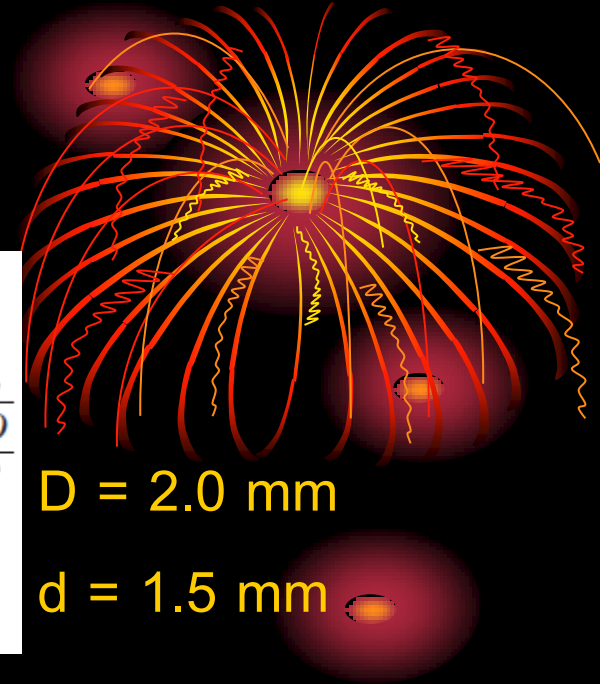
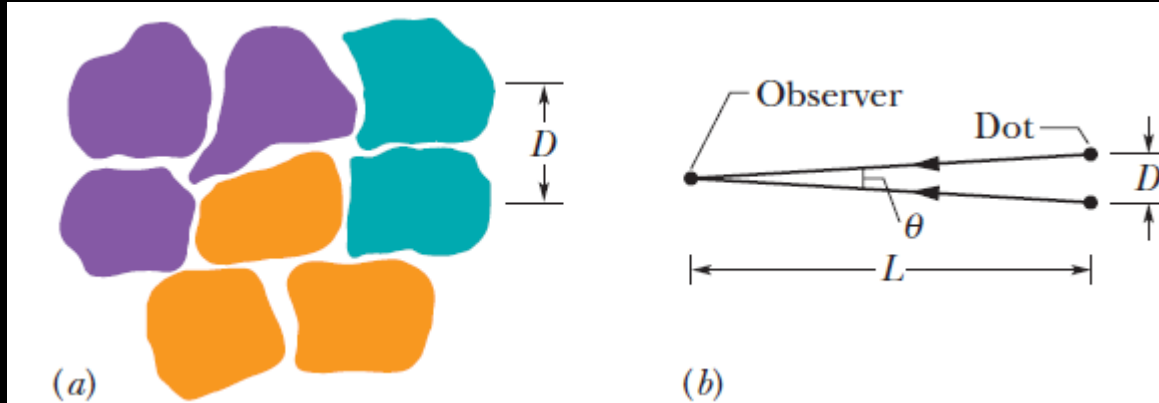
## 11-4.9 Diffraction - (繞射)



Why do the colors in a pointillism painting change with viewing distance?



# Ex.11-8 36-3 pointillism



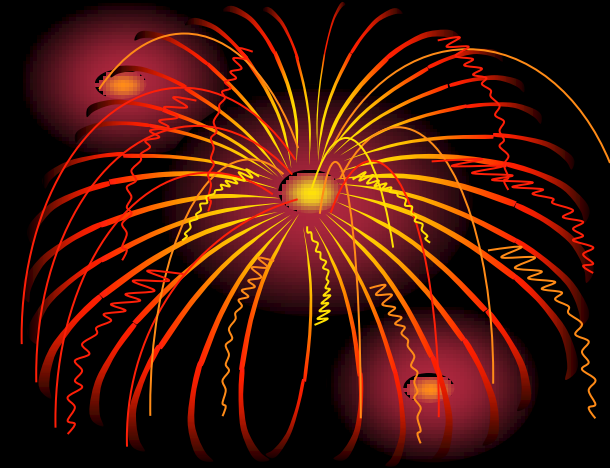
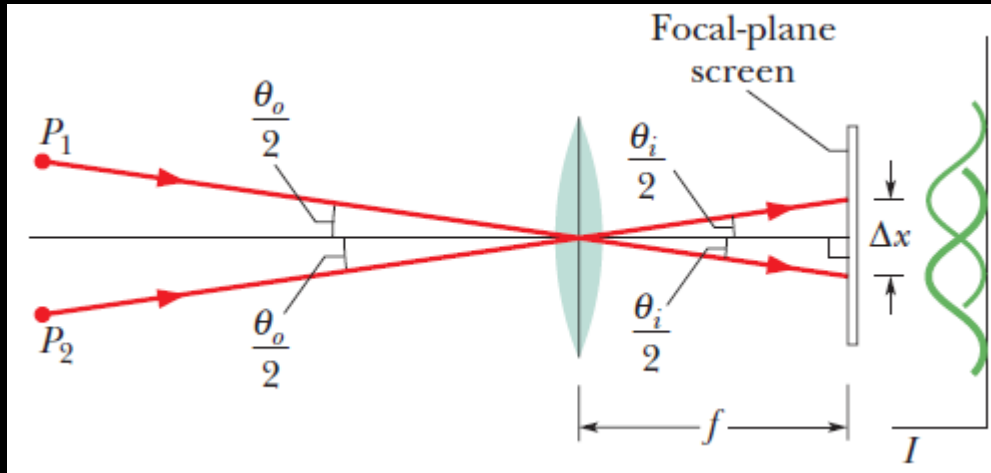
$$\theta_R = 1.22 \frac{\lambda}{d} \quad \theta = \frac{D}{L}$$

$$L = \frac{Dd}{1.22\lambda}$$

$$L = \frac{(2.0 \times 10^{-3} \text{ m})(1.5 \times 10^{-3} \text{ m})}{(1.22)(400 \times 10^{-9} \text{ m})} = 6.1 \text{ m.}$$



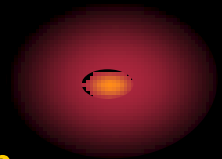
# Ex.11-9 36-4



$$d = 32 \text{ mm}$$

$$f = 24 \text{ cm}$$

$$\lambda = 550 \text{ nm}$$



$$\begin{aligned} \theta_o = \theta_i = \theta_R &= 1.22 \frac{\lambda}{d} \\ &= \frac{(1.22)(550 \times 10^{-9} \text{ m})}{32 \times 10^{-3} \text{ m}} = 2.1 \times 10^{-5} \text{ rad.} \end{aligned}$$

$$\Delta x = f\theta_i, \quad \Delta x = (0.24 \text{ m})(2.1 \times 10^{-5} \text{ rad}) = 5.0 \mu\text{m.}$$

# The telescopes on some commercial and military surveillance satellites

Resolution of 85 cm and 10 cm respectively



$$\frac{D}{L} = \theta_R = 1.22 \frac{\bullet}{d}$$

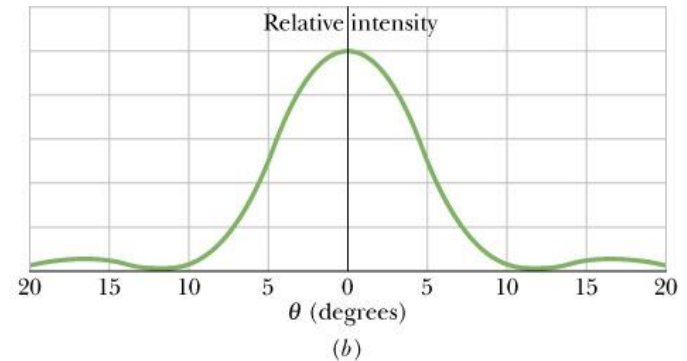
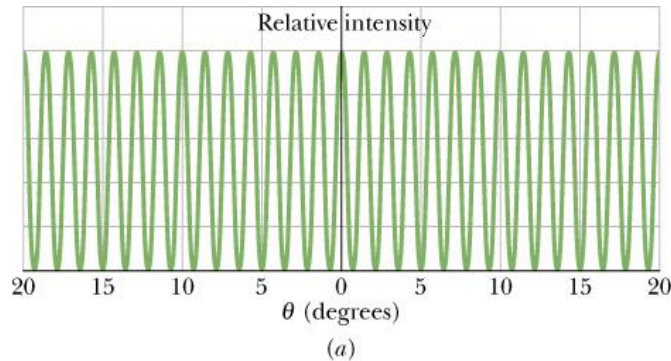
$$\lambda = 550 \times 10^{-9} \text{ m.}$$

$$(a) L = 400 \times 10^3 \text{ m} \quad , \quad D = 0.85 \text{ m} \rightarrow d = 0.32 \text{ m.}$$

$$(b) D = 0.10 \text{ m} \rightarrow d = 2.7 \text{ m.}$$

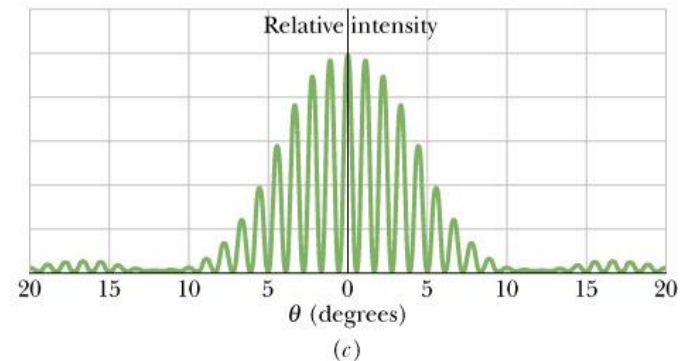
# Diffraction by a Double Slit

Single slit  $a \sim \lambda$



Two vanishingly narrow slits  $a \ll \lambda$

Two Single slits  $a \sim \lambda$

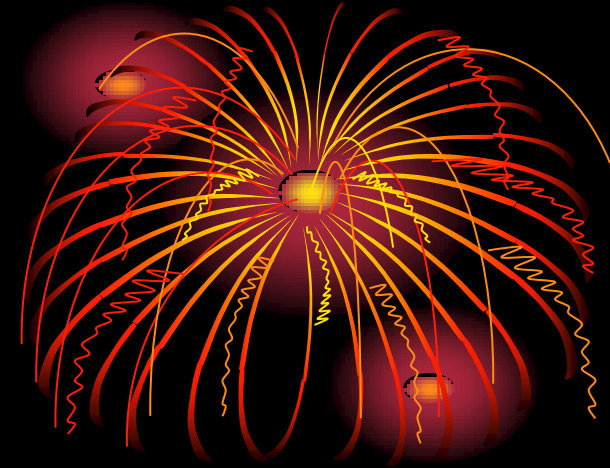
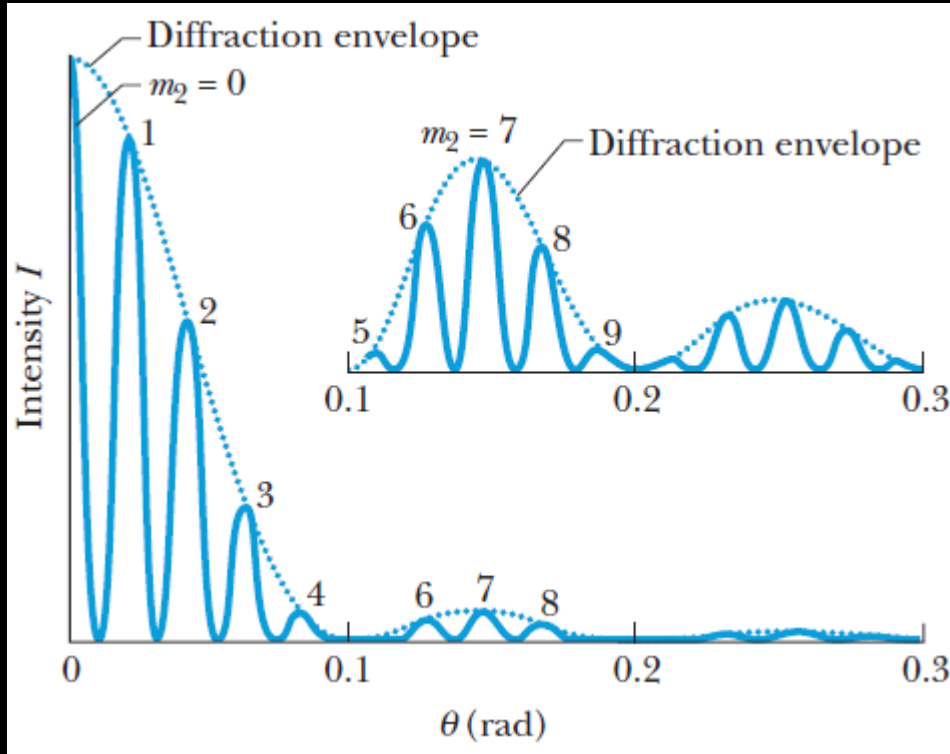


$$I(\theta) = I_m \left( \cos^2 \beta \right) \left( \frac{\sin \alpha}{\alpha} \right)^2 \quad (\text{double slit})$$

$$\beta = \frac{\pi d}{\lambda} \sin \theta$$

$$\alpha = \frac{\pi a}{\lambda} \sin \theta$$

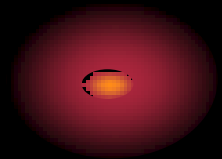
# Ex.11-10 36-5



$$d = 32 \mu\text{m}$$

$$a = 4.050 \mu\text{m}$$

$$\lambda = 405 \text{ nm}$$



$$a \sin \theta = \lambda$$

$$d \sin \theta = m_2 \lambda \quad \text{for } m = 0, 1, 2, \dots$$

$$a \sin \theta = 2\lambda.$$

$$m_2 = \frac{d}{a} = \frac{19.44 \mu\text{m}}{4.050 \mu\text{m}} = 4.8. \quad m_2 = \frac{2d}{a} = \frac{(2)(19.44 \mu\text{m})}{4.050 \mu\text{m}} = 9.6.$$

# Diffraction Gratings

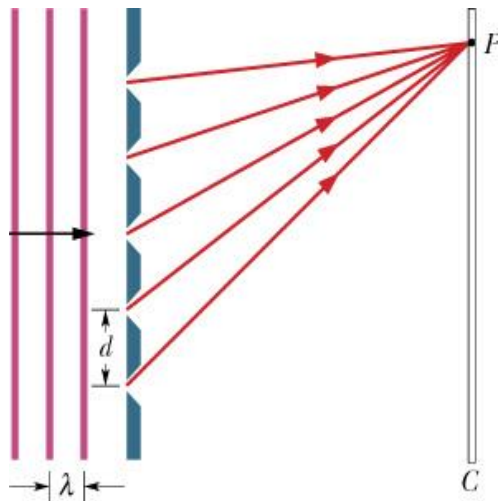


Fig. 36-18

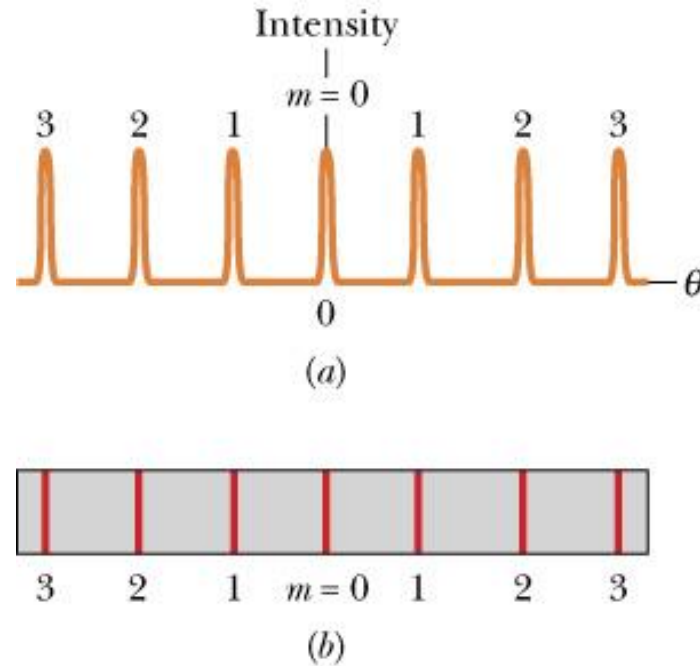


Fig. 36-19

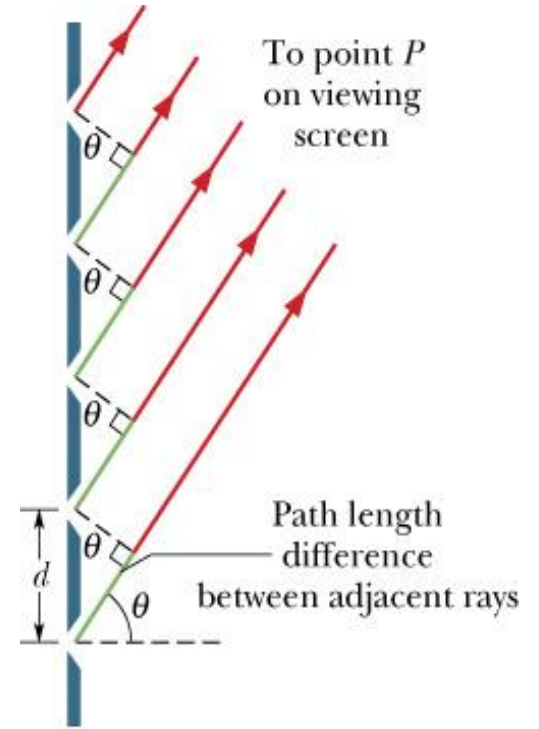


Fig. 36-20

$$d \sin \theta = m\lambda \quad \text{for } m = 0, 1, 2, \dots \quad (\text{maxima-lines})$$

# Width of Lines

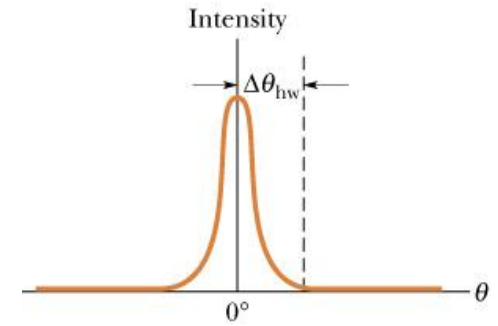


Fig. 36-21

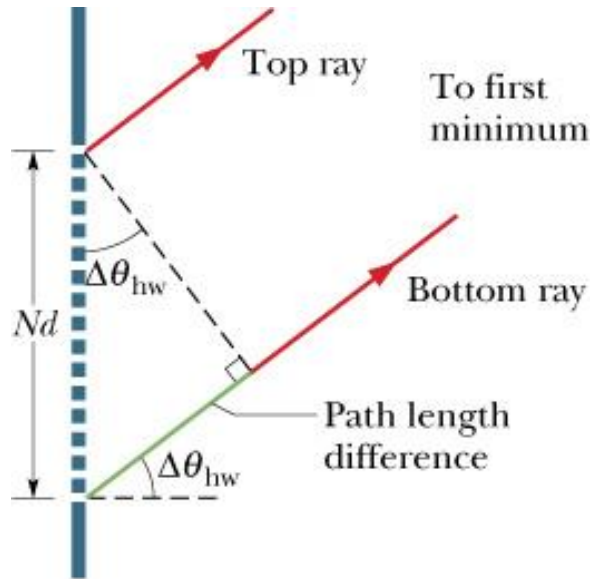


Fig. 36-22

$$Nd \sin \Delta\theta_{hw} = \lambda \quad , \quad \sin \Delta\theta_{hw} \approx \Delta\theta_{hw}$$

$$\Delta\theta_{hw} = \frac{\lambda}{Nd} \quad (\text{half width of central line})$$

$$\Delta\theta_{hw} = \frac{\lambda}{Nd \cos \theta} \quad (\text{half width of line at } \theta)$$

# Grating Spectroscopy

Separates different wavelengths (colors) of light into distinct diffraction lines

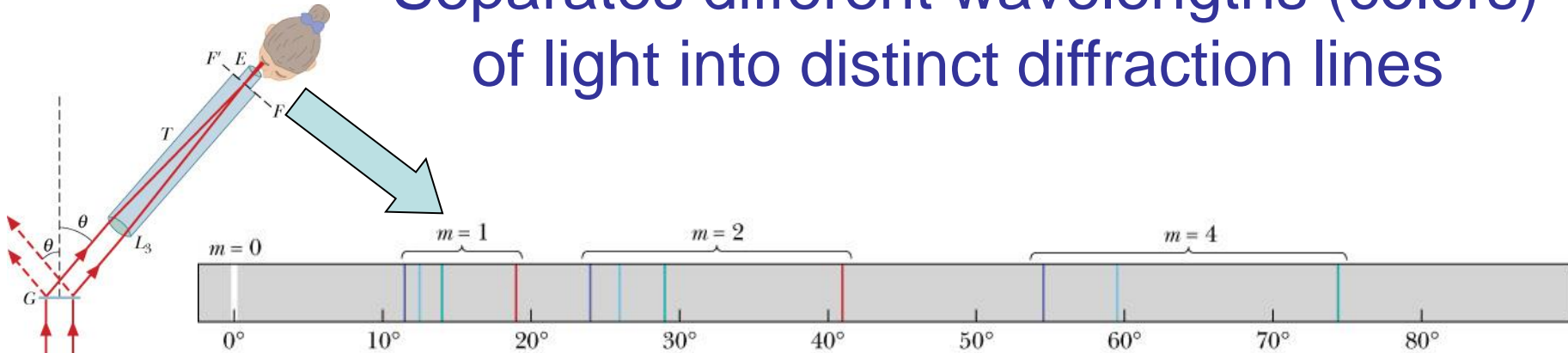
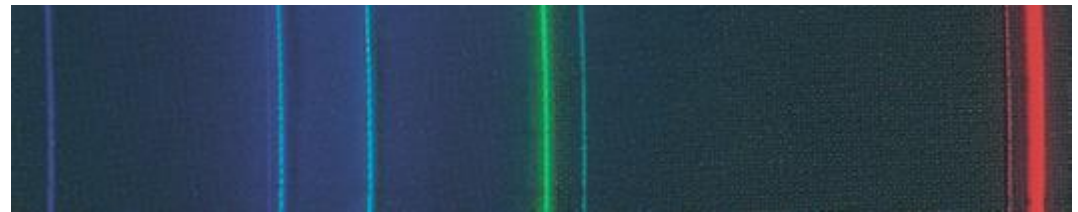


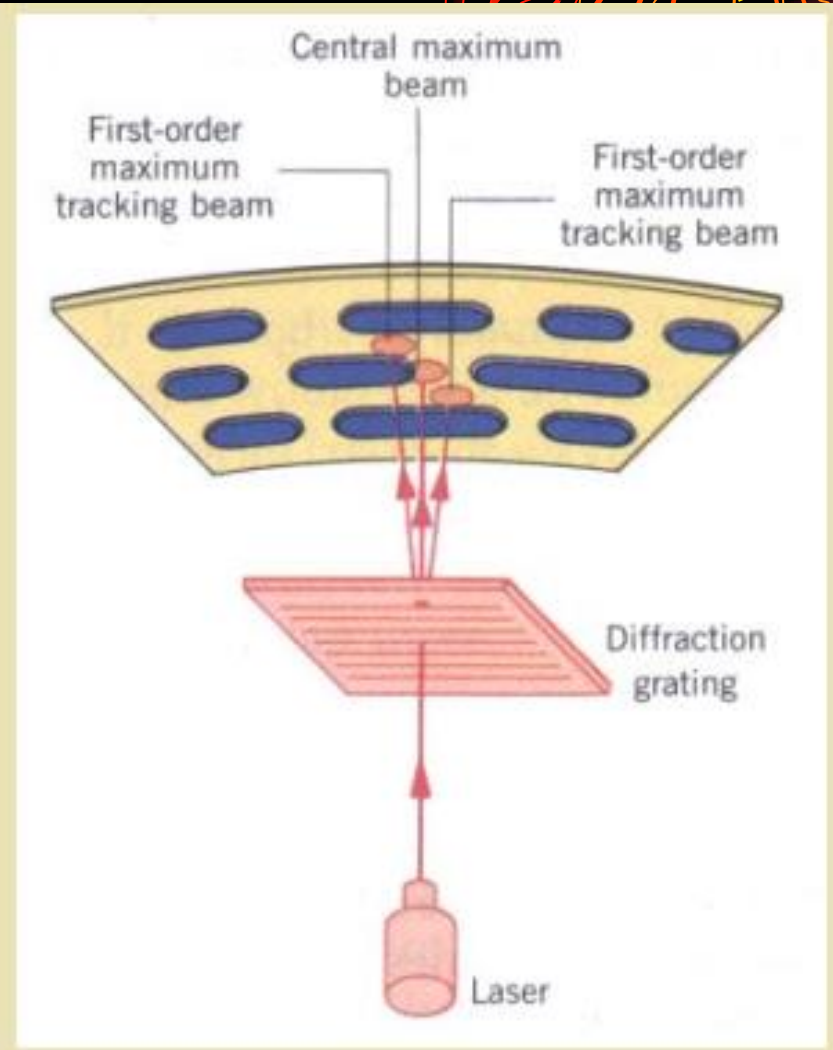
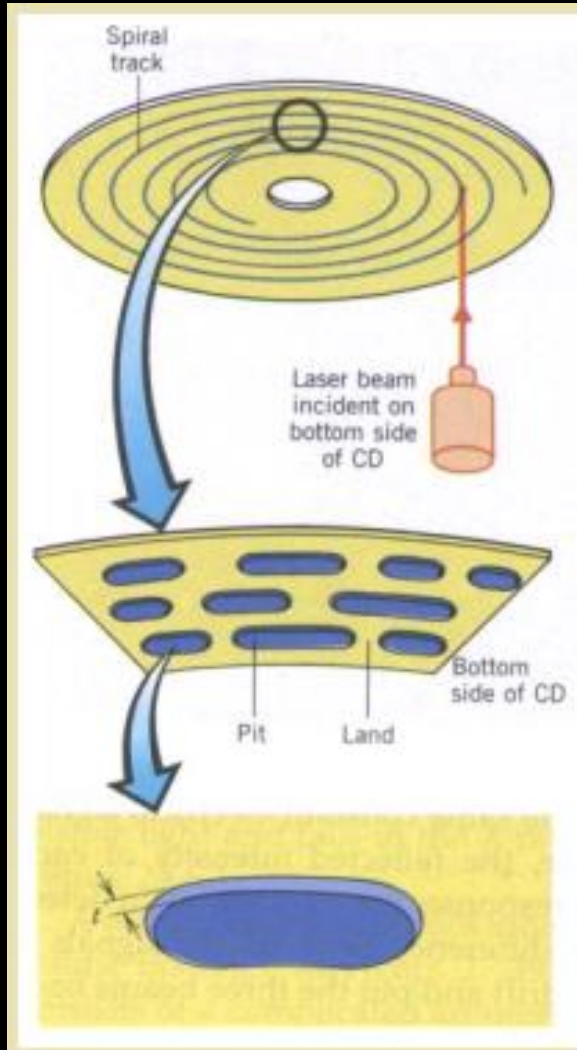
Fig. 36-24

Fig. 36-23



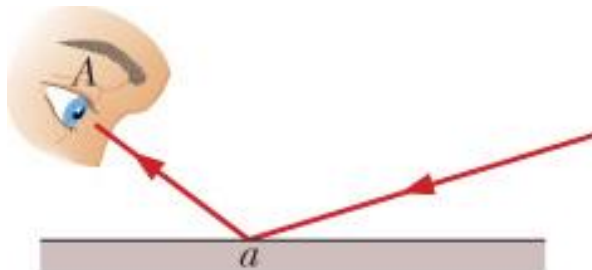


# Compact Disc

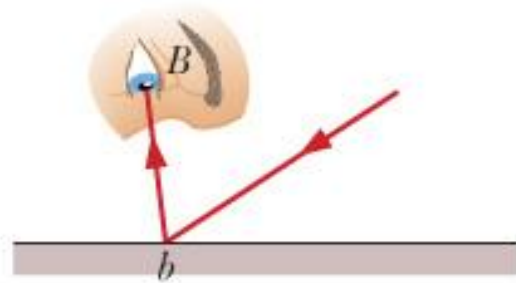




# Optically Variable Graphics



(a)

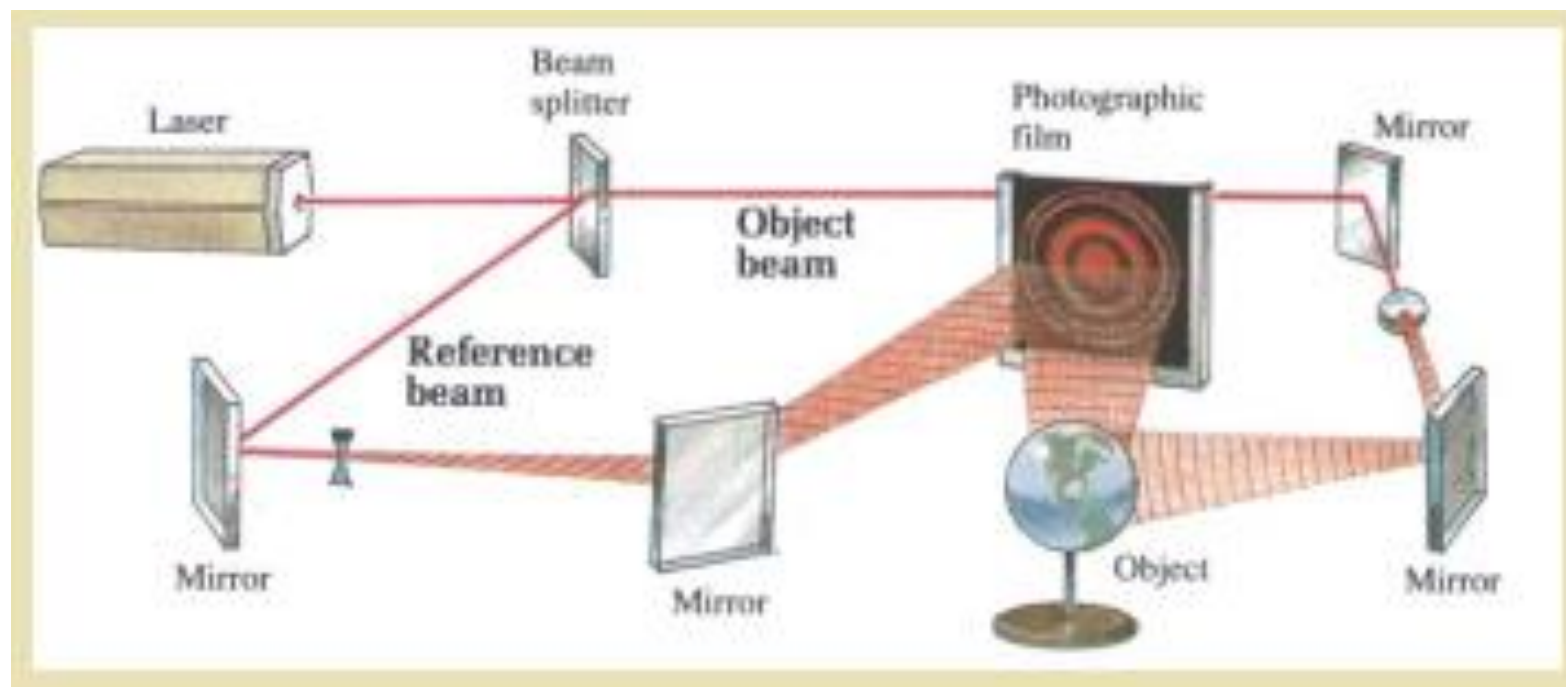


(b)

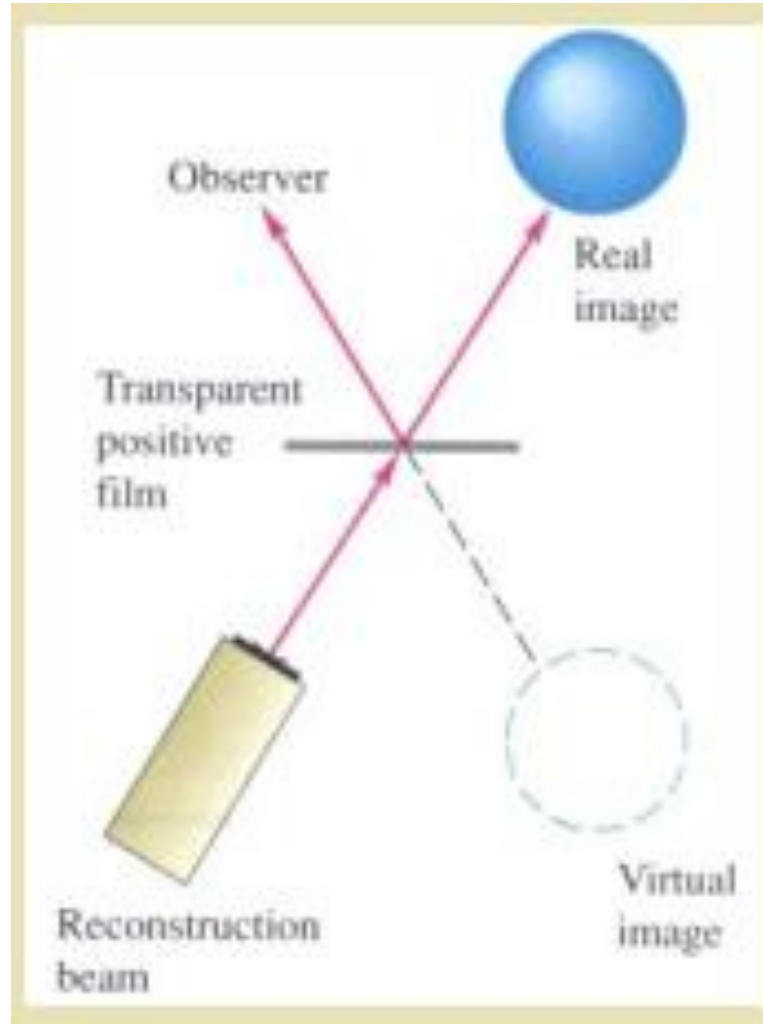


Fig. 36-27

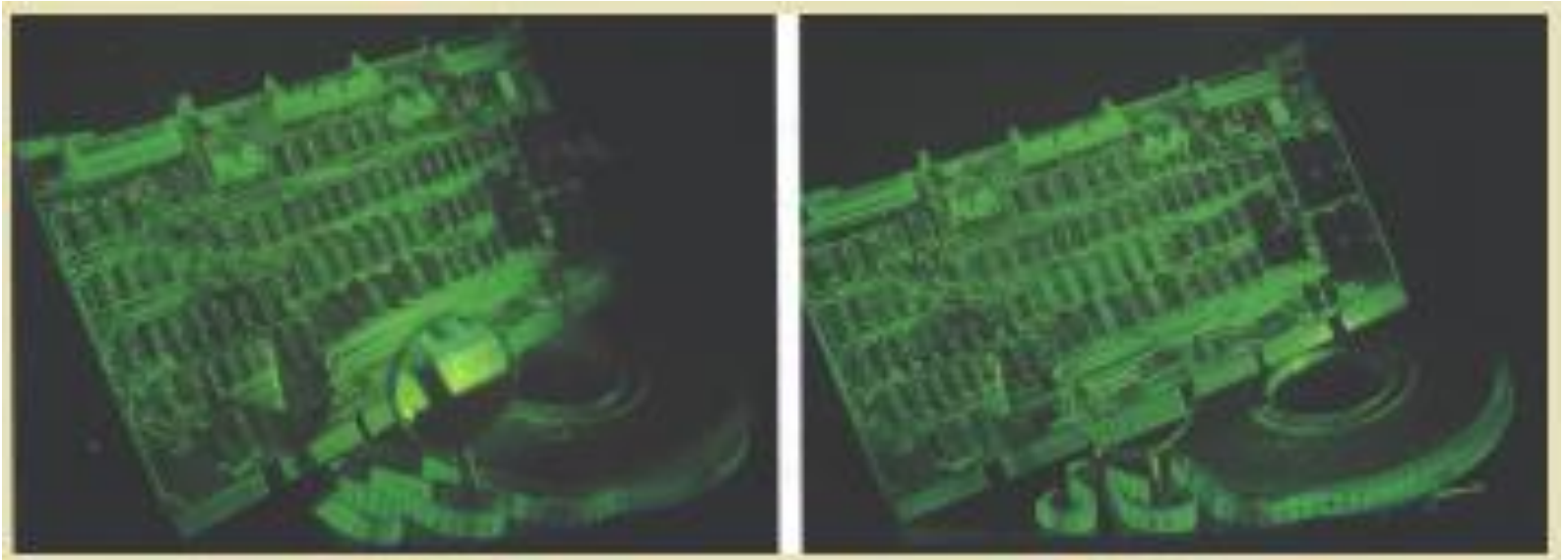
# 全像術



# Viewing a holograph



# A Holograph



# Gratings: Dispersion

$$D = \frac{\Delta\theta}{\Delta\lambda} \quad (\text{dispersion defined})$$

$$D = \frac{m}{d \cos \theta} \quad (\text{dispersion of a grating}) \quad (36-30)$$

Angular position of maxima  $d \sin \theta = m\lambda$

Differential of first equation  
(what change in angle  
does a change in  
wavelength produce?)  $d(\cos \theta) d\theta = m d\lambda$

For small angles

$$d\theta \rightarrow \Delta\theta \quad \text{and} \quad d\lambda \rightarrow \Delta\lambda$$

$$d(\cos \theta) \Delta\theta = m \Delta\lambda$$

$$\frac{\Delta\theta}{\Delta\lambda} = \frac{m}{d(\cos \theta)}$$



# Gratings: Resolving Power

$$R = \frac{\lambda_{\text{avg}}}{\Delta\lambda} \quad (\text{resolving power defined})$$

$$R = Nm \quad (\text{resolving power of a grating}) \quad (36-32)$$

Rayleigh's criterion for half-width to resolve two lines

$$\Delta\theta_{\text{hw}} = \frac{\lambda}{Nd \cos \theta}$$

Substituting for  $\Delta\theta$  in calculation on previous slide

$$\Delta\theta_{\text{hw}} \rightarrow \Delta\theta$$

$$\rightarrow \frac{\lambda}{N} = m\Delta\lambda$$

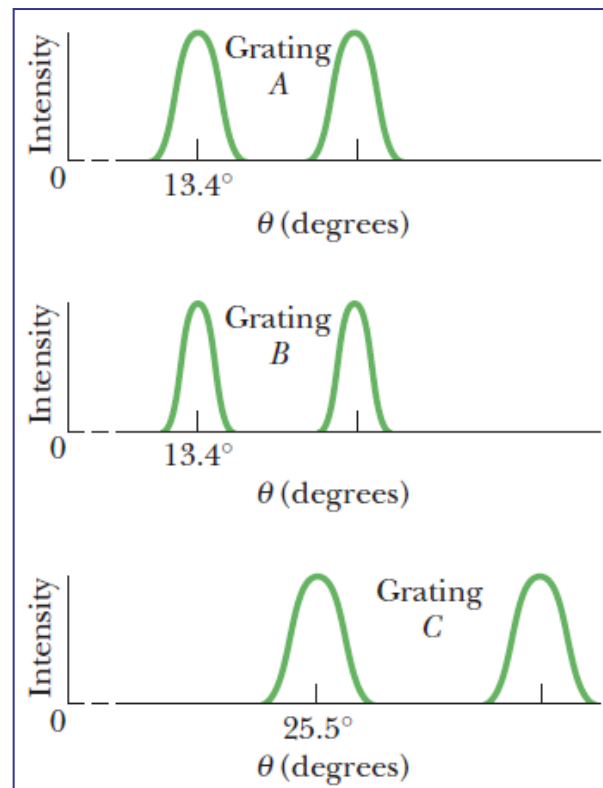
$$R = \frac{\lambda}{\Delta\lambda} = Nm$$



# Dispersion and Resolving Power Compared

Grating	$N$	$d$ (nm)	$\theta$	$D$ ( $^{\circ}/\mu\text{m}$ )	$R$
A	10 000	2540	$13.4^{\circ}$	23.2	10 000
B	20 000	2540	$13.4^{\circ}$	23.2	20 000
C	10 000	1360	$25.5^{\circ}$	46.3	10 000

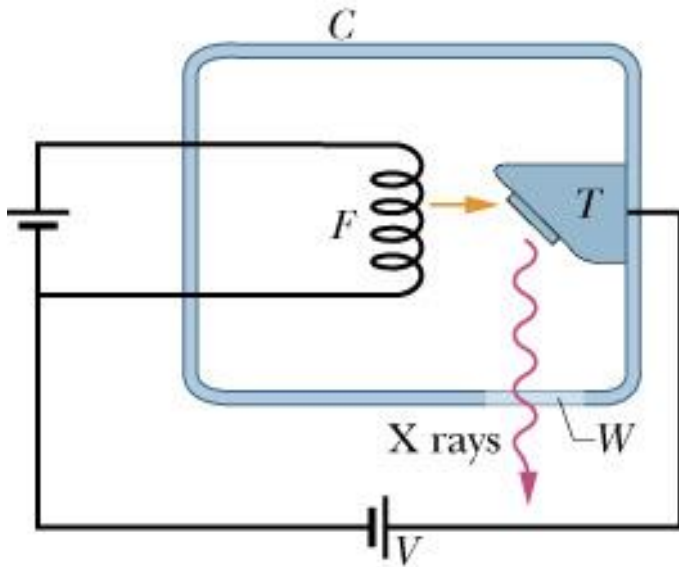
<sup>a</sup>Data are for  $\lambda = 589$  nm and  $m = 1$ .



# X-Ray Diffraction

X-rays are electromagnetic radiation with wavelength  $\sim 1 \text{ \AA}$   
 $= 10^{-10} \text{ m}$  (visible light  $\sim 5.5 \times 10^{-7} \text{ m}$ )

## X-ray generation



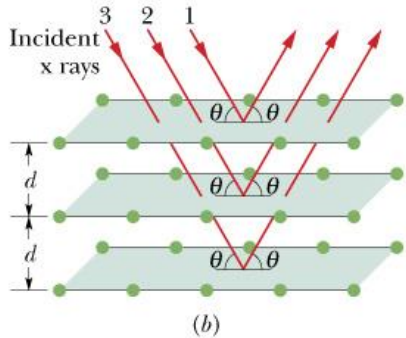
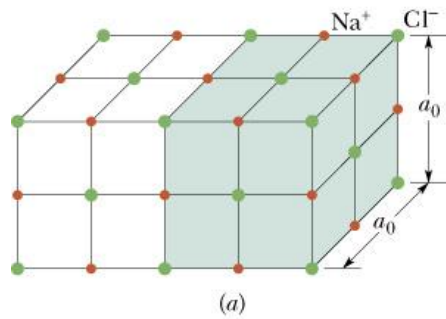
X-ray wavelengths too short to be resolved by a standard optical grating

Fig. 36-29

$$\theta = \sin^{-1} \frac{m\lambda}{d} = \sin^{-1} \frac{(1)(0.1 \text{ nm})}{3000 \text{ nm}} = 0.0019^\circ$$

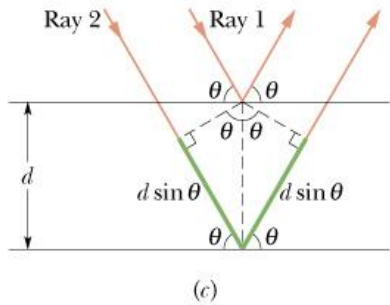


# Diffraction of x-rays by crystal



$d \sim 0.1 \text{ nm}$

→ three-dimensional diffraction grating



$$2d \sin \theta = m\lambda \quad \text{for } m = 0, 1, 2, \dots \quad (\text{Bragg's law})$$

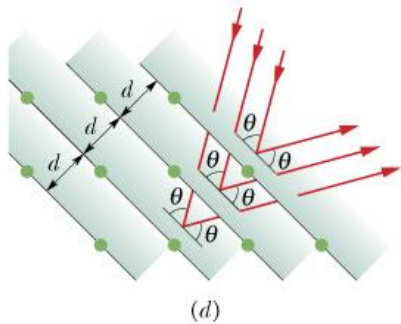
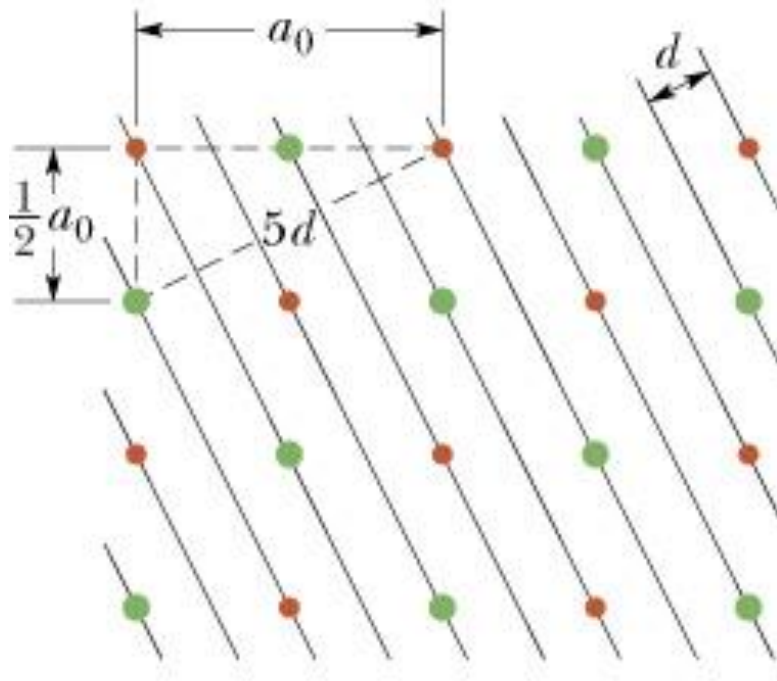


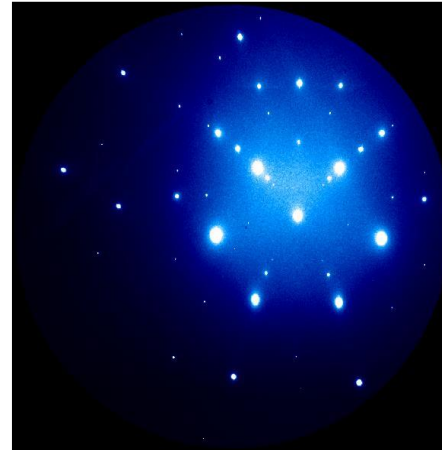
Fig. 36-30

# X-Ray Diffraction, cont'd

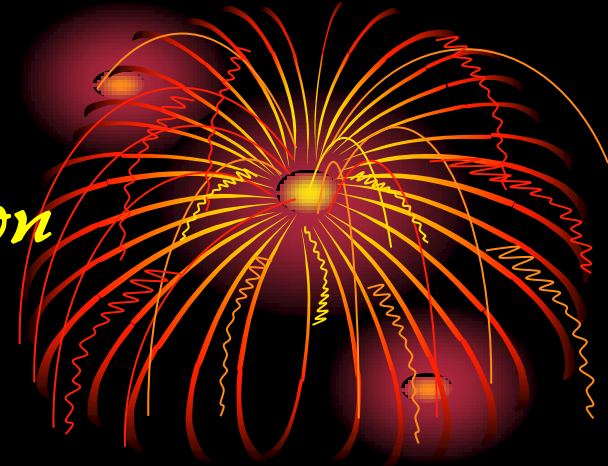


$$5d = \sqrt{\frac{5}{4} a_0^2} \quad \text{or} \quad d = \frac{a_0}{20} = 0.2236a_0$$

Fig. 36-31



# Structural Coloring by Diffraction



(a)

(b)

